

# The Future of US Particle Physics

## Report of the 2021 Snowmass Community Study

### Chapter 5: Cosmic Frontier

Aaron Chou, Marcelle Soares-Santos, Tim M. P. Tait

R. Adhikari, L. A. Anchordoqui, J. Annis, C. Chang, C.L. Chang, J. Cooley, A. Drlica-Wagner, K. Engel, K. Fang, B. Flaugher, J. Jaeckel, T. R. Lewis, T. Lin, H. Lippincott, V. Miranda, L. Newburgh, J. A. Newman, C. Prescod-Weinstein, G. Rybka, B. S. Sathyaprakash, D. Schlegel, D. Shoemaker, T. Slatyer, A. Slosar, K. Tollefson, L. Winslow, H.-B. Yu, T.-T. Yu



## Abstract

We present the Snowmass 2021 strategy to delve deep, search wide, and aim high for new discoveries at the Cosmic Frontier of particle physics. The Cosmic Frontier is the bedrock of High Energy Physics in the twenty-first century, providing evidence for beyond the Standard Model physics that motivates much of the Snowmass 2021 program across all frontiers. The scientific scope of the Cosmic Frontier encompasses four of the five science drivers of the field: dark matter, dark energy and cosmic acceleration, neutrinos, and exploring the unknown in search of new particles, new fields, and new principles of Nature. Covering a time frame of 20 years, our plan features a portfolio of small, medium, and large projects which are ready to produce a continuous stream of groundbreaking science results within this decade along with pathfinders for the next generation of projects to come in the following decade. In this report, we articulate the fundamental questions to be addressed for each science driver and identify key measurements required to achieve well-defined scientific thresholds for discovery. We describe the ecosystem of experiments designed to obtain those measurements as well as the associated developments in theory and technology. We plan to delve deep in sensitivity and search wide across many orders of magnitude in mass to discover the particle nature of dark matter. Moreover, we aim high, through billions of light years of cosmic history, to discover the time-evolution of dark energy, make the first experimental observations of the physics of inflation, and search for new physics at the highest energy scales.

**About Snowmass:** Snowmass is the U.S. High Energy Physics (HEP) Community Planning Exercise organized by the Division of Particles and Fields (DPF) of the American Physical Society (APS). Through the Snowmass process, the community comes together to formulate a 10-year scientific vision for the field within a 20-year global context. The Particle Physics Project Prioritization Panel (P5) takes the scientific input from Snowmass and develops a strategic plan that will inform decisions by the funding agencies. The Cosmic Frontier is one of the major Snowmass sub-communities, focused on cosmic probes of fundamental physics. Cosmic Frontier key topics include: dark matter, dark energy, and inflation. During the Snowmass 2021 process period, from April 2020 to July 2022, the Cosmic Frontier community produced white papers, participated in town-hall meetings and workshops, and gathered input more broadly from all 10 Snowmass frontiers and several related fields. As a result, seven topical reports were produced. The community consensus is consolidated in this frontier-level report, led by the three Cosmic Frontier conveners and co-authored by the topical conveners.

---

---

# Contents

8	<b>5 Cosmic Frontier</b>	<b>5</b>
9	5.1 Executive Summary	5
10	5.2 Delve Deep, Search Wide, Aim High: the Cosmic Frontier Strategy for Discovery	6
11	5.2.1 Fundamental Questions	7
12	5.2.2 State of the Cosmic Frontier	8
13	5.2.3 Science Opportunities	8
14	5.2.4 The Experimental Program	11
15	5.3 Dark Energy	13
16	5.3.1 The Next Stages of Dark Energy Experiments	15
17	5.3.2 A Stage V Spectroscopic Facility (Spec-S5)	16
18	5.3.3 New Opportunity: Gravitational Waves	17
19	5.3.4 Experimental Program & Facilities: Cross-Experiment Strategic Plan	19
20	5.4 Inflation	23
21	5.4.1 Determining the Energy Scale of Inflation	23
22	5.4.2 Reconstructing the Physics of Inflation	23
23	5.4.3 CMB-S4	25
24	5.5 Dark Matter	26
25	5.5.1 Waves of ultralight bosonic dark matter ( $10^{-22}$ eV to 1 eV)	30
26	5.5.2 Light particle dark matter (1 eV to 1 GeV)	36
27	5.5.3 Classic WIMPs (1 GeV to $\sim 100$ TeV)	40
28	5.5.4 Ultra-Heavy Dark Matter ( $> 1$ PeV to $100 M_{\odot}$ )	44
29	5.5.5 Delve Deep, Search Wide	46
30	5.6 Neutrinos	48
31	5.6.1 Cosmic Measurements of Neutrino Masses	49
32	5.6.2 New Opportunity: High-energy neutrinos	49

33	5.7	Exploring the Unknown: New Particles, New Fields, New Principles of Nature . . . . .	50
34	5.7.1	Dark Radiation . . . . .	50
35	5.7.2	The Highest Energy Particles . . . . .	50
36	5.7.3	Probes of Fundamental Physics with Gravitational Waves . . . . .	52

---

---

# Cosmic Frontier

## 5.1 Executive Summary

The Cosmic Frontier realizes the High Energy Physics (HEP) community’s goal of understanding the fundamental physics that governs the Universe and its constituents through a rich scientific research program designed to *delve deep, search wide, and aim high* towards discovery. Cosmic observations are exclusively responsible for our knowledge of the need to extend the Standard Model to describe dark matter, dark energy, and cosmic inflation. As a result, the Cosmic Frontier is uniquely positioned to shed light on these mysteries. The Cosmic Frontier research portfolio features large, medium, and small projects leveraging new developments in both technology and theory to produce a continuous stream of high-impact science over the coming decades.

Understanding the identity, nature, and origin of **dark matter** is one of the grandest challenges in physics and spans all Snowmass frontiers. Well-studied theoretical models provide a compelling scientific case to make broad and rapid inroads into unexplored dark matter parameter space via a *delve deep, search wide* strategy. The HEP community will delve deep to comprehensively explore the high priority science targets of weakly interacting massive particles and the QCD axion. Concurrently, they will also search wide to efficiently probe broad, logarithmic ranges of parameter space that includes hidden sectors and axion-like particles, deploying experimental tools incorporating cosmic probes and cutting edge technologies. Meeting the dark matter challenge will require a broad and coordinated program at multiple scales – including direct, indirect, and cosmic probe experiments – optimized for each decade in dark matter mass. This program of experiments will span many Snowmass frontiers. For example, while dark matter can only be concretely identified via direct observation of its presence in the cosmos, some of its properties can be best studied in controlled laboratory experiments with accelerator test beams or particle colliders. Given the importance of the science and the broad range of experimental opportunities, the overall investment in a coordinated dark matter program should be a priority and at a scale comparable to that of other large experimental programs.

Understanding cosmic acceleration, in both its present-day form of **dark energy** and in the early universe era of **inflation**, is a key science driver for our community. Cosmic surveys comprise experiments in the HEP program designed to study the physics of dark energy and inflation which would otherwise be inaccessible. Cosmic surveys *aim high* at observables spanning almost the entire 13.8 billion year history of our Universe. The current Golden Age of cosmic surveys includes optical/near-infrared survey projects that are under construction or recently commissioned, new cosmic microwave background projects that are actively underway, spectroscopy projects that are ready to be built this decade, as well as next-generation gravitational wave, line-intensity mapping, and cosmic microwave background projects for which pathfinders are now being carried out. Together, these experiments will make precision measurements of dark energy, determine the energy scale and dynamics of cosmic inflation, search for new light relics and other beyond-Standard-Model particles, and probe the physics of dark matter. Cosmic Frontier experiments are able to observe elements of the Standard Model at extreme energies and under extreme environmental conditions which are impossible to achieve in terrestrial laboratories. New messengers such as gravitational waves provide unobscured views of the Universe at its earliest moments, and could reveal new physics through dynamical processes such as early phase transitions which have never been probed to date.

## 5.2 Delve Deep, Search Wide, Aim High: the Cosmic Frontier Strategy for Discovery

The Cosmic Frontier (CF) comprises a broad set of activities aimed at understanding the fundamental physics which governs the evolution of the Universe and its constituents. For Snowmass 2021, the CF is organized into topical groups, each of which produced a report summarizing the open questions and scientific opportunities in its area:

- CF1 : particle-like dark matter [1];
- CF2 : wave-like dark matter [2];
- CF3 : cosmic probes of dark matter [3];
- CF4 : dark energy and cosmic acceleration in the modern universe [4];
- CF5 : dark energy and cosmic acceleration at and before cosmic dawn [5];
- CF6 : new facilities and complementarity in probes of dark energy and cosmic acceleration [6]; and
- CF7 : cosmic probes of fundamental physics [7].

Taken together, the topical working group reports present a broad program of inquiry that offers unique opportunities to understand the physics of dark matter, dark energy, and inflation, and to search for physics beyond the Standard Model (BSM) in a way that is highly complementary to the other Snowmass frontiers.

The Cosmic Frontier engages with all five of the 2014 P5 science drivers, with particular connections to: (i) identify the new physics of dark matter; (ii) understand cosmic acceleration: dark energy and inflation; and (iii) explore the unknown: new particles, interactions, and physical principles. Currently, the incontrovertible experimental evidence for physics beyond the Standard Model of particle physics consists of:

- dark matter;
- dark energy;
- cosmic inflation;
- the baryon asymmetry of the Universe; and
- neutrino masses and mixing.

With the exception of neutrino masses, all of these are the direct result of observational data from the Cosmic Frontier and neutrino masses involved solar neutrinos. Consequently, the techniques of the Cosmic Frontier are the only currently established handle to access several of these physics. The successful road map laid out by P5 in 2014 has resulted in a Golden Age of investigation into fundamental physics through CF projects, which have made great progress exploring the parameter space of dark matter models and pinning down the nature of dark energy. There is strong motivation to build and expand upon that vision to complete exploration of the physics of dark matter, dark energy and cosmic inflation, and to search for new physics that is too massive or too weakly-interacting to access terrestrially. A discovery in any of these areas would be transformational for particle physics, and would suggest new priorities for high energy physics (HEP) research across all the Snowmass frontiers in the future.

### 110 5.2.1 Fundamental Questions

111 CF projects engage with big fundamental questions at the core of the quest to extend/replace the Standard  
112 Model and aligned with the 2014 P5 science drivers. A non-exhaustive list of the central questions includes:

- 113 • What is the nature of dark matter? Is it a single particle or a set of multiple components? Does it  
114 belong to a dark sector?
- 115 • How does dark matter interact with the Standard Model? Does the Higgs serve as a portal between  
116 the two? Are there new bosonic mediators?
- 117 • Is the dark matter itself bosonic and does it exhibit macroscopic wave-like phenomena like oscillatory  
118 forces?
- 119 • What are the dynamics within the dark sector? Does dark matter have important self-interactions?
- 120 • How was the dark matter produced in the Early Universe?
- 121 • Does the nature of dark matter reveal further secrets informing physics beyond the Standard Model  
122 or cosmology?
- 123 • What is the nature of dark energy? Is the acceleration due to a new energy density component or does  
124 it demand a change in general relativity?
- 125 • If dark energy is an energy density, is it a cosmological constant, or a dynamical quantity changing  
126 with time?
- 127 • When did dark energy become important in the history of the universe?
- 128 • Why are there two eras of acceleration in the history of the universe?
- 129 • How is the inflationary paradigm realized in nature?
- 130 • What is the energy scale of inflation?
- 131 • Does inflation have dynamics that manifest themselves as an observable imprint on the primordial  
132 distribution of matter fluctuations?
- 133 • Did BSM degrees of freedom influence the thermal history of the universe?
- 134 • What is the scale of neutrino masses? Can cosmological observations sufficiently constrain it to  
135 distinguish between normal and inverted mass hierarchies?
- 136 • Can gravitational waves reveal new dynamics or early cosmological phase transitions?
- 137 • How can cosmic sources of high energy particles reveal their participation in new interactions or new  
138 physics?

## 5.2.2 State of the Cosmic Frontier

Since the previous Snowmass in 2013, the Cosmic Frontier has made great progress towards realizing the science goals prioritized by P5 in 2014. Searches for WIMP dark matter scattering with nuclei have advanced to generation two, and searches for axions have begun to probe the parameter space relevant to explain the strong CP problem. These searches ‘delve deep’ to study some of the high priority accessible targets in dark matter parameter space. At the same time, new techniques such as those sensitive to electron scattering are able to probe dark matter to sub-MeV masses and ‘search wide’, providing coverage for a wide parameter space consistent with dark matter production in the early universe.

The current U.S. priority on quantum information science has produced fruitful collaborations between HEP and researchers in other fields as ultra-weakly interacting dark matter is a natural science target for quantum sensing. Cross-disciplinary collaborations with the AMO and condensed matter fields have also produced novel dark matter search techniques to provide broad coverage of dark matter parameter space.

Indirect searches for dark matter using the Galactic center, Milky Way dwarf galaxies, and our local Galactic halo as observational targets have produced broad constraints on dark matter mass and self-annihilation cross section. However, despite a strong endorsement from the previous P5, indirect searches for dark matter with gamma-ray observatories have not been prioritized and the U.S. is ceding leadership in this area to other countries.

Investments in cosmological surveys are now coming to fruition – results from the Dark Energy Survey (DES) and BOSS/eBOSS surveys have cemented  $\Lambda$ CDM as a description of the expansion of the Universe, the Dark Energy Spectroscopic Instrument (DESI) is obtaining data at a phenomenal rate, and construction of the Vera C. Rubin Observatory is nearing completion. These experiments ‘aim high’, offering unprecedented opportunity to understand the nature of cosmic acceleration/dark energy as well as providing unique probes of the properties of dark matter such as self-interactions.

Major advances have also been made in Cosmic Microwave Background (CMB) science. The U.S. CMB community has coalesced around a single, large next-generation project, CMB-S4, which will observe the microwave sky with unprecedented statistical precision, targeting the physics of cosmic inflation, searches for BSM dark radiation, and measuring neutrino masses.

The detection of gravitational waves from merging binary neutron stars and black holes by the LIGO and Virgo Collaborations was a watershed moment for cosmic science. This new observational tool will be important for many areas of particle astrophysics and cosmology including measurements of cosmic acceleration at high redshift, searches for dark matter, and detection of relic gravitational wave radiation from early universe phase transitions.

## 5.2.3 Science Opportunities

We identify science opportunities across the broad themes of dark matter detection, identifying the correct model of cosmic inflation, precision cosmology for understanding both dark matter and dark energy, and the use of gravitational waves for these HEP science goals.

- **CMB probes of inflationary cosmology:** Completion of CMB-S4 will probe large classes of models of cosmic inflation and investigate the energy scale and dynamics of this ultra-high-energy sector. In a broad science program, it will also provide high-precision measurements of the matter power spectrum



178 which can be used for a variety of cosmology studies, constrain the neutrino mass scale with possible  
 179 sensitivity to the normal vs. inverted mass hierarchy, and will search for new physics in the form of  
 180 cosmic relic particles.

- 181 • **Direct detection of WIMP dark matter:** The next generation of direct dark matter experiments  
 182 will probe the WIMP paradigm down to and into the neutrino fog.
- 183 • **Direct detection of axion dark matter:** A portfolio of axion dark matter search experiments  
 184 enabled by new quantum sensing technologies will “delve deep” in searches for the ultraweak QCD  
 185 axion signal in most of its predicted band.
- 186 • **Cosmic probes of dark matter:** Cosmological and astrophysical observations, coupled with numer-  
 187 ical simulations and theory, will widely probe and constrain the fundamental nature of dark matter  
 188 over an expansive range of parameter space, some of which is inaccessible to terrestrial and laboratory  
 189 searches. In the coming decade, the Rubin Observatory Legacy Survey of Space and Time (LSST) and  
 190 DESI have the potential to revolutionize our understanding of microscopic properties of dark matter,  
 191 such as mass, self-interactions, interactions with radiation, and quantum wave features. On a longer  
 192 time horizon, the CMB-S4 and Spec-S5 projects have access to rich dark matter physics as well. The  
 193 inclusion of dark matter physics in the research programs of these experiments alongside studies of  
 194 dark energy and inflation should be strongly supported.
- 195 • **Indirect detection of dark matter via high energy astrophysics:** New  $> 100$  GeV gamma  
 196 ray observatories including SWGO and CTA will study dark matter via its annihilations or decays  
 197 in the galaxy, with sensitivity to the thermal-freezeout benchmark annihilation rate up to masses  
 198 approaching the 100 TeV scale, where model-agnostic limits on the early-universe cross section come  
 199 into play. Planned or proposed space missions will set stringent new bounds in the keV-GeV range,  
 200 and new cosmic-ray detectors will perform the first dedicated searches for low-energy antinuclei as a  
 201 background-free discovery channel. Planned and ongoing upgrades to existing observatories including  
 202 IceCube and AugerPrime will provide unique probes of ultra-heavy dark matter.
- 203 • **Multi-pronged detection of thermal dark matter in dark portal scenarios:** A variety of new  
 204 techniques including low-threshold microcalorimetry will be deployed to search for sub-GeV portal  
 205 dark matter. Concurrently, accelerator-based experiments will search for new bosons which serve as  
 206 the messengers between the dark and visible sectors.
- 207 • **Direct detection of ultralow mass dark matter using cross-disciplinary technologies:** En-  
 208 gagement with the fields of quantum information science and atomic, molecular, and optical physics  
 209 has brought new ideas and technologies to bear on searching for ultra-low-mass wave-like dark matter,  
 210 which produce periodic perturbations detectable via ultra-sensitive instruments such as atomic clocks  
 211 and interferometers.
- 212 • **Precision cosmology with existing telescopes:** Cosmic surveys including DESI and LSST will  
 213 provide strongly-constraining data sets for precision cosmology. One percent measurements of the dark  
 214 energy density to  $z \sim 1.5$  will provide crucial tests of dynamical dark energy models. Measurements  
 215 of the overall matter power spectrum will enable tests for new dynamics in the early universe and  
 216 for precision constraints on the total mass of neutrinos. By probing both the growth rate of density  
 217 perturbations and the expansion history of the universe, these experiments will provide critical tests  
 218 of whether cosmic acceleration is in fact due to a dark energy-like component of the Universe or if it  
 219 instead requires a non-Einsteinian theory of gravity on cosmic scales. Increased statistics for Milky  
 220 Way dwarf galaxies and studies of smaller dark halos will provide information on dark matter masses  
 221 and self interactions and also provide new targets for indirect detection searches. For these surveys to  
 222 reach their full potential, funding for innovative science analyses – including cross-survey measurements

– will be needed. New data from other facilities will be needed as a complement to unlock the full constraining power of LSST, including follow-up observations of supernovae and gravitational wave standard sirens as well as measurements of spectroscopic redshifts for deep training samples of objects to enable precision photometric redshift measurements.

- **Stage V spectroscopic probes of cosmic acceleration and inflationary cosmology:** The development of a powerful new Stage V spectroscopic facility will allow precision measurements of the cosmic expansion history to be extended to the limits of the modern universe ( $z \lesssim 5$ ), simultaneously providing new constraints on the inflationary epoch of cosmic acceleration. Accessing high redshifts greatly expands the cosmic volume surveyed, enabling structures on the largest scales where inflationary signatures are expected to manifest to be explored. A Stage V spectroscopic facility also enables a search for new light relics and other beyond the Standard-Model particles including measurement of the sum of the masses of neutrinos, and probes of the physics of dark matter via a variety of precision studies of nearby galaxies and of stars in the Milky Way halo.
- **New pathfinder experiments for Stage VI precision cosmology:** Pathfinder experiments could enable progress on powerful new probes of cosmology such as line intensity mapping methods in 21 cm and millimeter wave molecular lines. Those new probes can extend matter power spectrum studies to still higher redshifts and greater volumes than Stage V spectroscopy can reach. For example, LuSEE-Night is a mission to the far side of the Moon being carried out in partnership between DOE and NASA as a pathfinder for a future Stage VI 21-cm experiment.
- **New experimental windows on fundamental physics:** Pathfinder experiments including multi-messenger cosmology with gravitational waves, new techniques for precision measurements in astrometry, spectroscopy and photometry, and new sky surveys from space at radio frequencies below Earth ionospheric cut-off have potential to unlock hitherto unused probes of fundamental physics through cosmological measurements.
- **Ultra-high-energy cosmic rays as a probe of new physics below the fermi distance:** The new level of precision achieved by AugerPrime, IceCube, and next-generation experiments will provide opportunities to study non-perturbative QCD and physics beyond the Standard Model using natural accelerators well beyond the reach of human-made ones. The already-established anomalous muon production in extensive air showers (which has eluded explanation in models tuned to LHC data) guarantees that discoveries will be made.
- **Gravitational waves as a new probe of particle physics and cosmology:** Gravitational waves have emerged as a new and unique probe of dark sector physics, which may only be minimally coupled to the standard model via gravity. In addition to searching for relic gravitational waves from early-universe phase transitions (including leptogenesis models or the QCD axion), current and future gravitational wave observatories can potentially perform precision measurements of cosmic acceleration through mergers of neutron stars and black holes, standard sirens, which are new standard distance indicators. These large interferometers also have a large effective collection area, enabling searches for heavy dark matter.
- **High-energy neutrinos as a new venue to study neutrino interaction and mixing:** The cosmic neutrinos detected by IceCube provide a new opportunity to learn neutrino physics at TeV-PeV energies. The upcoming IceCube Upgrade and future GeV cosmic neutrino experiments will measure the flavor composition and tau neutrino appearance with high precision. Future >100 PeV neutrino experiments may uncover cosmogenic neutrinos and probe BSM features at extreme energies.

#### 5.2.4 The Experimental Program

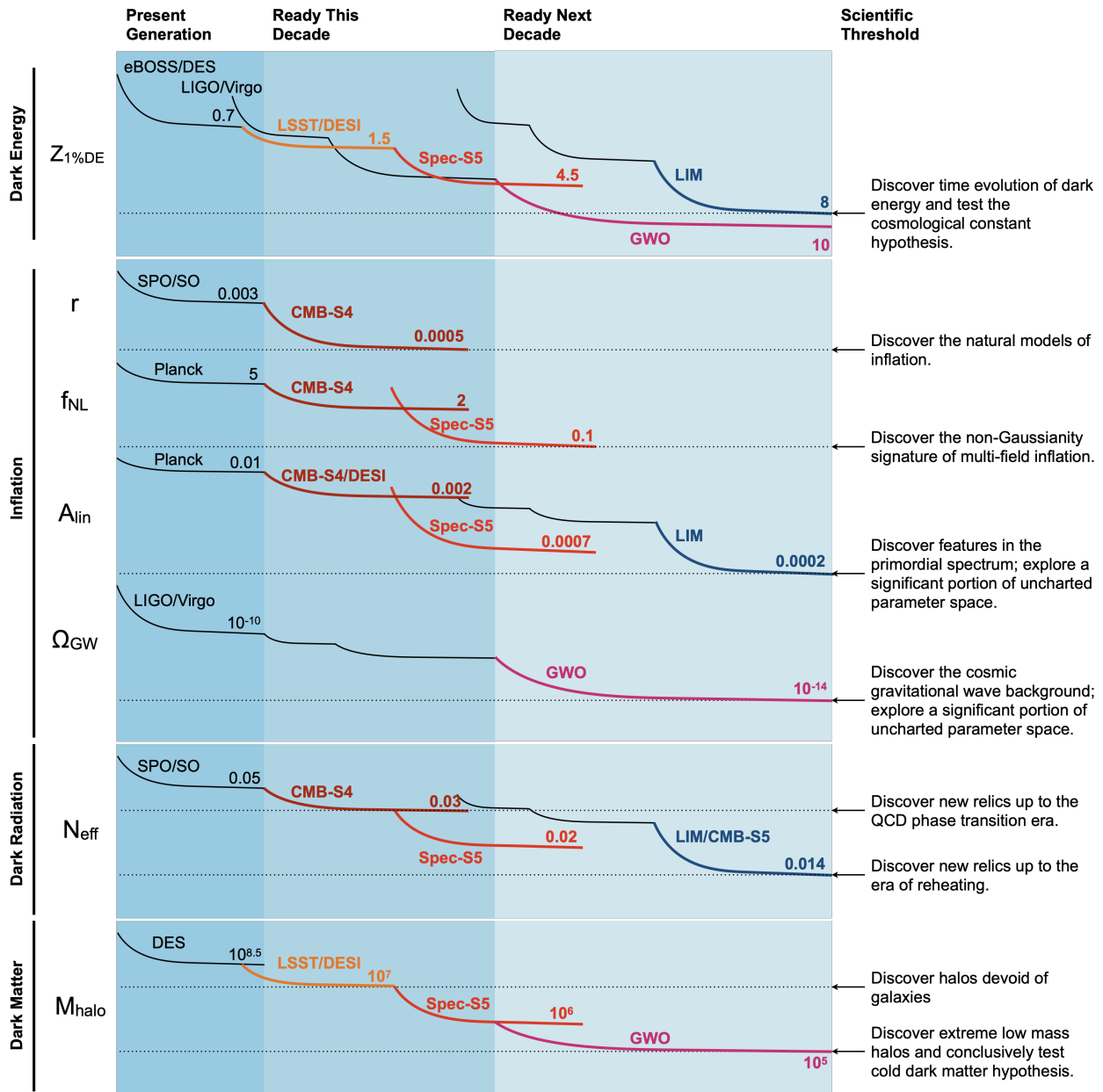
The CF strategy detailed in this report encompasses a science-rich program with small, medium, and large experiments which are planned on both near-term and long-term timescales. The community consensus is that this coming decade must feature strong support for this program in full to ensure that HEP will enjoy the new era of discovery and groundbreaking precision measurements that have been promised for the next two decades, as well as building even more powerful experiments to reach key science thresholds. An overview of the CF program is provided in this section. Details about the projects and activities mentioned here are provided in later sections of this report.

**Cosmic surveys:** Figure 5-1 summarizes the cosmic survey projects and their discovery potential across multiple science areas of our program. For each fundamental question, we have identified a key measurement with a clear threshold for discovery. Together, the planned projects should deliver results spanning large swaths of cosmic frontier science, including studies of cosmic acceleration/dark energy, inflation, searches for indications of new particles (e.g., dark radiation), and exploration of dark matter properties. Cosmic surveys generally provide datasets that can be used for many purposes, addressing multiple HEP science drivers simultaneously. The figure shows a rough timeline for when each major project could come into play and gradually reach its full discovery power. The large CMB-S4 and Spec-S5 projects should play a key role in advances over the next decade and beyond. The figure also includes potential future surveys that will be enabled by R&D and small pathfinder initiatives; these projects would employ new technologies and techniques such as 21-cm and mm-wave Line Intensity Mapping (LIM), Gravitational Wave Observatories (GWO), and future broadband CMB imaging. This planned sequence of experiments ensures that CF will have a continuous stream of data to support scientific breakthroughs without gaps for construction of our largest facilities.

In the Cosmic Frontier survey ecosystem each type of program brings unique strengths to the portfolio, including shape measurements that can only come from imaging, precise and robust redshift measurements that require spectroscopy, and multi-messenger data that exploits new gravitational wave detection capabilities; by also spanning a range in wavelength (e.g., Optical/IR vs. mm-wave vs. 21-cm), information from different redshift ranges can be accessed. While each survey will deliver groundbreaking results on its own, no single experiment can meet all of the discovery thresholds without complementary information from other programs.

As a result, the CF survey program is designed to leverage the strengths of multiple experiments simultaneously. One example of this complementarity is that going beyond discovery of the energy scale of inflation (which CMB-S4 could accomplish on its own) to conclusively probing its physics via precision measurements of the  $f_{\text{NL}}$  parameter, a combination of information from a Stage V Spectroscopic Facility (Spec-S5) with CMB-S4 would be needed. Other science goals such as studying the time evolution of dark energy, conclusively testing the cold dark matter hypothesis, searching for features in the primordial power spectrum, and discovering the cosmic gravitational wave background will require the next-generation of surveys for which pathfinders and R&D work should be supported in this decade. These and many other science opportunities realized by the CF's survey ecosystem have been discussed in detail in the topical group reports [3–6]. In any funding scenario to be considered by the next P5, the Snowmass community consensus is that support for the CF cosmic surveys program should be at a level that enables discovery across the entire scope of the science program.

**Searching for dark matter:** Figure 5-2 shows two snapshots of the key dark matter discovery parameter space, which has greatly expanded since the last Snowmass. They illustrate the broad coverage spanning many decades of dark matter parameters which will be achieved by the ambitious experimental program envisioned in this report, comprising both a broad portfolio of small/medium-scale direct and indirect



**Figure 5-1.** The landscape of cosmic survey projects and their discovery potential. For each major science topic, we identified a key measurement threshold for discovery. Bold lines show medium and large projects ready to produce science in this decade or the next. In particular, CMB-S4 is a large project that is currently at CD-0 stage. Labeled thin lines represent the results from current surveys and unlabelled thin lines are pathfinders for the next generation of projects. Note that the bold line labels Spec-S5, GWO, LIM, and CMB-S5 do not represent specific projects. Instead, they represent categories of future experiments for which multiple proposed projects may currently exist in the community (namely, Stage V Spectroscopic Facility, next-generation gravitational wave observatory, 21cm/mm-wave line intensity mapping, and next-generation CMB, respectively). Please refer to the text for details about these projects and their key measurements.

311 detection experiments optimized for each decade in dark matter mass, as well as an expanded cosmic probes  
312 program to identify the properties of dark matter via cosmic surveys, which is highly synergistic with the  
313 other science these surveys will undertake. The top panel shows current experimental limits via the grey  
314 shaded regions, while the bottom panel shows the projected sensitivity over the same parameter space after  
315 completion of our program. In many cases, utilizing new detector technologies will require cross-disciplinary  
316 collaborations to access the expertise, technology, and facilities of neighboring fields of study such as AMO,  
317 condensed matter physics, quantum information science, and gravitational physics. Similarly, for indirect  
318 probes of dark matter scattering, annihilation, or decay in galactic halos, HEP provides a critical role in  
319 developing instrumentation for astroparticle observatories. Understanding the astrophysical backgrounds  
320 will require collaboration with the broader non-HEP astronomy and astrophysics community.

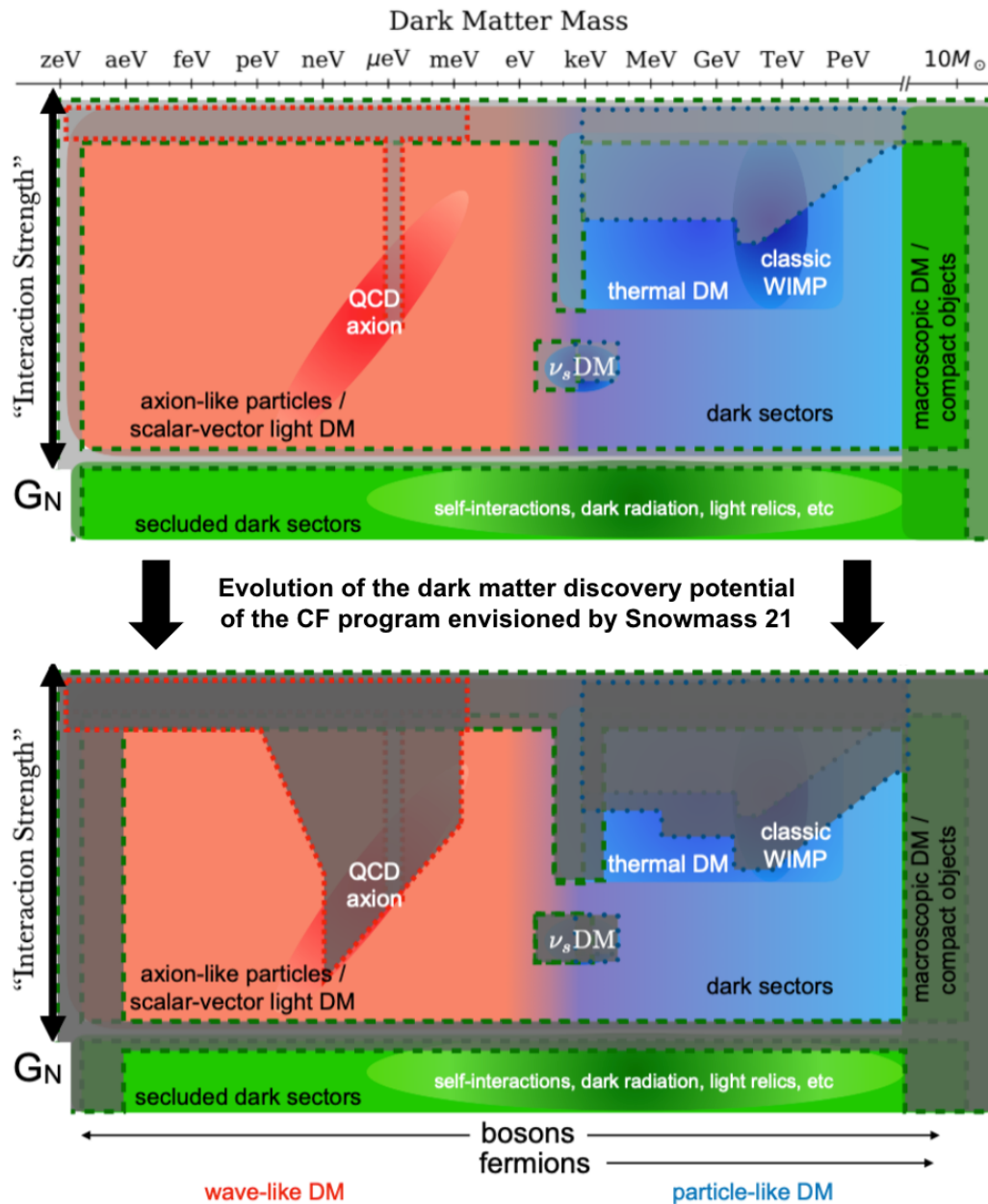
321 Accomplishing the full scientific scope of the cosmic frontier program will require that robust support for  
322 ongoing, new, and future projects be well-matched by funding mechanisms for research groups to do science  
323 with data from those projects. This includes research support for teams within experimental collaborations  
324 and beyond (e.g., theory and simulation developments as well as detector R&D studies that may not be part  
325 of an existing experiment in our portfolio today, but that will enable the next generation of experiments).  
326 In particular, a strong feature of the CF program is that multi-project analyses can yield results of greater  
327 impact than the sum of their parts. Thus, new mechanisms in support of such multi-project initiatives have  
328 been identified as an important component of the CF strategy for discovery.

## 329 5.3 Dark Energy

330 The accelerating expansion of the modern universe over the last 5 billion years of cosmic history represents  
331 one of the most formidable scientific problems of our time. Its discovery was awarded the physics 2011  
332 Nobel prize. Similar to what has happened after the discovery of the Higgs boson by our colleagues in the  
333 energy frontier, this discovery has launched our field into a new era of precision measurement studies aiming  
334 to understand the physics of cosmic acceleration — the explanation for which must clearly lie beyond the  
335 Standard Model.

336 Developing an understanding of the cause of cosmic acceleration in the modern universe should revolutionize  
337 our understanding of fundamental physics. Either we will discover a new form of energy whose source lies  
338 beyond the Standard Model of particle physics – some form of dark energy; or we will motivate a new  
339 theory of gravity beyond Einstein’s general relativity. If dark energy is the solution, it then accounts for  
340 two-thirds of the total present mass-energy density of the universe and permeates the empty space between  
341 galaxies. It hinders the growth of density perturbations such as those which can form clusters of galaxies, and  
342 makes distant astrophysical sources appear dimmer than expected by increasing their distances. Detailed  
343 measurements of both the expansion history of the universe and of the rate at which cosmic structures grow  
344 are central to understanding the nature of cosmic acceleration; in combination, they allow us to distinguish  
345 explanations associated with dark energy from those which require modifications to theories of gravity.

346 Cosmology research in the twenty-first century is dominated by large cosmic survey experiments carried out  
347 by world-wide collaborations comprising hundreds of members from dozens of institutions. These large teams  
348 have developed specialized equipment and novel analysis tools which enable us to carry out observations of  
349 unprecedented scope and precision. For example, the Dark Energy Survey (DES) collaboration uses a 570-  
350 megapixel camera, the Dark Energy Camera (DECam), installed on a 4-meter telescope in Chile; the Dark  
351 Energy Spectroscopic Instrument (DESI) collaboration uses a system of 5000 robotic fiber positioners on a  
352 4-meter telescope in Kitt Peak, Arizona; and the Legacy Survey of Space and Time (LSST) Dark Energy  
353 Science Collaboration (DESC) will use a 3.2 gigapixel camera on a 8.5-meter telescope at the Vera Rubin  
354 Observatory in Chile.



**Figure 5-2.** Cartoon (not to scale) of the dark matter sensitivity of the Delve Deep, Search Wide strategy outlined in this report compared to current dark matter limits. The range of possible dark matter masses with well-studied origin stories now spans many orders of magnitude while the interaction strengths with standard model particles can range from electroweak scale down to purely gravitational strength. Broad, logarithmic coverage of this expanded dark matter parameter space is attainable with a range of new techniques which have been developed, but will require an ambitious and significantly expanded program of mid- and small-scale experiments as well as novel cosmic probes using existing and planned survey instruments. For details on each of the experimental projects and initiatives, please refer to the dark matter section of the report.

355 As a result of this robust series of cosmic survey experiments, the community has advanced in leaps and  
 356 bounds towards our goal of understanding cosmic acceleration via precision measurements of key cosmological  
 357 parameters. These key parameters include:

- 358 • **The equation of state parameter of dark energy ( $w$ ).** — If  $w$  is exactly equal to  $-1$  and does  
 359 not change throughout the cosmic history, then dark energy acts as a cosmological constant – the  
 360 simplest, yet non-trivial, dark energy model. One of the key questions that our community wants to  
 361 shed light upon is whether or not dark energy is consistent with this model. Determining the dark  
 362 energy equation of state with percent-level precision over a wide redshift range is key for addressing  
 363 many fundamental questions related to dark energy and cosmic acceleration. Over the past twenty  
 364 years of CF research cosmic surveys have made huge progress in this area. We now can measure  $w$   
 365 well at  $z < 0.7$ ; observations so far are consistent with a cosmological constant model, but the time  
 366 evolution of the cosmic equation of state is only poorly constrained so far. Our next leap will be to  
 367 study the signatures of dark energy to the limits of the modern universe ( $z \lesssim 5$ ), which will provide  
 368 novel insights into the earliest times when it began to become significant.
  
- 369 • **The rate of expansion of the universe today ( $H_0$ ).** — This parameter is related to dark energy  
 370 because it provides the low-redshift anchor to measurements of the cosmic expansion history. As a  
 371 result, the value of  $H_0$  has an impact on many constraints on  $w$ .  $H_0$  has been determined with good  
 372 precision by experiments relying on data from two vastly different eras of the universe. Early-universe  
 373 measurements (using for example CMB observations) would be expected to agree with values derived  
 374 from galaxies, supernovae and other observables in the dark energy-dominated modern universe. A  
 375 discrepancy between the two regimes would indicate either that the model is incomplete or that there  
 376 are systematic issues in the distance determinations used to anchor sets of  $H_0$  measurements. At this  
 377 point, measurements of  $H_0$  with precision of 1-2% have been made with a variety of techniques; the  
 378 tension between determinations tied to the low-redshift versus high-redshift distance scales is starting  
 379 to exceed the 5-sigma level. Several new measurements that will help uncover the origins of this  
 380 discrepancy over the next two decades are described in this report.
  
- 381 • **The amplitude of fluctuations in the density of matter in the universe ( $S_8$ ).** — Precision  
 382 measurements of the amplitude of density fluctuations in the universe and their growth over time are  
 383 highly complementary to measurements of the cosmic expansion history (which should directly depend  
 384 on the dark energy equation of state). For instance, if general relativity provides the correct description  
 385 of gravity on large scales, the growth rate of density perturbations can be predicted directly from the  
 386 expansion history; however, if cosmic acceleration occurs due to departures from Einsteinian gravity,  
 387 this relationship is broken. Early and late-universe measurements of the amplitude of the matter power  
 388 spectrum, as measured by the  $S_8$  parameter, are currently in tension. The significance of this tension  
 389 is lower than for  $H_0$ , but the ongoing CF experimental program is rapidly improving the precision of  
 390 measurements in this area, and percent-level determinations of  $S_8$  in both the modern and CMB eras  
 391 will soon be possible.

### 392 5.3.1 The Next Stages of Dark Energy Experiments

393 The dark energy component of the CF program has long been planned to proceed via a series of cosmic  
 394 survey *Stages*, each greater in scope and precision measurement capability than its predecessor, as defined  
 395 in the Dark Energy Task Force (DETF) report [8]. The current state-of-the-art results are coming from  
 396 DETF Stage III surveys (e.g., BOSS, eBOSS and DES); DETF Stage IV experiments are now underway

397 (e.g., DESI) or soon to begin (LSST). These new experiments should provide percent-level constraints on  $w$   
 398 (for models where its value does not evolve with time), driven primarily by measurements at  $z < 1.5$ .

399 Since the last Snowmass report, CMB experiments have become a significant part of the CF program. The  
 400 next-generation project, CMB-S4, was ranked highly by the previous Snowmass/P5 process and is currently  
 401 at the CD-0 stage. The data obtained by CMB experiments plays a vital role in exploring the fundamental  
 402 HEP questions of cosmic acceleration in both the modern universe (dark energy) and the early universe  
 403 (inflation).

404 Looking forward, within the next two decades, our community seeks to build CMB-S4 as well as to develop  
 405 new Stage V and Stage VI projects. This suite of experiments will make precision measurements of  $w$ ,  $H_0$ ,  
 406 and  $S_8$  and other cosmological parameters with data sets that probe not only the current era in which dark  
 407 energy is firmly dominant, but also detailed measurements going well beyond the transition between the  
 408 dark matter- and dark energy- dominated eras and new tests of consistency with early universe data.

409 In this report, we categorize the next stages of dark energy and cosmic acceleration experiments as follows:

- 410 • Stage IV – In this category we include those experiments which meet the DETF Stage IV definition  
 411 (such as DESI and Rubin LSST), as well as fourth-generation CMB experiments. Both classes of  
 412 experiment complement each other, as the science of early and late-time cosmic acceleration are strongly  
 413 intertwined. A Stage IV CMB project (CMB-S4) is currently under development and should play a key  
 414 role in cosmic frontier science. Stage IV cosmic survey experiments will provide precision constraints  
 415 on dark energy, predominantly over the redshift range  $z \sim 1.5$  (vs.  $z \sim 0.7$  for Stage III). Novel  
 416 analysis techniques that can take advantage of the leap in constraining power provided by Stage IV  
 417 experiments, both when they are considered on their own and especially when analyzed in combination,  
 418 should reduce systematic uncertainties and establish the true significance of the emerging tensions in  
 419  $H_0$  and  $S_8$ .
- 420 • Stage V – In this category we include those experiments that would represent a substantial improvement  
 421 ( $3 - 4\times$  or better) in capabilities or cosmological constraining power over the equivalent Stage IV  
 422 experiment. These projects will take the current precision cosmology program to the next level,  
 423 providing percent-level constraints on the contribution of dark energy to the total density of the  
 424 universe at redshifts up to  $z \sim 5$ ; enabling new probes of the physics of inflation, with the potential  
 425 to rule out broad swaths of models and to measure inflationary energy scales; and exploring in depth  
 426 possible explanations for the origins of the current  $H_0$  and  $S_8$  tensions.
- 427 • Stage VI – This category encompasses a future generation of dark energy experiments which can  
 428 further increase constraining power over Stage V and begin to test cosmology in the ‘dark ages’ of  
 429 the Universe at  $5 \lesssim z \lesssim 10$ . Such projects will require smaller pathfinder efforts now to develop key  
 430 technologies and methods, which should enable them to become ready to reach project status  $\sim 10$   
 431 years from now. These ambitious experiments will incorporate novel technologies and take advantage  
 432 of new opportunities, driving the field to its next leap in dark energy precision studies. Higher-redshift  
 433 experiments should be able to explore even larger scales than Stage V, expanding the sensitivity to  
 434 inflationary signatures, as well as providing sensitive tests for the presence of early dark energy.

### 435 5.3.2 A Stage V Spectroscopic Facility (Spec-S5)

436 The greatest near-term opportunity to revolutionize our understanding of cosmic acceleration both in the  
 437 modern universe and the inflationary epoch will be provided by a new Stage V Spectroscopic Facility (Spec-



438 S5) which would combine a large telescope aperture, wide field of view, and high multiplexing. The technology  
 439 for such a facility is at this point well-developed, building substantially on the legacy of DESI and Rubin  
 440 Observatory. A Spec-S5 can simultaneously provide a dense sample of galaxies at lower redshifts to provide  
 441 robust measurements of the growth of structure at small scales, as well as a sample at redshifts  $2 < z < 5$   
 442 to measure cosmic structure at the largest scales, spanning a sufficient volume to probe primordial non-  
 443 Gaussianity from inflation and to search for features in the inflationary power spectrum on a broad range  
 444 of scales, while also testing dark energy models in poorly-explored regimes, determining the total neutrino  
 445 mass, and strongly constraining the effective number of light relics. A Spec-S5 would also be able to probe  
 446 the nature of dark matter using the kinematics of stars in the Milky Way halo and measurements of the  
 447 matter power spectrum at small scales.

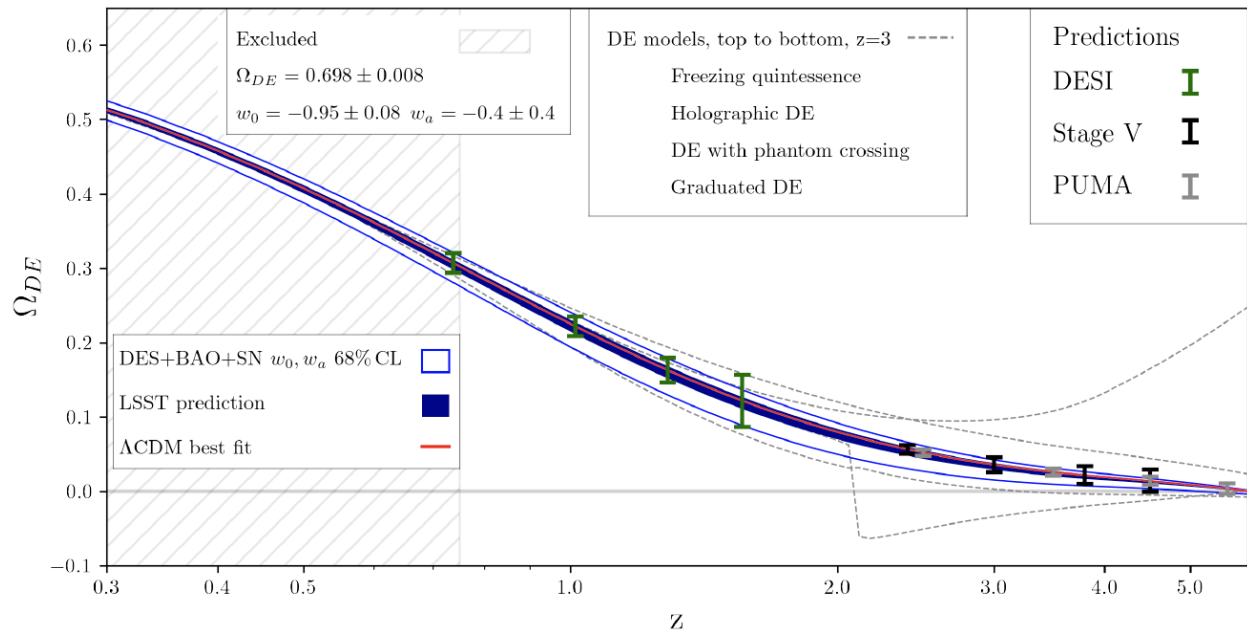
448 Multiple options for facilities that would qualify as a Spec-S5 have been proposed. These projects are  
 449 described in the CF4 report [4], which also outlines the key factors to be considered when evaluating candidate  
 450 facilities. The key requirement is that a Spec-S5 should represent a significant advance over what is possible  
 451 with Stage IV facilities, should enable transformational progress on all of the Stage V science goals, and  
 452 open up new discovery space, including:

- 453 • Measuring the contribution of dark energy to the total mass-energy density in the universe with  
 454 uncertainties below 2% of  $\Omega_{\text{total}}$  out to  $z = 5$ , redshifts at which dark energy becomes dynamically  
 455 negligible in  $w_0/w_a$  models, while simultaneously constraining modified gravity models over this range;  
 456 cf. Figure 5-3. Spec-S5 measurements will also provide sensitive tests for early dark energy (at  
 457  $500 < z < 50000$ ), as illustrated in Figure 5-4.
- 458 • Exploring the physical nature of the tensions in current measurements of  $H_0$  (e.g., via BAO-based mea-  
 459 surements spanning an unprecedented range of redshifts) and  $S_8$  (with dense low- $z$  samples providing  
 460 many cross-checks for systematics by enabling measurements using many separate galaxy populations  
 461 as well as multiple lensing methods).
- 462 • Testing for signatures of non-Gaussianity from inflation at a sufficient precision to be able to rule  
 463 out all non fine-tuned multi-field models, while simultaneously searching for inflationary primordial  
 464 features that signal the breaking of scale invariance and constraining contributions from dark radiation  
 465 (sub-eV particles thermalized in the Universe). The Spec-S5 studies of inflation and dark radiation are  
 466 points of synergy with CMB-S4 and will help the community to build a unified picture of the physics  
 467 underlying both the early and late-time eras of cosmic acceleration.

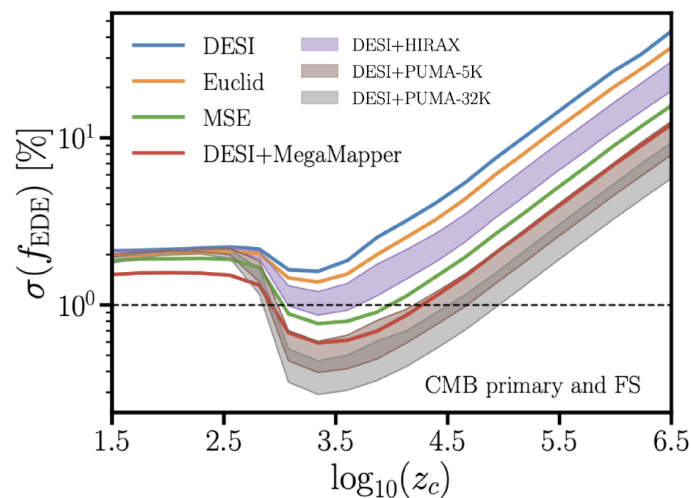
468 A major strength of Spec-S5, in any of its proposed implementations, will be its ability to advance our  
 469 understanding on multiple science fronts simultaneously, while also enhancing the science outcomes from  
 470 other experiments in the CF portfolio. Cross-correlation measurements that combine Spec-S5 data with other  
 471 surveys can be used to improve photometric redshifts for LSST as well as to unlock additional cosmological  
 472 information (generally with reduced systematic uncertainties). Construction of the instrument for a Spec-S5  
 473 facility in collaboration with other partners could build upon HEP strengths and experience with DESI.  
 474 Realizing Spec-S5 in any of the funding scenarios that P5 might consider would play a key role in advancing  
 475 the scientific goals of the CF community.

### 476 5.3.3 New Opportunity: Gravitational Waves

477 The emergence of gravitational wave observatories (GWO) sensitive enough to detect sources at cosmic  
 478 distances has revolutionized humanity's view of the universe. The HEP community, and the Cosmic Frontier



**Figure 5-3.** Projected sensitivity to the fractional dark energy density relative to closure density as a function of redshift. Solid lines indicate extrapolations to high redshift of current conventional models fitted to data at low redshift  $0.1 < z < 0.75$ . Dashed lines indicate predictions of new models designed to address various tensions in the datasets. Data points with error bars indicate how the 2% level sensitivity obtained in ongoing and future high redshift surveys can directly constrain and discriminate between these models.



**Figure 5-4.** Projected sensitivity of current and next-generation high redshift cosmic surveys to the fractional dark energy density relative to closure density near the CMB recombination epoch. The signal of early cosmic acceleration is an alteration to the growth of structure as measured via features in the matter power spectrum of the large scale structure. In particular, these probes will test models of early dark energy designed to explain the “Hubble tension,” a discrepancy between high redshift CMB measurements and low redshift optical survey measurements of the Hubble expansion rate which is growing in statistical significance.

479 sub-community in particular has been quick in realizing the implications of this new development and has  
 480 been working towards incorporating GWO projects in our portfolio. Gravitational-wave standard sirens  
 481 (merging compact object binary systems) allow measurement of the luminosity distance of the source  
 482 and, together with redshift measurements, can be used to measure  $H_0$  via the distance-redshift relation.  
 483 Measurement of the Hubble parameter using standard sirens does not require a cosmic distance ladder  
 484 and is model-independent: the absolute luminosity distance is directly calibrated by the theory of general  
 485 relativity. Aiming to develop standard sirens into fully-fledged probes of dark energy, the DES collaboration  
 486 launched a search and discovery program for the electromagnetic (EM) counterparts of events detected by  
 487 the current GWO projects LIGO/Virgo. DES participated in the first multi-messenger discovery, of the  
 488 binary neutron star merger GW170817, and contributed to the first standard siren-based measurements  
 489 of  $H_0$ . Approximately fifty additional multi-messenger binary neutron star observations would be needed  
 490 to reach the required precision to weigh in the Hubble tension debate [9]. The community is currently  
 491 planning upgrades to the existing GWO network as well as a next-generation GWO network. One proposed  
 492 next-generation project, currently being led by the US community, is known as Cosmic Explorer. The  
 493 CF community plans to use standard sirens from a Cosmic Explorer-like GWO as a powerful sample of  
 494 independent distance indicators going all the way up to  $z = 10$ . Combined with deep optical-to-near-  
 495 infrared Stage IV and Stage V EM observatories, we can reach the required precision for Stage VI dark  
 496 energy science.

497 Gravitational waves can probe dark energy and cosmic acceleration throughout the entire history of the  
 498 universe with an observable that is novel and largely independent from the traditional observables employed  
 499 thus far in the field. The next-generation GWO network will also have access to the binary black hole  
 500 population when the universe was still in its infancy, to the equation of state of matter at neutron star  
 501 cores at supranuclear densities and quark deconfinement phase transitions in hot merger remnants, and the  
 502 ability to measure the properties of dark energy and dark matter, to stochastic gravitational-waves from  
 503 early-universe phase transitions, and to the highly warped space-time in the strong-field and high-velocity  
 504 limit of general relativity.

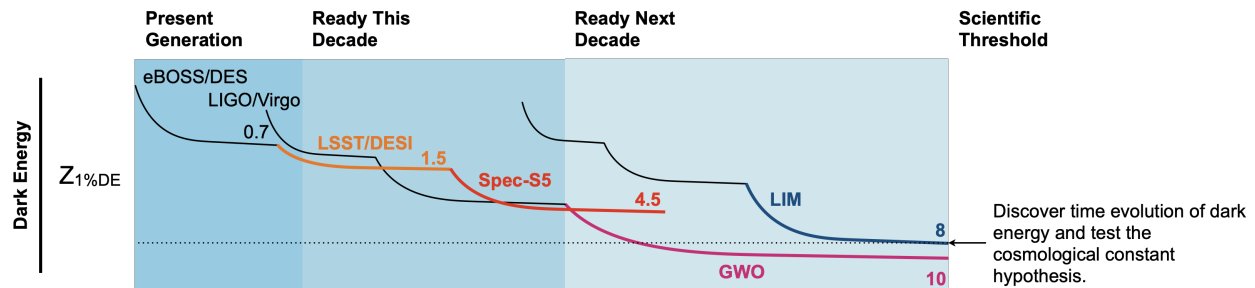
505 The CF community plans to incorporate GWO in its portfolio of tools for discovery with a long term strategic  
 506 vision. We will pursue EM counterparts of events detected by the growing GWO network while launching  
 507 new pathfinder (R&D) efforts to enable the HEP community to participate in the next-generation GWO  
 508 project in a leading role. The new detector's sensitivity, roughly 10 times better than the planned LIGO  
 509 upgrade, requires significantly larger facilities and a number of technological upgrades. Both are challenging  
 510 requirements that the HEP community is well-equipped to meet, given our experience.

511 This is likely a once-in-a-century opportunity for the HEP community to make new breakthroughs in an  
 512 entirely new class of experiments and utilize this new opportunity to advance on our scientific drivers at a  
 513 much faster pace than previously anticipated.

### 514 5.3.4 Experimental Program & Facilities: Cross-Experiment Strategic Plan

515 In light of the opportunities for new breakthroughs in dark energy and cosmic acceleration over the next  
 516 decades, the CF community has converged on a strategic plan with both near-term and longer-term action  
 517 items, built upon a set of highly complementary projects/facilities and research initiatives. Figure 5-5  
 518 summarizes how cosmic surveys will advance over Stages IV-VI using a key statistic: the maximum redshift  
 519 where we can measure the contribution of dark energy to the total mass-energy density of the universe with  
 520 percent-level precision. Stage III has achieved this at  $z < 0.7$ , and Stage IV should reach this goal for  $z \lesssim 1.5$ .  
 521 In contrast Stage V and VI experiments should extend our reach to  $z \sim 5$  and  $z \lesssim 8-10$ , respectively. These

522 thresholds will allow us to definitively test whether simple cosmological constant or time-varying  $w_0/w_a$   
 523 models are sufficient descriptions of cosmic acceleration, or if instead more exotic explanations with more  
 524 complicated time evolution (such as those invoked to explain existing tensions, such as the models illustrated  
 525 in Figure 5-3) will be needed. In either scenario, Stages V and VI will lead to a breakthrough comparable  
 526 to, if not surpassing, the discovery of dark energy itself.



**Figure 5-5.** Timeline illustrating the expected improvement in the redshift range over which the contribution of dark energy to the total mass-energy of the universe is strongly constrained. Their performance is quantified by  $z_{1\%DE}$ , the highest redshift at which the fraction of the critical density in dark energy,  $\Omega_{DE}$ , can be measured with percent-level precision (uncertainties of 0.01 or less). We set the threshold for discovery at  $z = 8$ , which will allow us to rigorously test the cosmological constant hypothesis far beyond the dark energy-dominated era of the universe. Bold lines show medium and large projects ready to build this decade or the next; thin lines represent the current results and/or pathfinders for the next generation of projects.

527 Figures 5-3 and 5-4 show projections for the sensitivity of example proposed Stage V and Stage VI surveys  
 528 relative to current Stage IV surveys for two key parameters: the energy density of dark energy ( $\Omega_{DE}$ ), and  
 529 the contribution of early dark energy to the overall cosmological model ( $f_{EDE}$ ). The program devised by  
 530 the CF community will enable us to test a wide variety of dark energy models that have been proposed to  
 531 explain the observed tensions in  $H_0$  and  $S_8$ .

532 We summarize here the expected progression of projects and the key areas of support needed to address  
 533 the questions of cosmic acceleration in both the modern and inflationary eras as described in the CF4 [4],  
 534 CF5 [5], and CF6 [6] reports (noting that there is no significance to the ordering of items):

535 **Near-future efforts:** In the near term, it will be important to complete the Stage IV projects and fund the  
 536 science efforts needed to take advantage of these powerful new facilities, including obtaining complementary  
 537 data which will make them more powerful. Key needs to ensure the success of this program are to:

- 538 • Build and operate CMB-S4.
- 539 • Begin operation of Rubin LSST.
- 540 • Continue operation of DESI (via a new DESI-II program) to constrain dark energy in new domains  
 541 and as a step towards a Stage V spectroscopic facility (Spec-S5).
- 542 • Fund science analyses of the Stage IV datasets.
- 543 • Establish funding mechanisms for small and medium scale projects to follow up transients discovered  
 544 by LSST and standard sirens discovered by gravitational wave observatory (GWO) facilities, as well  
 545 as to enable photometric redshift training spectroscopy for LSST.

546 **Longer-term efforts:** In order to enable longer-term progress, it will be vital to establish a flagship Spec-  
 547 S5 program, as well as to explore new opportunities which can enable other Stage V and even Stage VI  
 548 experiments. Key needs for this progress are to:

- 549 • Build and operate Spec-S5.
- 550 • Support R&D and pathfinder studies for a next-generation GWO.
- 551 • Support R&D and pathfinder studies for a next-generation CMB experiment (at the Stage V or VI  
 552 level).
- 553 • Support R&D and small projects to develop technologies and methods that can enable a future surveys  
 554 (e.g., LIM) and new precision probes of cosmology (e.g., redshift drift measurements).
- 555 • Support operations of Rubin Observatory beyond the initially planned 10 year LSST survey; the best  
 556 future use of this facility should be evaluated later this decade after the first LSST analyses have been  
 557 done.

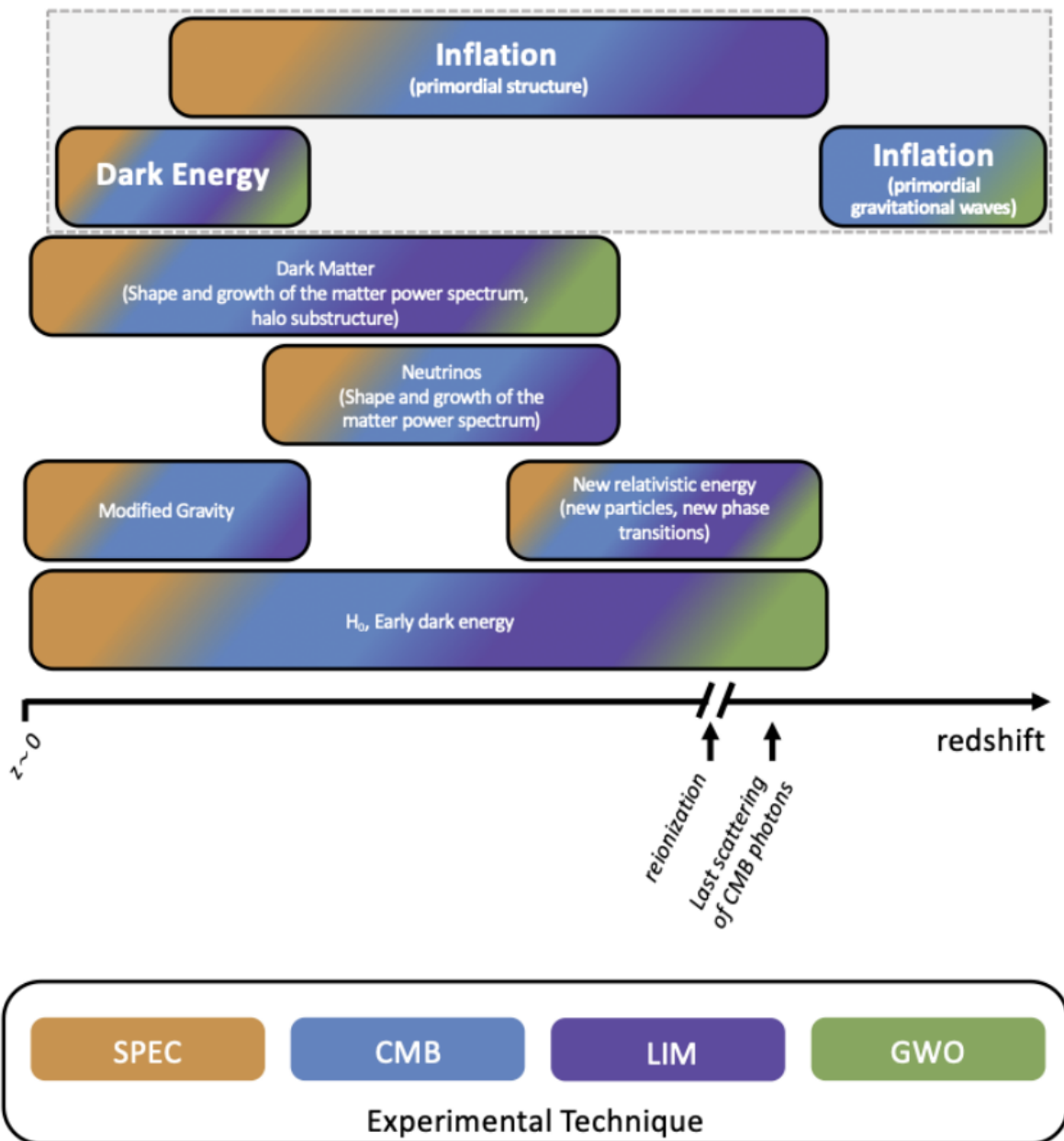
558 **Taking advantage of complementary experiments:** The experiments in our program will probe dark  
 559 energy physics in a variety of different ways, enabling cross-checks for and control of systematic uncertainties  
 560 to obtain robust and rigorous results. For example, standard siren measurements based on sources from  
 561 gravitational wave observatories will be subject to very different systematics from observables such as  
 562 galaxy clustering or supernovae that have been commonly employed in precision cosmology studies to date.  
 563 Furthermore, different experiments provide complementary information about the universe that yields more  
 564 powerful constraints on cosmology when analyzed in combination.

565 However, such combined analyses present more challenges (particularly organizationally) than those which  
 566 only involve one science collaboration. Key needs to ensure the success of multi-experiment analyses are to:

- 567 • Create funding streams and support for cross-survey analyses.
- 568 • Develop and support coordination between large facilities for optimized design, timely execution, and  
 569 joint analyses.
- 570 • Create data archive centers for long term storage where data from cosmological surveys is replicated,  
 571 continuously available, and easily accessible with software and computing resources for developing joint  
 572 constraints on dark energy, inflation and dark matter.
- 573 • Create and support development of a diverse set of simulated data sets that could be used in joint  
 574 analyses. Support should include common access to data and super-computing resources.

575 The CF strategy for medium and large cosmic survey facilities is designed to take advantage of the in-  
 576 terconnected nature of the science questions we aim to approach. The same facilities which can provide  
 577 compelling tests of the nature of cosmic acceleration can also probe the physics of inflation, neutrino masses,  
 578 dark matter, and dark radiation, as described in other sections of this report. This is illustrated in Figure  
 579 5-6, taken from the CF6 report [6], which depicts the different areas of cosmic frontier science addressed  
 580 by wide-field spectroscopic facilities (such as Spec-S5), CMB experiments (such as CMB-S4), line intensity  
 581 mapping (which may play a key role in Stage VI), and gravitational wave observatories.

582 By incorporating both near-term and long-term projects at a range of scales, this experimental program  
 583 should ensure that high-impact science results are continuously produced, avoiding lengthy gaps during  
 584 project construction. In sum, the ecosystem of cosmic surveys proposed by the community to address the



**Figure 5-6.** A high-level summary of the key scientific opportunities. The horizontal extent of each box corresponds to the redshift-range of the tracer, while the coloring indicates the experimental technique used to measure the signal. The dashed grey box emphasizes dark energy and inflationary probes. Figure taken from the CF6 report [6].

585 fundamental question of the source of cosmic acceleration is carefully thought out to efficiently address some  
 586 of the greatest scientific questions of our time, extending across all areas of cosmology. Accomplishing this

587 program in its entirety will ensure that the HEP community will continue its path of leadership in studies  
588 of the accelerating universe.

## 589 5.4 Inflation

590 There is strong evidence that the Universe went through a hyper-accelerated period of expansion during its  
591 earliest moments. This period of cosmic inflation reconciles a number of otherwise puzzling observations via  
592 the simple mechanism of an epoch in which cosmic evolution is dominated by a large quantity of vacuum  
593 energy. The consequences of inflation offer important clues to the necessary modification of the Standard  
594 Model of particle physics, in order to describe the dynamics at work. In addition, during this time the  
595 Universe effectively conducted the most extreme high energy physics experiment imaginable.

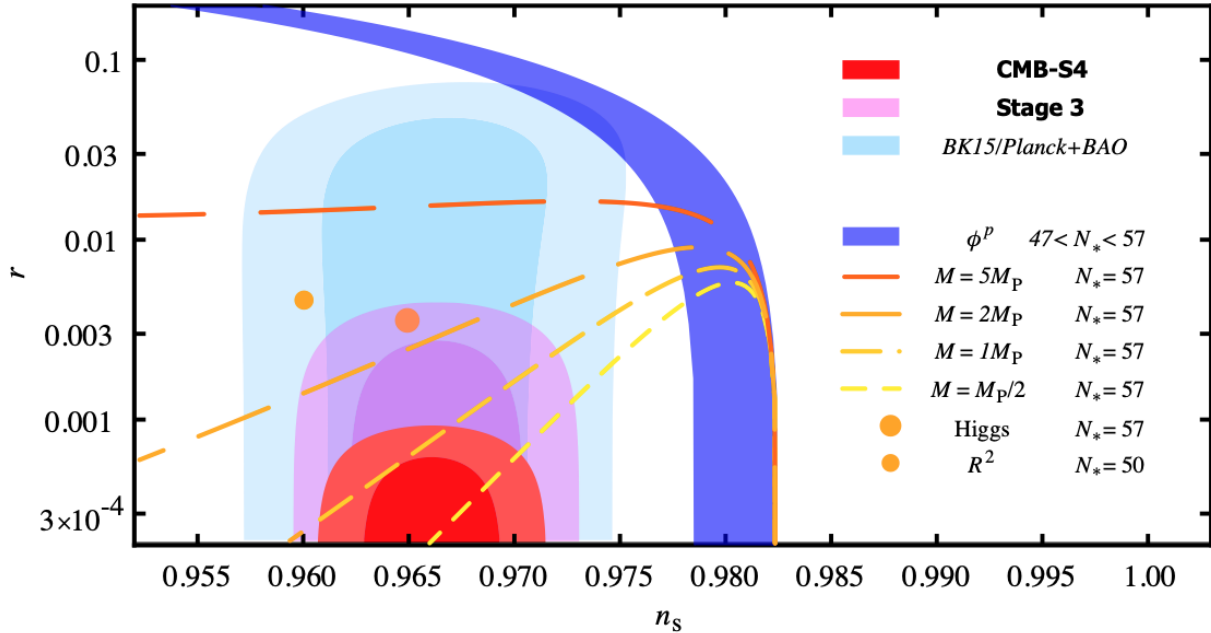
### 596 5.4.1 Determining the Energy Scale of Inflation

597 Future observations are poised to reveal the clues necessary to select among models of inflation. The  
598 amplitude of the associated gravitational waves (generally characterized by the tensor to scalar ratio,  $r$ ),  
599 which can be inferred from the primordial  $B$ -modes of the CMB, encodes the absolute energy scale at which  
600 inflation takes place; projected sensitivity is shown in figure 5-7. CMB-S4 aims to improve sensitivity by  
601 more than an order of magnitude, covering much of the parameter space of natural models of inflation.  
602 The shape of the primordial power spectrum and its deviation from Gaussian statistics encodes inflaton  
603 interactions; a large-scale structure survey from a Stage-V spectroscopic survey and pathfinding missions  
604 based on radio and millimeter intensity mapping instruments would enable determinations of  $f_{NL}$  and  $A_{lin}$   
605 and explore the range of dynamics. Inflationary physics may further imprint itself on the spectrum of the  
606 relic background of gravitational waves,  $\Omega_{GW}$ . A cartoon summary of the potential for these advances is  
607 shown in Figure 5-8.

### 608 5.4.2 Reconstructing the Physics of Inflation

609 A large part of the future cosmic surveys program focuses on discovering new dynamics at ultra high energy  
610 scales using the inflaton beam as a cosmic collider. During the epoch of inflation, the cosmic particle horizon  
611 recedes with acceleration equal to the Hubble parameter  $H_I = \sqrt{\rho_I}/M_{\text{Planck}} \leq 10^{14}$  GeV. Interactions of  
612 quantum fluctuations with the accelerating horizon create a white spectrum of Hawking radiation in inflaton  
613 waves, gravitational waves, and possibly also waves of low mass dark matter particles such as axions or dark  
614 photons. While the gravitational waves propagate without rescattering to the present day and create  $B$ -  
615 mode polarization patterns in the CMB sky, and searches for ultralight dark matter are discussed elsewhere  
616 in this report, this section will focus on the hot fireball of inflaton waves which is traditionally measured in  
617 the temperature power spectrum of the CMB. These inflatons can scatter on each other via new dynamical  
618 interactions that could arise within the inflaton sector itself through its nonlinear potential, or could more  
619 generally be present in other high scale sectors that are accessible to this ultra high energy beam. Since, by  
620 the process of reheating, the inflaton waves eventually decay to produce the matter and radiation content of  
621 the universe, the spectrum of this primordial blackbody radiation is imprinted upon the large scale structure  
622 of the distribution of matter in the universe.





**Figure 5-7.** The tensor-to-scalar ratio  $r$  plotted vs. the matter power spectral index  $n_s$ .  $r$  is observed from B-mode polarization patterns in the CMB sky, and is a measure of the absolute energy scale of cosmic inflation in the  $10^{16}$  GeV range. The light blue region represents current constraints, the purple region the projected constraints from ongoing experiments including the South Pole Telescope and the Simons Observatory, and the red region the targeted sensitivity of CMB-S4. Dots and lines represent various models of inflation which will be tested.

623 While new particle resonances in propagators are naturally spectroscopically measured via the 2-point  
 624 correlation function, dynamical interaction terms in the Hamiltonian can be elucidated via 3-wave or 4-  
 625 wave mixing measurements. In particular, the 3-point correlation function of the matter distribution can  
 626 be measured in the fluctuations of the CMB sky, in redshift-resolved distributions of galaxies and clusters  
 627 obtained via optical spectroscopic surveys, or even in redshift-resolved distributions of neutral hydrogen gas  
 628 (or other species) obtained via the new technique of line intensity mapping. Just as the event topology and  
 629 angular distribution of particles and jets in a collider experiment provide information about the underlying  
 630 scattering event, the shape (local, equilateral, folded) of the 3-point triangle that is chosen for the analysis  
 631 can be used to distinguish between different types of interaction terms in the ultrahigh energy Lagrangian.  
 632 The strength of the interaction can be characterized by taking the Fourier transform of the 3-point function  
 633 to form the power bispectrum and measuring its spectral amplitude  $f_{\text{NL}}$ . As inflaton wave energies in  
 634 the originally Gaussian thermal spectrum are redistributed by repeated scattering events, primordial non-  
 635 Gaussianities manifest as non-vanishing amplitudes for these higher order correlation functions.

636 Looking at the night sky, triangles formed at each distance scale probe scattering events at a different center-  
 637 of-mass energy. The collection of such events in the sky forms a high statistics data set of independent  
 638 scattering events at each energy. Higher angular or redshift resolution surveys will naturally provide higher  
 639 statistics as typically characterized by the total number of resolved wave modes  $N_{\text{modes}}$ .

640 While CMB experiments including CMB-S4 will provide pristine measurements of the dynamical interaction  
 641 amplitudes using data at a single slice in time or redshift, current optical spectroscopic surveys such as DESI  
 642 also provide high resolution redshift information to measure the matter distribution in all 3 dimensions.  
 643 Future instruments and facilities aim to increase the scattering statistics by a factor 10-100 relative to the



644  $N_{\text{modes}}$  achieved by DESI, an improvement factor measured by the "primordial figure-of-merit." For example,  
 645 the first factor of 10 would come from a Stage V Spectroscopic Facility targeting  $2 < z < 5$  galaxies using  
 646 the LSST dataset in order to reconstruct the full 3-d matter distribution. The statistics could be improved  
 647 by a factor of 100 using a relatively inexpensive line intensity mapping (LIM) technique to probe the same  
 648 physics and by using 21 cm observations to map the large scale distributions of neutral hydrogen in the  
 649 universe via 21 cm emissions. An ambitious program, exemplified by the LuSEE-Night pathfinder, seeks to  
 650 develop the LIM technique and understand the systematics to determine whether the full statistical power  
 651 could in fact be attained.

652 A cartoon summary of the potential for these advances is shown in Figure 5-8.

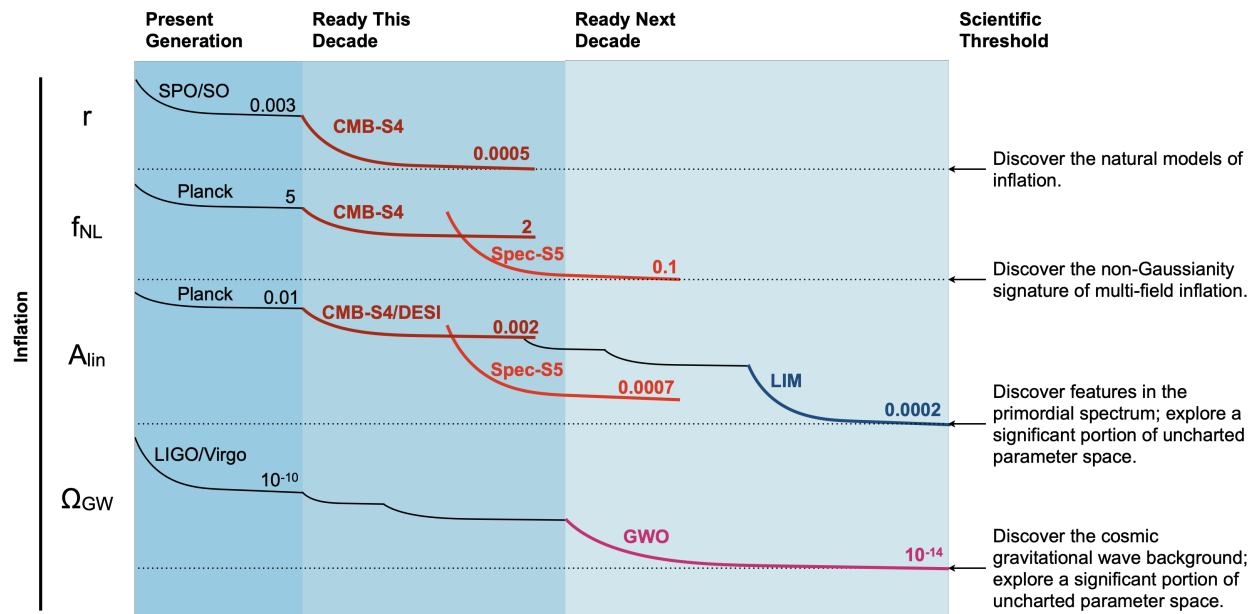


Figure 5-8. Timeline of medium and large projects for inflation science.

### 653 5.4.3 CMB-S4

654 CMB-S4 [10] is a Stage IV cosmic microwave background project that plans to field multiple telescopes at  
 655 the South Pole and in the Atacama desert, Chile. See the Snowmass 2021 White Paper [11] for a discussion  
 656 on the broader experimental context. CMB-S4 has an enormously broad science case with key science goals  
 657 including searching for primordial gravitational waves through the B-mode signal in the CMB as predicted  
 658 from inflation (detecting  $r > 3 \times 10^{-3}$  at  $5\sigma$  or limiting  $r \leq 10^{-3}$  at 95% confidence if  $r$  is very small)  
 659 and searching for the imprint of relic particles including neutrinos (measuring  $\sigma_{N_{\text{eff}}}$  with uncertainty  $\leq 0.06$   
 660 at 95% confidence). CMB-S4 will also offer unique insights into dark energy and tests of gravity on large  
 661 scales, find large samples of high-redshift galaxy clusters, elucidate the role of baryonic feedback on galaxy  
 662 formation and evolution, open a window onto the transient Universe at millimeter wavelengths, and explore  
 663 objects in the outer Solar System, among other investigations.

664 The current CMB-S4 instrumental design calls for 500,000 polarization-sensitive bolometers that measure  
 665 the sky at frequencies from 20–280 GHz. The superconducting detectors will be read out using time-domain

666 multiplexed electronics and will be distributed between a set of telescopes at two sites: two 6-meter cross-  
667 Dragone reflecting telescopes in Chile, a 5-meter three-mirror-astigmatic reflecting telescope at the South  
668 Pole, and eighteen 0.5-meter refracting telescopes at the South Pole, grouped as triplets on six mounts,  
669 with each mount sharing a cryogenic system. This preliminary baseline design for CMB-S4 uses proven  
670 technology scaled up to much higher detector counts. The already-established integrated project office has  
671 addressed the technical challenges of the required scale-up with a detailed design and implementation plan  
672 that includes a full work-breakdown structure. The project office also maintains a register of risks and a  
673 detailed cost estimation plan. The preliminary, technology-limited project schedule contains nearly nine  
674 thousand milestones and catalogs the dependent relationships between them. A dedicated group within  
675 the project oversees the production of prototype detectors and leads the development of a unified detector  
676 fabrication plan that covers multiple fabrication sites.

## 677 5.5 Dark Matter

678 Astronomical and cosmological observations of the gravitational influence of dark matter provide one of  
679 the strongest indications of new physics beyond the Standard Model [12]. Ascertaining the nature of these  
680 mysterious new particles, their interactions both with Standard Model particles and with themselves, and  
681 their cosmological origin is perhaps the grand challenge of this generation. While the ambitious program  
682 to understand dark matter will require tools and techniques from across the HEP frontiers [13], the Cosmic  
683 Frontier is unique in that its experiments seek to detect and measure dark matter in its natural habitat –  
684 the halo of our Galaxy, the halos of distant galaxies, and the large-scale structure of the Universe. While  
685 other frontiers may hope to discover new particles that could play the role of dark matter, only the Cosmic  
686 Frontier can establish that a given discovery is, in fact, associated with the dark matter in the Universe.

687 Our understanding of the landscape of dark matter theories has evolved significantly in the past several years,  
688 as theoretical exploration has better defined the boundaries of what models are consistent with observations.  
689 As of Snowmass 2013 [14], the classification of dark matter candidates was largely based on the particle  
690 physics features of the underlying models (see Figure 5-9). Since then, focus has shifted toward exploring wide  
691 ranges of the possible phenomena in an effort to understand how well existing experimental searches cover  
692 the space of possibilities, and how new experimental opportunities provide sensitivity to regions of theory-  
693 space that are not captured by the current program [15]. There is great freedom to construct microphysical  
694 descriptions of dark matter, and a vast landscape of theoretical extensions of the Standard Model have been  
695 proposed. These models range from very simple extensions of the Standard Model containing a single new  
696 particle to complex dark sectors containing multiple dark matter states, composite dark matter blobs, or even  
697 towers of dark particles that could constitute several different components of dark matter simultaneously  
698 [16].

699 Cosmic observations set some of the strongest constraints on the properties of dark matter. Dark matter must  
700 interact gravitationally, must be produced sufficiently non-relativistically that it clusters to form galaxies,  
701 must be sufficiently long-lived that it is present in the Universe today, and must not interact frequently with  
702 the Standard Model or it would produce signals that would have been observed. Going forward, cosmic  
703 probes offer unique opportunities to learn about the properties of dark matter via the influence that it exerts  
704 on ordinary matter [3]. Measurements of the distribution of dark matter, including observables such as  
705 the matter power spectrum, the characteristics of dark matter halos (e.g., the mass spectrum, distribution,  
706 and density profiles), and the abundances of compact objects have become precise enough to place bounds  
707 on fundamental properties of dark matter such as particle mass and interaction strengths with itself and  
708 with the Standard Model. Since the last Snowmass study, there have been a series of breakthroughs in  
709 modeling cosmic structure formation in novel dark matter scenarios [17]. These studies set the basis for



732 itself can change the expectations for early cosmology, and thus the predicted abundance as a function of  
 733 the model parameters [26–32].

734 Experimental efforts over the past two decades have focused on experiments to search for the weakly interact-  
 735 ing massive particle (WIMP) and the QCD axion, both of which are strongly and independently motivated  
 736 by other unsolved scientific mysteries. In vanilla models of GeV-TeV mass dark matter associated with  
 737 solving the electroweak gauge hierarchy problem [33], WIMPs with electroweak couplings to Standard Model  
 738 particles would be naturally populated at approximately the correct density by the freeze out mechanism  
 739 described above. The QCD axion inevitably arises in Peccei-Quinn-type models in which the strong-CP  
 740 problem (the otherwise inexplicable vanishing of the neutron electric dipole moment) is solved by promoting  
 741 the strong CP-violating phase to a dynamical axion field [34]. The axion is produced via misalignment during  
 742 the QCD phase transition, which causes the field roll to vanishing angle while releasing the original vacuum  
 743 energy as axion dark matter. Much progress has been made in experimental searches on both fronts. Large  
 744 second-generation WIMP detectors based on a variety of scattering targets, combined with indirect searches  
 745 for high-energy annihilation products, have excluded  $Z$ -mediated couplings to standard model particles for  
 746 WIMP masses up to  $\sim \text{TeV}$ , and are now probing weaker couplings such as those mediated by Higgs boson  
 747 exchange. Concurrently, axion experiments have finally achieved sensitivity to the predicted QCD axion  
 748 coupling strengths by deploying various microwave quantum sensing technologies and have begun to slowly  
 749 scan the axion mass parameter space using resonant cavity detectors. A cartoon encapsulating the current  
 750 experimental situation is shown in Figure 5-10.

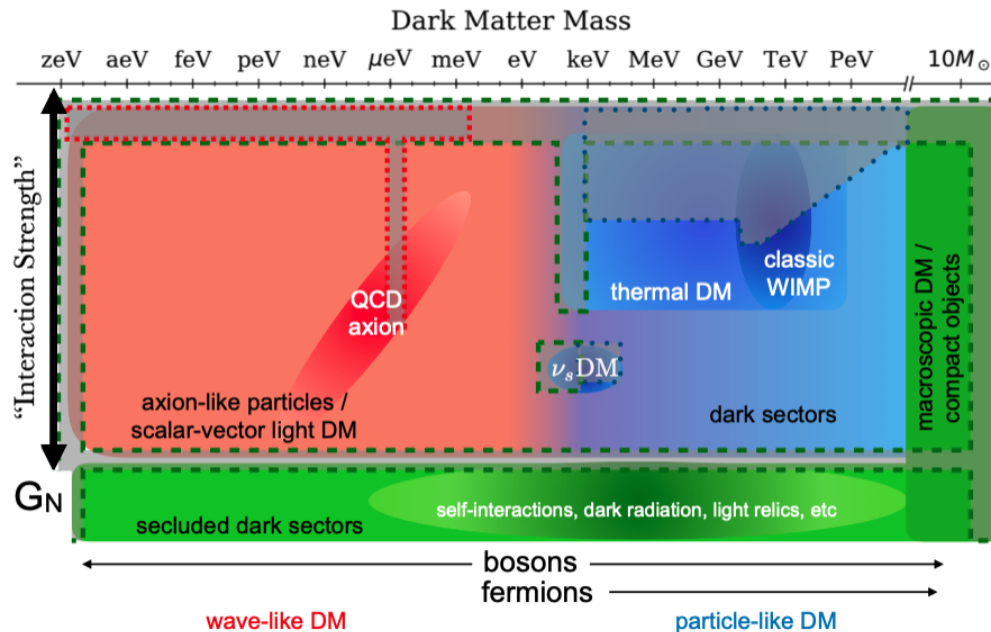


Figure 5-10. Cartoon (not to scale) of current dark matter limits.

751 While the WIMP and the QCD axion remain perhaps the most strongly motivated dark matter candidates,  
 752 the breadth of dark matter searches has expanded enormously as new tools and sensing techniques have  
 753 been identified that are able to probe large swaths of previously unexplored parameter space. Many of the  
 754 new techniques have been proposed by theorists identifying and studying models beyond WIMPs and QCD  
 755 axions, which are compatible with all experimental and observational data and furthermore have their own  
 756 well-formed dark matter origin stories.

757 A particularly compelling class of new models are those of portal dark matter in which a hidden dark sector  
 758 interacts weakly with the Standard Model sector through low dimension gauge singlet Standard Model  
 759 operators coupled to messenger particles. For example, in the case where the messenger is a vector particle  
 760 connected to the Standard Model via kinetic mixing with the hypercharge field strength portal, the result is  
 761 a massive dark photon whose interactions are predominantly proportional to electric charge scaled down by  
 762 the kinetic mixing parameter  $\varepsilon$ . Similarly, a dark Higgs can result from a scalar or pseudo-scalar messenger  
 763 whose potential contains mixing with the Standard Model Higgs mass portal. Such theories offer natural  
 764 scenarios in which the dark matter may have much smaller interactions with the Standard Model, realizing  
 765 the correct freeze out relic density for masses in the eV to GeV range, far below the weak scale. Detecting  
 766 the scattering of such low-mass dark matter requires much lower threshold calorimetric detectors sensitive  
 767 to recoil energies in the range  $\mu\text{eV} - \text{keV}$ .

768 A generic prediction of hidden-sector dark matter is that dark matter carries its own forces, which may  
 769 produce novel signals in direct and indirect detection experiments. Such a force may also change cosmic  
 770 structure formation, leading to signatures beyond the prevailing cold dark matter paradigm. For example, a  
 771 dark force could operate at the most fundamental level with a range of  $\mathcal{O}(10^{-12})$  cm, but it would change the  
 772 dark matter distribution within  $\mathcal{O}(10^{22})$  cm in galactic halos, which could be detected in cosmic observations.  
 773 In particular, dark matter self-interactions provide a compelling mechanism to produce diverse dark matter  
 774 distributions ranging from tiny dwarf galaxies to huge galaxy clusters as inferred from observations [35, 36],  
 775 which are a long-standing puzzle in cold dark matter [37]. In the next decade, Rubin LSST will measure the  
 776 dark matter distribution at unprecedented small scales and hence it will provide a unique opportunity for  
 777 probing interactions in the hidden sector [3, 38]. Furthermore, detection of additional relativistic degrees of  
 778 freedom by CMB-S4 would immediately imply the existence of a hidden sector [10].

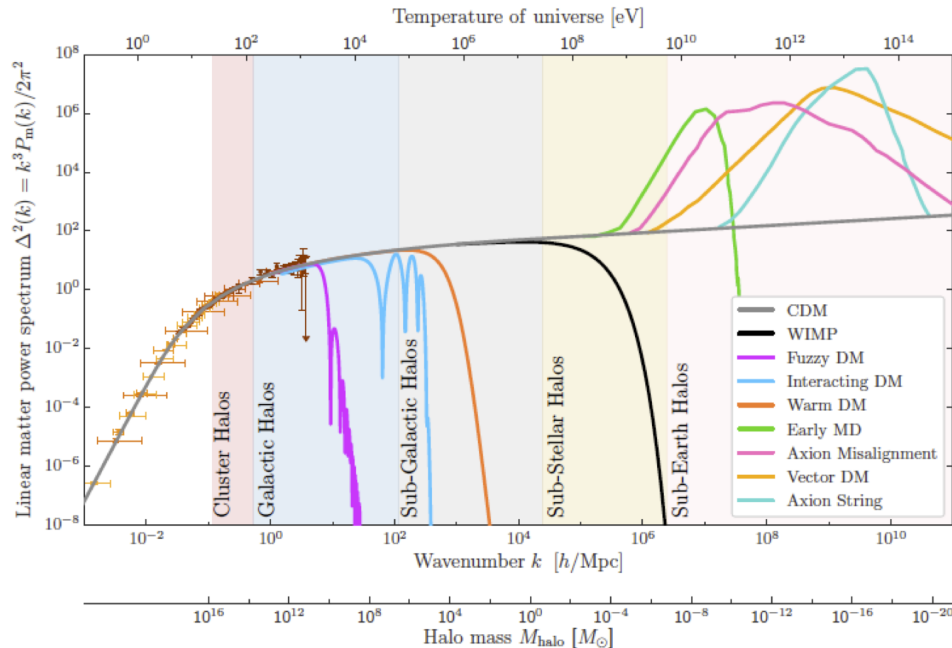
779 The range of possible dark matter masses spans at least 50 orders of magnitude from  $\sim 10^{-22}$  to  $\sim 10^{28}$  eV  
 780 for individual particles with the lower end constrained by the smallest dwarf galaxies and the upper end  
 781 constrained by the Planck mass. While composite baryonic matter (e.g., MACHOs) has been ruled out by  
 782 measurements of the CMB and Big Bang Nucleosynthesis, exotic macroscopic compact objects, including  
 783 primordial black holes, are still allowed to comprise some fraction of the dark matter up to mass scales  
 784 greater than a Solar mass ( $\sim 10^{57}$  GeV). For direct detection experiments, where the dark matter interacts  
 785 directly with the detector, different detection techniques are needed to match the energy scale and overall  
 786 detector size to the expected energy/momentum transfer and event rates for the traversing dark matter for  
 787 each regime of dark matter mass. As such, the following broad classification of types of dark matter helps  
 788 to distinguish between classes of experimental techniques that can be deployed as part of the “search wide”  
 789 strategy to achieve comprehensive coverage of dark matter parameter space.

- 790 •  $10^{-22}$  eV  $\lesssim m_\chi \lesssim 1$  eV: “Ultralight” Dark Matter
- 791 •  $1$  eV  $\lesssim m_\chi \lesssim 1$  GeV: “Light” particle Dark Matter
- 792 •  $1$  GeV  $\lesssim m_\chi \lesssim 100$  TeV: “Heavy” particle Dark Matter
- 793 •  $m_\chi \gtrsim 100$  TeV: “Ultra-Heavy” Dark Matter (UHDM)

794 In the following sections, the status and prospects for the detection of dark matter in each category will be  
 795 discussed.

### 5.5.1 Waves of ultralight bosonic dark matter ( $10^{-22}$ eV to 1 eV)

797 Provided that gravitational clumping of a single dark matter species is responsible for the formation of all  
 798 galaxies and galactic substructures, one may take from the kpc sizes of the dark matter halos hosting the  
 799 smallest observed dwarf galaxies that the mode volume of this dark matter  $(\Delta x)^3 = 1/(\Delta p)^3 \gtrsim (\text{kpc})^3$ .  
 800 Estimating the momentum dispersion from the escape velocity from this local gravitational potential well,  
 801  $\Delta p = mv_{\text{escape}}$ , a lower bound  $m \gtrsim 10^{-22}$  eV can be derived. Dark matter with smaller masses than this  
 802 “fuzzy dark matter” limit cannot be spatially localized within the volume of the smallest observed galaxies.  
 803 In addition, for dark matter mass  $m \leq 100$  eV, the mode occupation number must exceed unity in order for  
 804 there to be enough dark matter to account for the local gravitational well that formed the dwarf galaxy. One  
 805 may therefore infer from these cosmic probes that dark matter in the mass range  $10^{-22}$  eV –  $10^2$  eV must be  
 806 bosonic since the Pauli exclusion principle would prevent sufficient fermionic dark matter from fitting inside  
 807 these small galaxies. At these low masses, the local Galactic dark matter takes the form of a gas of classical  
 808 sine waves, each with large mode occupation number.

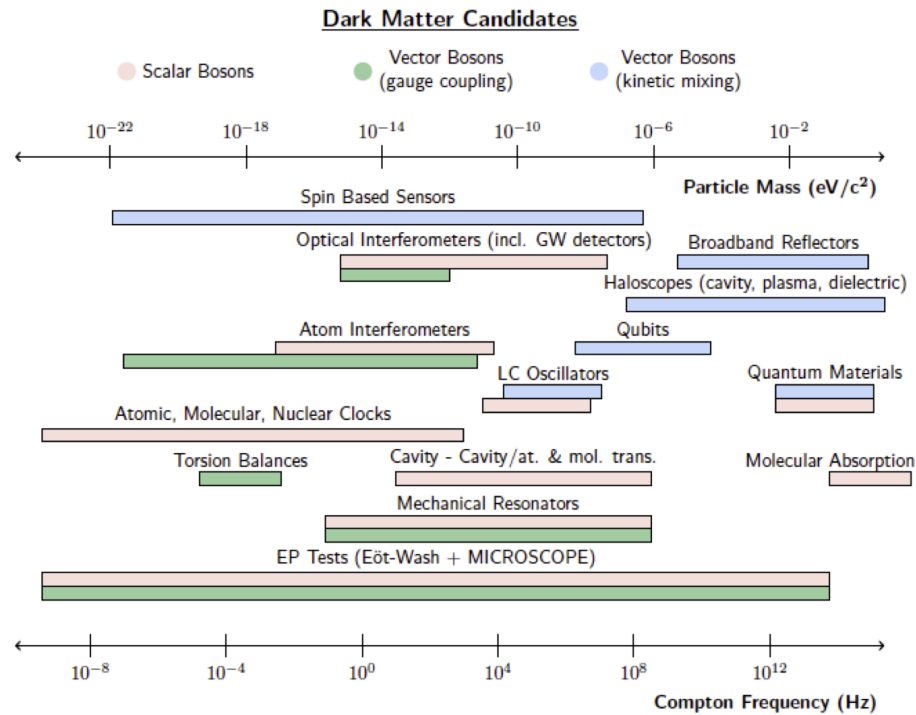


**Figure 5-11.** Measurements of the sizes of Milky way dwarf galaxies and dark halos provide information about dark matter masses and interactions.

809 In the near future, cosmic surveys will gain sensitivity to even smaller dark matter halos through measure-  
 810 ments of dwarf galaxies, stellar streams, strong lensing, and the Lyman- $\alpha$  forest [39]. While the canonical  
 811 collisionless cold dark matter model predicts that that the mass spectrum of dark matter halos extends far  
 812 below the mass at which halos are known to form luminous galaxies (Figure 5-11), ultra-light bosonic dark  
 813 matter that has a mass close to the fuzzy dark matter limit ( $\sim 10^{-22}$  eV) or fermionic dark matter that  
 814 is close to the Pauli exclusion limit ( $\sim 100$  eV) would suppress the formation and density profiles of these  
 815 small halos. Thus, cosmic surveys such as Rubin LSST have an opportunity to discover signatures of these  
 816 models through a measured absence of small halos, or further constrain the fundamental mass scale of dark  
 817 matter if small dark matter halos are observed. These are just two examples of dark matter models where



818 fundamental dark matter properties alter the distribution of dark matter in the universe in testable ways.  
 819 Similar arguments apply for a large range of specific dark matter particle models, as shown in Figure 5-11.



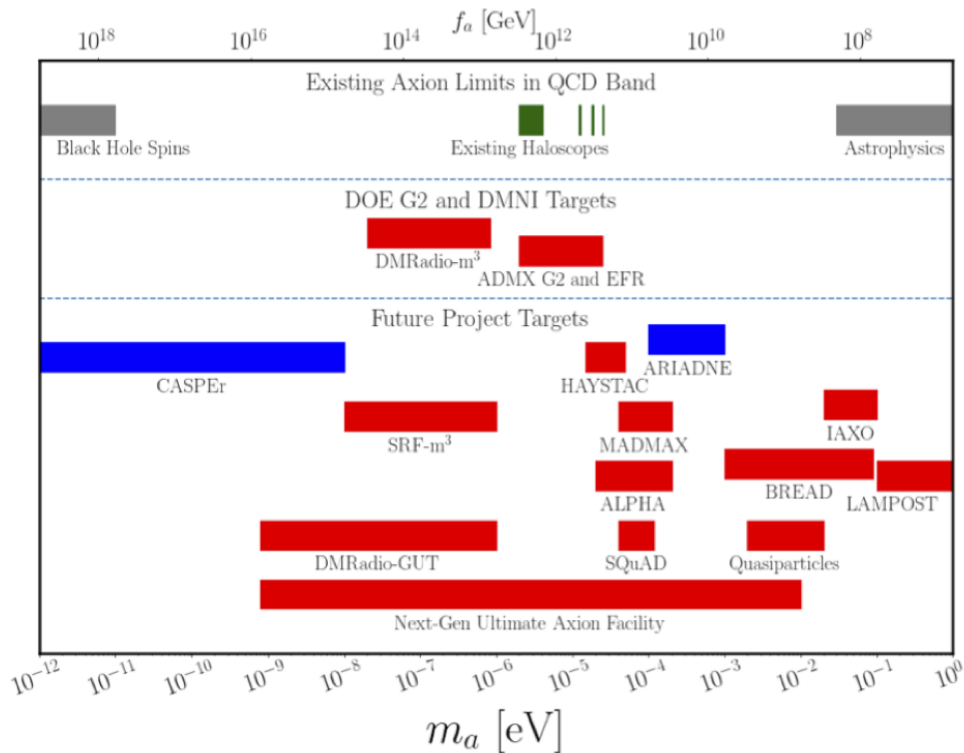
**Figure 5-12.** Ultra low mass dark matter waves still have oscillation period less than 1 year. Direct detection of periodic forces exerted by these background waves utilizes various atomic, molecular, and optical techniques. These technologies can be transferred to HEP applications where sensitivities often must be further improved beyond state of the art in order to meet the dark matter challenge.

820 Specialized terrestrial searches can also look for modulating signals due to the oscillation of the dark matter  
 821 sine wave on experimentally accessible time scales, as shown in figure 5-12. For example, the lowest mass  
 822 ( $m \sim 10^{-22}$  eV) wave may produce a detectable time-dependent classical force on test masses over an  
 823 oscillation period of approximately one year. Another class of theories couples scalar moduli dark matter  
 824 (which are ubiquitous in string theory compactifications) to the sizes of the Standard Model parameters such  
 825 that the oscillation of the dark matter wave mimics a modulation of the fundamental constants. Techniques  
 826 from the field of atomic, molecular, and optical (AMO) physics can immediately be brought to bear on  
 827 this new detection challenge, where instead of a slow drift of the fundamental constants, one optimizes the  
 828 apparatus and analysis to search for narrow line, single frequency perturbations. Tools of choice include  
 829 atomic clocks, atom interferometers, torsion balances, nuclear magnetic resonance, electric dipole moment  
 830 experiments, and gravitational wave interferometers.

831 While much of this work remains in the domain of AMO physics, in some cases, capabilities of the HEP  
 832 labs can greatly enhance the state of the art for these techniques beyond university-scale experiments in the  
 833 process of adapting them to the dark matter challenge. For example, the MAGIS pathfinder experiment will  
 834 utilize an access shaft of the NuMI beamline at Fermilab to achieve a world record 100 m vertical baseline  
 835 for atom interferometry and probe dark matter interactions with period equal to the free-fall time. An  
 836 upgraded km-scale interferometer is envisioned for the Sanford lab. Accelerator lab expertise in large scale

837 experimental deployments of beam pipes, vacuum systems, etc. may also prove useful in construction and  
 838 operations of future large optical interferometers.

839 **5.5.1.1 A high priority target – the QCD axion ( $10^{-12}$  eV to  $10^{-2}$  eV)**



**Figure 5-13.** A high priority target is the QCD axion which explains the observed vanishing of the neutron electric dipole moment as well as the origin of the dark matter from high temperature phase transitions. By virtue of a see-saw mechanism which generates the axion mass, the QCD axion model makes testable predictions for the interaction strengths as a function of mass, and a useful benchmark targets an dimensionless coupling strength of 1/3. This plot shows a suite of ongoing and future experiments which will test the QCD model by providing broad coverage of axion mass regions at the predicted coupling strengths.

840 Perhaps the most significant technological development in the field of terrestrial dark matter searches since  
 841 the last Snowmass is the demonstration by ADMX and HAYSTAC of experimental sensitivity to the invisible  
 842 QCD axion. The large improvement in signal/noise ratio achieved by these experiments was made possible  
 843 by the deployment of new quantum sensing technologies including dilution refrigerators, quantum-limited  
 844 amplifiers, and squeezed state receivers. The ability to finally begin testing the QCD axion dark matter model  
 845 has renewed both experimental and theoretical interest in the field of ultralight dark matter detection with  
 846 several new techniques proposed to cover nearly the entire range of allowed QCD axion dark matter masses.  
 847 These experiments would concurrently provide sensitivity to large areas of model-independent parameter  
 848 space for more general axion-like particles and dark photons.

849 The QCD axion which emerges from the Peccei-Quinn (PQ) solution to the strong CP problem – the 70-  
 850 year-old mystery of why the neutron has a vanishing electric dipole moment despite being made up of a



851 Fermi-sized collection of charged quarks. The PQ model begins with a standard wine-bottle phase transition  
 852 in which a global  $U(1)$  symmetry is broken at some high energy scale  $f_a$  which results in a massive Higgs mode  
 853 and a massless Goldstone mode called the axion. This axion field is coupled to the gluons of QCD and thus  
 854 identified with a dynamical CP-violating angle. Later, during the QCD phase transition at energy  $\Lambda_{\text{QCD}}$ ,  
 855 a gluon instanton condensate slightly tilts the wine-bottle potential, causing the axion vacuum expectation  
 856 value to roll to its minimum potential energy. At this true minimum, the CP-violating angle vanishes and  
 857 thus zeroes out the neutron electric dipole moment. The axion gains a tiny mass  $m_a = \Lambda_{\text{QCD}}^2/f_a$  via this  
 858 see-saw mechanism between the QCD scale and the PQ scale. Furthermore, during this vacuum relaxation  
 859 process, the potential energy associated with the original CP-violating value of the misaligned vacuum axion  
 860 field is released as ultracold dark matter.

861 The QCD axion search program is simplified in that the coupling to SM fields is proportional to the axion  
 862 mass, both being proportional to  $1/f_a$  as a result of the see-saw mechanism. A couple of benchmark models  
 863 (KSVZ and DFSZ) with order unity coupling coefficients determine a well-defined target region which cuts  
 864 a diagonal band in a more general coupling vs. mass parameter space. As shown in figure 5-13, coverage  
 865 of axion parameter space can therefore be represented via a single parameter, the axion mass, with a lower  
 866 bound of  $\sim 10^{-12}$  eV coming from the requirement that  $f_a < M_{\text{Planck}}$  for the validity of the effective field  
 867 theory. Cosmic probes of astrophysics in extreme environments further constrain the allowed QCD axion  
 868 mass range to lie between  $\sim 10^{-12}$ – $10^{-2}$  eV [40]. For example, stronger axion-nucleon couplings predicted  
 869 at larger axion masses would have caused supernova 1987a to cool more quickly than observed and would  
 870 lead to anomalous cooling of neutron stars [41, 42], while stronger axion-photon or axion-electron couplings  
 871 would be inconsistent with stellar evolution as observed in globular clusters [43–45]. Detections of black  
 872 holes with non-vanishing spins also provide constraints on axions with masses near  $10^{-12}$  eV, since axions  
 873 would populate a gravitationally bound state around black holes extracting energy and angular momentum  
 874 through superradiance [46]. Future observations of extreme astrophysical environments, along with improved  
 875 theoretical understanding of the Standard Model physics at play in these environments, promise to further  
 876 improve our sensitivity to axion physics.

877 The mass parameter space can further be split into pre-inflationary and post-inflationary axion production  
 878 mechanisms. The mass determines the characteristic cosmological time scale at which the initial axion  
 879 potential energy in the tilted potential can be released as dark matter. Because the photon and baryon  
 880 densities are redshifting away, the dark matter production time determines its relative proportion of the  
 881 overall cosmological energy density. At a mass of around  $10^{-5}$  eV, the release of an initial potential energy  
 882 density of order  $\Lambda_{\text{QCD}}^4$  would produce the observed dark matter energy fraction. At higher masses, the  
 883 energy would be released too early while the baryon density is still too high, and so the dark matter would  
 884 be underproduced. However, the potential energy could be bound to networks of topological defects which  
 885 would also prevent early release. At lower masses, the dark matter is released too late and would be  
 886 overproduced. A solution to this overproduction problem is to posit that the initial PQ phase transition  
 887 happened prior to cosmic inflation. In this case, the earth could live in an inflated patch of the universe  
 888 which had a downward statistical fluctuation in initial axionic potential energy density with correspondingly  
 889 less energy released as dark matter.

### 890 5.5.1.2 Joint probes of axion and cosmic inflation parameter space

891 The interplay of axion production histories with cosmic inflation enables complementary probes of the joint  
 892 parameter space using both axion experiments and inflation experiments. For example, if a low mass axion  
 893  $m_a \leq 10^{-5}$  eV is discovered in a direct detection experiment, this implies that the PQ symmetry is broken  
 894 at a high energy scale  $f_a = \Lambda_{\text{QCD}}^2/m_a \geq 10^{12}$  GeV. The requirement that this symmetry breaking happened  
 895 prior to cosmic inflation then forces the energy scale of inflation to be quite low,  $E_I \leq f_a$ . Furthermore,

896 even for very low axion masses for which  $f_a \rightarrow M_{\text{Planck}}$ , radiation of massless axions from the inflationary  
 897 horizon would produce isocurvature perturbations in the CMB which are strongly constrained by the Planck  
 898 experiment. In general, the detection of a low mass axion requires that  $E_I \leq 10^{14}$  GeV to avoid overproducing  
 899 axion isocurvature, but then concurrently, the Hawking radiation in gravitational waves from the horizon  
 900 would also produce too little tensor mode power to be observed in the CMB B-mode spectrum. CMB  
 901 experiments should then instead focus on improving precision on isocurvature measurements to study this  
 902 low energy scale inflation. Conversely, if inflation occurs at a high energy scale  $E_I \geq 10^{16}$  GeV, the higher  
 903 amplitude gravitational waves may first be measured in CMB experiments via the B-mode polarization  
 904 spectrum, and low mass axions would be excluded due to overproducing isocurvature. The axion direct  
 905 detection program should then focus on higher mass axions in which the PQ symmetry is broken after  
 906 inflation so that no massless isocurvature modes can be produced.

### 907 5.5.1.3 Wave dark matter status and strategies for direct detection

908 In the U.S., the DOE-supported ADMX-G2 resonant cavity experiment has excluded the QCD axion in the  
 909 axion frequency range 645 MHz - 1.1 GHz (2.66-4.2  $\mu\text{eV}$ ) at even the more pessimistic DFSZ axion-photon  
 910 coupling strength, and will continue to scan upward to 2 GHz using nearly quantum limited preamplifiers  
 911 in their radio receiver. The NSF-supported HAYSTAC experiment has deployed a squeezed state quantum  
 912 receiver and demonstrated KSVZ coupling sensitivity in a narrow band near 4.1 GHz (17  $\mu\text{eV}$ ).

913 Two new small-scale U.S. axion experiments are being developed within the DOE's Dark Matter New  
 914 Initiatives (DMNI) program for small projects. The first is ADMX-Extended Frequency Range (ADMX-  
 915 EFR) which will increase the frequency scan rate by deploying and simultaneously operating multiple cavities  
 916 within a single magnet bore, targeting higher frequencies 2-4 GHz. The other is DMRadio- $\text{m}^3$  which will  
 917 search over a wide range of lower frequencies and lower axion masses  $< 250$  MHz (1  $\mu\text{eV}$ ) using a toroidal  
 918 antenna incorporated into a lumped element LC circuit to couple to the axion waves. In each case, the total  
 919 signal power will scale with the volume of the magnetized region which serves as the detector target for  
 920 axion-photon conversion. Decisions on new starts within the DMNI program are expected to happen before  
 921 the P5 report.

922 Meanwhile, private foundation-funded R&D continues on the CASPER-Electric and CASPER-Wind nuclear  
 923 magnetic resonance experiments which aim to measure the dark matter-induced modulation of the nucleon  
 924 electric dipole moment and the axion-nucleon spin coupling respectively. The ac effects on the spin energy  
 925 levels can be resonantly detected by modulating an applied electric field for the ac EDM measurement or by  
 926 tuning the spin precession frequency to match the axion wave frequency for the spin coupling measurement.  
 927 Both experiments are currently targeting lower frequencies/masses below 1 MHz and have a longer term  
 928 goal of reaching sensitivity to the tiny signals predicted for the QCD axion. This R&D is critical as only  
 929 by observing the axion in multiple, predicted detection channels can it be distinguished from more general  
 930 dark matter models.

931 Various quantum sensing technologies are being transferred to the dark matter axion search. For example,  
 932 the low frequency experiments are limited by thermal noise. In this case, single quantum mixers used can be  
 933 used to transduce signal photons to higher frequency ranges where quantum limited amplification is available.  
 934 The low level of quantum noise relative to thermal noise then allows the resonant search to occur at constant  
 935 ratio of signal to thermal noise over many Lorentzian linewidths to recover some of the speed advantage of a  
 936 broadband search – the same strategy as was used in the Weber bar technique originally used to search for  
 937 gravitational waves. Another strategy is to utilize an energized superconducting rf cavity (SRF) and use the  
 938 ac magnetic field instead of a dc magnetic field to upconvert the low frequency axion wave into an rf signal  
 939 photon.

940 At signal frequencies higher than those being probed by ADMX, the experimental challenge is that the  
941 size of the resonant cavities must shrink with increasing frequency to match the smaller wavelength of  
942 the signal photons, and so the signal power shrinks according to this rapidly decreasing target volume.  
943 While ADMX-EFR will pack many cavities into a single magnet to try to recover the lost volume, other  
944 ideas like ALPHA, the Axion Longitudinal Plasma Haloscope will try to create a large volume metallic  
945 metamaterial whose plasma frequency can be tuned into resonance with the dark matter wave. Alternatively,  
946 instead of attempting to maintain or increase the signal level, another strategy is to drastically reduce  
947 noise by using photon counting detectors which can have arbitrarily low dark count rates. The resulting  
948 noise from background count fluctuations can be far below the zero-point noise that is incurred by the  
949 currently employed frequency-resolved power readout techniques. New, ultra-low noise single microwave  
950 photon detection techniques based on superconducting qubits and on Rydberg atoms are being transferred  
951 from neighboring fields of quantum computing and AMO. The Superconducting Qubit Advantage for Dark  
952 Matter (SQuAD) prototype has demonstrated a factor of nearly 40 noise reduction relative to zero-point  
953 noise which, combined with high-Q dielectric cavities could enable QCD axion searches up to 20 GHz. R&D  
954 is also underway with the Rydberg Atoms at Yale (RAY) experiment to perform single photon counting at  
955 even higher microwave frequencies.

956 Finally, broadband techniques are also being developed which use large impedance mismatches at metallic  
957 or dielectric plates immersed in magnetic field to create axion to photon transition radiation. This radiation  
958 can be focused, even by the plate geometry itself, onto low noise single photon detectors. In Germany, a  
959 large R&D effort is underway to develop the MADMAX experiment which seeks to deploy a set of aligned  
960 plates to create multiple scattering targets within a large, high-field dipole magnet bore and read out signal  
961 power coherently using HEMT amplification. In the U.S., smaller-scale experiments such as BREAD and  
962 LAMPOST are performing R&D funded by DOE-OHEP's QuantISED program to achieve high sensitivity  
963 to smaller signal levels via ultra low noise photon counting detectors such as quantum capacitance detectors,  
964 MKIDs, and SNSPDs. This new quantum photon counting technology covering 10 GHz - 100 THz may find  
965 use in various other areas of HEP outside of dark matter.

966 As many of these direct detection technologies are nearing maturity and are already or will be shovel-ready  
967 in the next decade, there is a key opportunity to cover most of the remaining  $\sim 10$  orders of magnitude in the  
968 QCD axion mass parameter space while concurrently exploring much of dark photon parameter space. The  
969 detector and sensor component of most of the experimental techniques falls firmly in the small experiment  
970 category and so a portfolio of small experiments could be strategically devised to divide and conquer the  
971 remaining axion parameter space. However, a critical path and long lead time item for most experiments is  
972 the large bore, high field magnet needed to induce the interaction between axion waves and signal photons.  
973 While resonant cavity experiments are viable in the frequency range 1-30 GHz and have been operated with  
974 commercial research solenoids or MRI magnets, the frequency scan speed has been in practice limited by the  
975 availability of magnet space. For example, if ADMX could operate two magnets instead of one, then nearly  
976 identical radios deployed in each magnet could be tasked to scan through independent frequency ranges to  
977 cover twice as much mass range in the same amount of time. Unlike most other techniques, these frequency  
978 scanning experiments enjoy linear scaling of science return with project cost which is usually dominated by  
979 magnet cost.

980 While the resonant experiments are constrained in size to the wavelength scale of the signal, lumped element  
981 experiments targeting lower frequencies and broadband antenna targeting higher frequencies each rely on  
982 using larger, customized magnets to intercept enough of the dark matter flow to achieve sensitivity to the  
983 QCD axion. A serious effort to comprehensively cover the many decades in allowed axion mass parameter  
984 space in finite operations time will thus require a significant investment in magnet facilities both in large,  
985 custom magnets for signal-limited experiments and a fleet of smaller commercial magnets for the resonant  
986 cavity experiments. Given that a Mu2e or CMS-scale solenoid may take a decade to construct, and even

987 smaller superconducting magnets take several years to fabricate, a possible path forward is for the magnets to  
 988 be planned, constructed, and operated as user facilities in which multiple experiments could be concurrently  
 989 or sequentially deployed. A HEP magnet facility would focus on large bore magnets for longer term tenants,  
 990 in contrast to the current portfolio of small bore, ultra high-field magnets that provided for shorter term  
 991 condensed matter experiments by the NSF’s National High Magnetic Field Facility.

#### 992 5.5.1.4 Complementarity with gravitational wave searches for low-mass cosmic bosons

993 Gravitational waves have been recently shown to be an excellent new cosmic probes of ultralight bosonic  
 994 matter. For example, if the Peccei-Quinn phase transition is sufficiently strongly first order, then a stochastic  
 995 background of gravitational waves could be created by the boiling of the vacuum during this phase transition.  
 996 Measuring the spectrum of this stochastic background with LIGO or with future GW observatories would  
 997 determine the energy scale of the PQ transition and then via the axion model, uniquely determine the mass of  
 998 the axion dark matter. The axion direct detection program could then be focused on a narrow mass window  
 999 rather than needing multiple experiments with different technologies to cover each of the  $\sim 10$  decades the  
 1000 currently allowed mass window.

1001 Searches for narrow line gravitational wave signals also constrain ultralight bosons of any type due to the  
 1002 prediction of the formation of gravitational atoms around spinning black holes. Just as in the hydrogen atom,  
 1003 particles bound to a black hole are confined to orbits of quantized angular momentum. Since bosons do not  
 1004 have a Pauli exclusion principle, these orbits become classically occupied with boson condensates which are  
 1005 created from the gravitational superradiance effect when the boson Compton wavelength is comparable to  
 1006 the black hole size. Transitions of bosons between different orbitals then results in narrow line gravitational  
 1007 wave radiation. A range of current and planned GW observatories can conduct model-independent searches  
 1008 for low-mass bosons from  $10^{-21}$ – $10^{-11}$  eV.

### 1009 5.5.2 Light particle dark matter (1 eV to 1 GeV)

1010 Above the  $\sim$ eV scale, the dark matter typically manifests as individual quanta. A priority target is the  
 1011 sterile neutrino, whose existence is suggested by the need to incorporate neutrino masses in the Standard  
 1012 Model, and which typically becomes sufficiently long-lived to play the role of dark matter at masses below  
 1013 around 10 keV. If sufficiently mixed with the active neutrinos, sterile neutrinos with masses below a few keV  
 1014 would thermalize as warm dark matter, and would suppress structure formation due to free-streaming to a  
 1015 degree that is inconsistent with cosmic observations. Evading this cosmological bound requires a mechanism  
 1016 to produce the dark matter in a very cold state, and appropriately weak coupling to the Standard Model such  
 1017 that it does not kinetically thermalize. The power of these constraints from structure formation highlight  
 1018 the key role that cosmic probes play in defining the viable parameter space for dark matter models.

#### 1019 5.5.2.1 Direct Searches

1020 For sub-GeV-mass dark matter, direct detection technologies generally rely on individual dark matter  
 1021 particles scattering or being absorbed on various target media in detectors with low energy and momentum  
 1022 threshold. In scattering processes, the signal energy is at most the initial dark matter kinetic energy and the  
 1023 momentum transfer must be less than twice the initial dark matter momentum. For typical Galactic dark  
 1024 matter velocities of  $v = 10^{-3}c$ ,  $E_{\max} \approx 10^{-6}M_{\text{dm}}$ , and  $p_{\max} \approx 10^{-3}M_{\text{dm}}$ , detectors with energy thresholds  
 1025 in the meV–keV range are needed to search for scattering of dark matter in the mass range keV–GeV. At

1026 these lower dark matter masses, kinematic matching for efficient billiard-ball scattering requires using lower  
1027 mass nuclei, electrons, or collective modes such as phonons or magnons as scattering targets.

1028 The universe itself provides an exceptional calorimeter in the form of the primordial baryon–photon plasma  
1029 prior to recombination. Scattering between light dark matter and nucleons and/or electrons would transfer  
1030 heat and momentum between the Standard Model and dark matter. This would have the effect of heating the  
1031 dark matter and resulting in a suppression of the matter power spectrum similar to that seen in warm dark  
1032 matter, which would be observable in the CMB, Lyman- $\alpha$  forest, and dwarf galaxies [47–51]. These cosmic  
1033 probes provide access to the full range of light dark matter masses at large scattering cross sections that  
1034 inaccessible to terrestrial experiments due to shielding by the Earth’s atmosphere. Expected improvements  
1035 will come from current and near-future cosmic survey experiments like DESI, Rubin LSST, CMB-S4, and  
1036 Spec-S5 [3, 18].

1037 In terrestrial experiments, noble liquid detectors have achieved signal thresholds of 20 eV when detecting  
1038 electrons in the ionization channel, demonstrating the lowest background rate per unit target mass in this  
1039 signal energy range. Various ways of doping the target medium are being explored in order to increase  
1040 scintillation and ionization yield and to improve kinematic response to lower mass particles with low-A  
1041 nuclei. Odd-A nuclei also provide sensitivity to spin-dependent interactions.

1042 Semiconductor-based detectors have a threshold near the typical band gap energies of 1 eV in the electron  
1043 scattering channel, providing sensitivity to dark matter masses near 1 MeV, approximately an order of mag-  
1044 nitude below that of the noble liquid detectors. A recent significant development has been the demonstration  
1045 of sub-electron charge resolution using skipper CCDs which perform repeated, non-destructive measurements  
1046 of the charge in individual pixels to average away the readout noise. The ongoing SENSEI and DAMIC-  
1047 M experiments are planning a factor of 10–100 scale-up in silicon CCD target mass in a jointly proposed  
1048 experiment Oscura, part of the the DOE’s DMNI program. Meanwhile, semiconductor crystal detectors  
1049 such as SuperCDMS HVeV and EDELWEISS are also reaching single electron sensitivity using voltage bias  
1050 across the semiconductor diode to increase the single electron signal energy.

1051 To further lower thresholds for the electron scattering channel, novel target materials with collective electronic  
1052 excitations with meV scale bandgaps are being investigated. For example, superconducting nanowire single  
1053 photon detectors can also be used to detect energy deposited from dark matter scattering events with trigger  
1054 thresholds near 250 meV for localized energy deposits which quench the bias supercurrent. Dark matter can  
1055 also have electric couplings to the oscillating dipoles exhibited by optical phonon modes in polar crystals.  
1056 The 10-100 meV scale bandgap of optical phonons provides both a low detection threshold and a large phase  
1057 velocity for good kinematic matching to the relatively fast 300 km/s speeds of the galactic dark matter  
1058 distribution. Engineered semiconductor materials with small Dirac cone band gaps such as heavy fermion  
1059 materials are also being developed. A key challenge is scaling up these 2-d materials to sufficiently large  
1060 target masses to provide deep coverage of dark matter parameter space. Collective plasmon modes including  
1061 those in new heavy fermion materials may also be better kinematically matched to low mass dark matter  
1062 scattering.

1063 In the phonon channel, the semiconductor and crystalline detectors such as CRESST, EDELWEISS, Su-  
1064 perCDMS, and MINER are targeting eV or sub-eV resolution by reducing the readout noise in transition  
1065 edge sensors. To achieve strong phonon response for lower mass dark matter scattering, SPICE, part of the  
1066 TESSERACT program also funded through DOE DMNI is exploring targets made of polar crystals which  
1067 have strong optical phonon response to dark matter interactions mediated by dark photons. The optical  
1068 phonon band gap of around 10–100 meV provides the low detection threshold when paired with a sufficiently  
1069 low noise readout. Beyond the NTDs and TESs traditionally used to read out microcalorimeters, new sensor  
1070 technologies being explored include kinetic inductance detectors, magnetic microcalorimeters, SNSPDs, and  
1071 superconducting qubits, with some of this next-generation technology being funded through various DOE

1072 quantum programs. While development of ultra-low-threshold quantum sensors may justifiably be funded  
 1073 through dedicated quantum programs, the deployment of these sensors and the studies of dark matter  
 1074 detector response including both measurements and simulations, and the characterization of new low energy  
 1075 backgrounds are uniquely HEP activities that could best be carried out in small, pathfinder experiments.  
 1076 As in the case of the axion searches described above, the DOE DMNI program currently provides funding  
 1077 to develop a limited number of proposed techniques. Given the variety of new sensing technologies, an  
 1078 opportunity exists to cover broad swaths of unexplored sub-GeV dark matter parameter space by expanding  
 1079 the portfolio of small experiments.

1080 In addition to sensors designed to measure excitations of internal bulk phonon modes, a complementary  
 1081 approach is to monitor the center-of-mass motion of a composite object [52]. Optical or electrical readout  
 1082 of such mechanical sensors, in either the classical or quantum regime, offers a novel approach to detecting  
 1083 particle dark matter on a wide range of mass scales, including both heavy [53] and light [54] candidates.  
 1084 These sensors are naturally directionally dependent because they monitor changes to three-momentum of  
 1085 the mechanical system, and can reach energy thresholds well below the eV scale [54].

1086 Beyond the technical capabilities of any particular detector, the distribution of deposited energies is also  
 1087 sensitive to the microphysics, being largely determined by nature of the mediator connecting the dark  
 1088 matter to the Standard Model. Heavy mediators manifest as non-renormalizable interactions which favor  
 1089 larger energy transfer (up to the kinematic limit), whereas light mediators enhance lower transferred energy.  
 1090 Figure 5-14, shows four cases of mediators, corresponding to heavy mediators with preferred coupling to  
 1091 electrons or nuclei, and light mediators coupled to electrons or in the form of dark photons. In each panel,  
 1092 existing bounds are shaded in tan, potential advances in the the probed parameter space of dark matter  
 1093 mass versus coupling/cross section in the near term ( $\sim 5$  years) based on existing research investments are  
 1094 shaded green, whereas those achievable in the far term (longer time scales) are shaded blue. Each shaded  
 1095 region represents the combined footprint of several experiments.

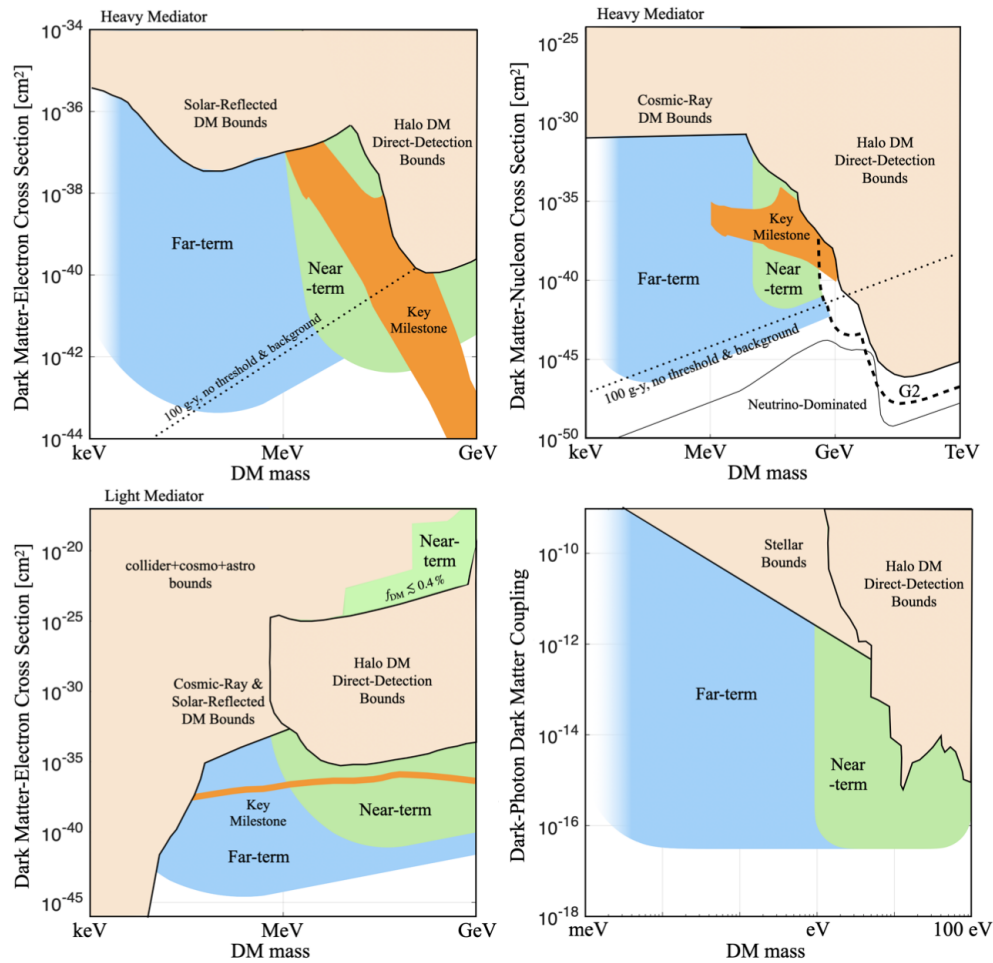
### 1096 5.5.2.2 Indirect Searches

1097 Low-mass particle dark matter can be detected through its decay or annihilation into Standard Model  
 1098 messengers. Below a GeV, the available Standard Model channels include photons, neutrinos, electrons,  
 1099 muons, and the lightest hadrons. Because of poorly understood astrophysical backgrounds and confounding  
 1100 factors such as Galactic magnetic fields (which scramble the incident direction of charged particles), photons  
 1101 with energies  $\gtrsim$  keV are typically considered among the most promising messengers.

1102 Indirect searches are uniquely powerful for probing the dark matter lifetime. A compilation of constraints  
 1103 based on the searches for X-ray and gamma-ray signals across many decades of mass is shown in Figure 5-15.  
 1104 It is striking that such constraints require the dark matter to live many orders of magnitude longer than the  
 1105 current age of the Universe.

1106 In thermal freeze-out scenarios, indirect detection probes the same interactions that fix the dark matter  
 1107 abundance in the early universe. In the sub-GeV mass range, current limits already generically rule out the  
 1108 simple thermal freeze-out scenario for  $s$ -wave annihilation, unless the dark matter annihilation products are  
 1109 almost exclusively neutrinos, significantly constraining the space of viable dark matter candidates that were  
 1110 once in thermal contact with the Standard Model (e.g., [55]). In particular, constraints from measurements  
 1111 of the CMB [23] have a key role in establishing the viable portals by which dark sectors could communicate  
 1112 with the Standard Model.

1113 At the low end of this energy range, X-ray telescopes have placed stringent constraints on sterile neutrinos,  
 1114 complementary to cosmic probes of warm dark matter [56]. Upcoming X-ray telescopes (XRISM [57, 58],

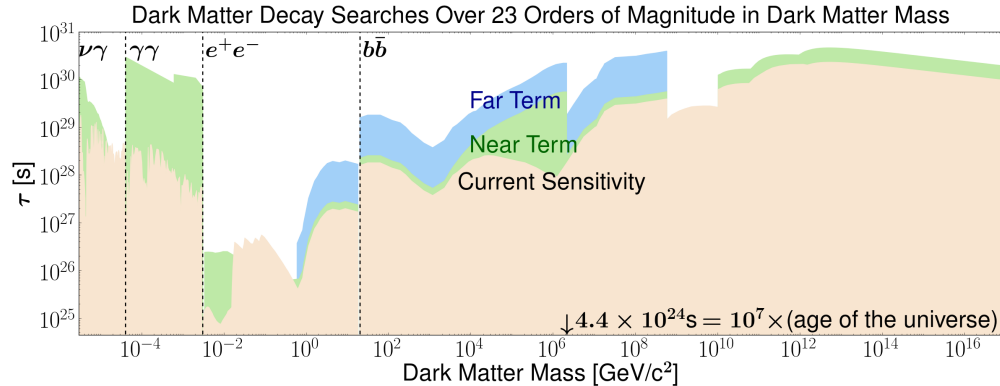


**Figure 5-14.** Predicted near-term and far-term sensitivities to dark matter with mass in the range 1 eV–1 GeV through different types of mediators connecting to the Standard Model. Key milestones indicate representative target regions in which the dark matter is thermally produced for each scenario. Near-term projections are based on demonstrated sensitivities to electron recoil and nuclear recoil in the 1 eV and 10 eV range respectively. As the maximum energy deposit in scattering is the initial kinetic energy of the dark matter of order  $10^{-6} M_{\text{DM}}$ , far term projects with lower energy thresholds will be enabled with ongoing R&D in novel target materials with band gap below the 1 eV scale of semiconductors and of chemistry. Examples include optical phonons in polar materials with band gap 10–100 meV and superconductors and other quantum materials with sub-meV gap. Concurrently, demonstration of low threshold sensor technology, including those based on quantum sensors will be required.

1115 Micro-X [59], Athena [60], HEXP-P [61]) will further improve constraints on decaying/annihilating light DM  
 1116 in general, and sterile neutrinos in particular; some will have sufficient energy resolution to seek to resolve  
 1117 DM-sourced spectral lines.

1118 At higher energies, there is currently a sensitivity gap in the MeV–GeV gamma-ray band. The last major  
 1119 experiment was NASA’s Imaging Compton Telescope (COMPTEL)[62], which operated from 1991–2000.  
 1120 Several proposed future experiments aim to address this gap: the Compton Spectrometer and Imager (COSI)  
 1121 [63] has a planned launch date in 2025, and will survey the gamma-ray sky at energies of 0.2–5 MeV; AMEGO-





**Figure 5-15.** Constraints on the lifetime of dark matter based on null searches for produced X-rays and gamma rays, as well as planned improvements in the near term (green) and far term (blue). The sharp boundaries are artifacts of analysis, rather than instrumental thresholds.

1122 X [64] has a planned satellite launch date in the late 2020s and will probe energies from 25 keV–1 GeV using  
 1123 a silicon pixel tracker, CsI calorimeter, and plastic anti-coincidence detector; and e-ASTROGAM [65] would  
 1124 target the 0.3 MeV–3 GeV energy range using a similar approach. Other proposals focused on the MeV  
 1125 band include SMILE [66–69], GRAMS [70], and GammaTPC, which rely on gaseous or liquid time projection  
 1126 chambers (TPCs), and GECCO [71] (based on a novel CdZnTe imaging calorimeter with a deployable coded  
 1127 aperture mask).

1128 Investment in such experiments would enable both data-driven studies of backgrounds relevant for lower  
 1129 and higher energy indirect searches, as well as providing greater sensitivity to decaying and/or annihilating  
 1130 light dark matter in the MeV–GeV mass range. They have the potential for sufficient sensitivity to probe  
 1131 thermal freeze-out even when the dominant annihilation is  $p$ -wave [56, 72] (which is suppressed at non-  
 1132 relativistic velocities and thus challenging to observe via indirect detection), further complementing terrestrial  
 1133 experiments focusing on this mass range.

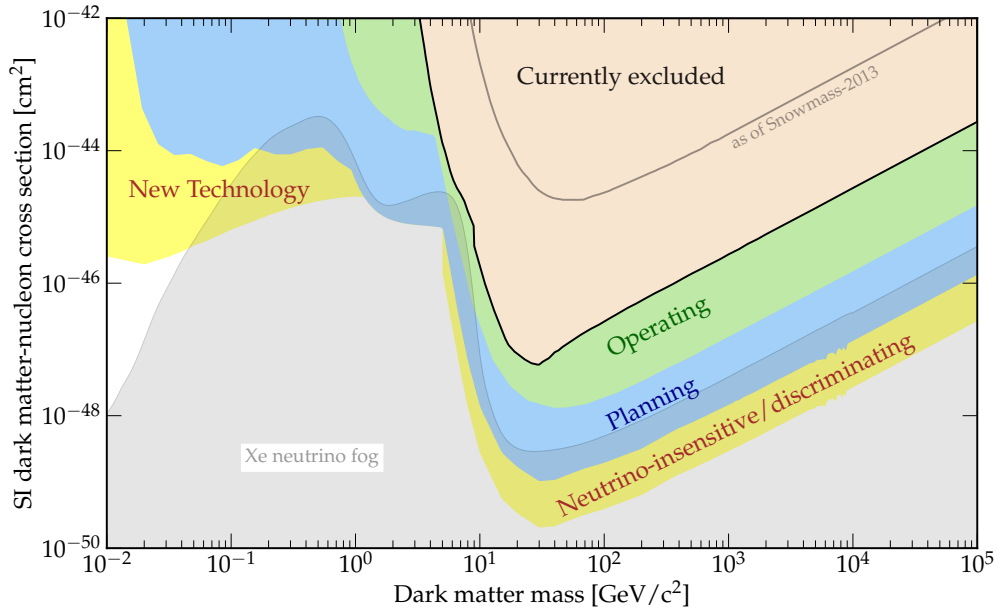
### 1134 5.5.3 Classic WIMPs (1 GeV to $\sim 100$ TeV)

1135 Dark matter masses between  $\sim$  GeV and hundreds of TeV populate the classic WIMP parameter space,  
 1136 where theories aimed at explaining the electroweak hierarchy typically reside. Such theories have been the  
 1137 subject of extensive theoretical exploration [73–78], driven both by the connection to electroweak physics as  
 1138 well as the fact that roughly electroweak-sized couplings naturally lead to a relic abundance close to the one  
 1139 required by cosmological observations through freeze-out. As a result, searches for WIMPs remain extremely  
 1140 well-motivated, with many interesting models inhabiting viable and testable parameter space [79, 80].

#### 1141 5.5.3.1 Direct Searches

1142 For these masses, the most efficient scattering is with nuclei, which can be described as an expansion in  
 1143 the dark matter velocity by terms in an effective field theory [81–84] Two classes of interactions dominate  
 1144 in the non-relativistic limit: spin-dependent (SD) scattering, which requires a nuclear target with non-zero  
 1145 spin, and spin-independent (SI) scattering, which is coherently enhanced in many models for large- $A$  targets.



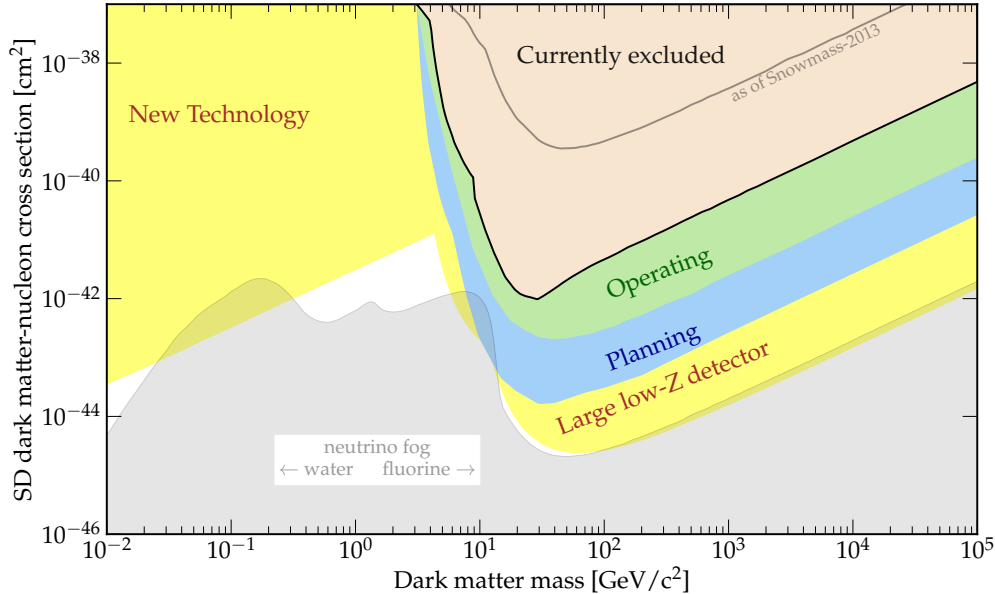


**Figure 5-16.** Combined Spin-independent dark-matter nucleon scattering cross section space. Current 90% c.l. constraints are shaded beige, while the reach of currently operating experiments are shown in green (LZ, XENONnT, PandaX-4T, SuperCDMS SNOLAB, SBC). Future experiments are shown in blue (SuperCDMS, DarkSide-LowMass, SBC, XLZD, ARGO) and yellow (Snowball and Planned $\times$  5). The neutrino fog for a xenon target is shaded light grey. (Reproduced from Ref. [87]).

1146 Detectors employing heavy nuclei to target SI scattering have seen rapid advance in their size and control of  
 1147 backgrounds, achieving many decades of improved sensitivity over the past twenty years. Such experiments  
 1148 are now within an order of magnitude of sensitivity to the expected background neutrinos produced in  
 1149 the atmosphere, Sun and supernovae. This ‘neutrino fog’ is interesting to measure in its own right, but  
 1150 will require new technologies and analysis techniques to be distinguished a dark matter signal, such as the  
 1151 expected annual modulation of the DM rate [85] and directional detectors (see *e.g.* the review found in [86]).  
 1152 Directional DM detectors would further open up the tantalizing potential to map out the local DM structure  
 1153 of the Galaxy.

1154 Currently, the most stringent SI bounds for  $\gtrsim 10$  GeV dark matter masses are from liquid xenon (LXe) and  
 1155 liquid argon (LAr) detectors employing time projection chamber (TPC) technology. Charge and phonon  
 1156 detectors are able to reach lower DM masses and are expected to reach sensitivity for dark matter masses  
 1157 between 0.5–5 GeV to within a decade of the neutrino fog, limited by cosmogenically activated isotopes  
 1158 including  $^3\text{H}$  and  $^{32}\text{Si}$ ,  $^{210}\text{Pb}$ , and by dark counts. Phase change detectors such as bubble chambers can  
 1159 provide powerful background-discriminating technology and the combination of scalability, low threshold,  
 1160 and background discrimination *at low threshold* could allow noble-liquid bubble chambers to explore the  
 1161 neutrino fog in the 1–10 GeV WIMP mass range [88].

1162 Figure 5-16 shows the current, operating, and future projected 90% CL constraints for WIMP DM interacting  
 1163 spin-independently [87]. The beige regions show the current best exclusion for each mass, combining  
 1164 previously reported limits [89–104], as collected by Ref [105]. A measure of the incredible progress of  
 1165 the past decade is provided by the comparison with the grey line from Snowmass 2013 [106]. The projected



**Figure 5-17.** Combined Spin-dependent dark-matter nucleon scattering cross section space for scattering with neutrons or protons. Current 90% c.l. constraints are shaded beige, whereas the reach of currently operating experiments are shown in green (LZ, XENONnT). Future experiments are shown in blue (PICO-500, XLZD) and yellow (Snowball, PICO-100 ton). The neutrino fog for a water or fluorine target is shaded light grey. (Reproduced from Ref. [87]).

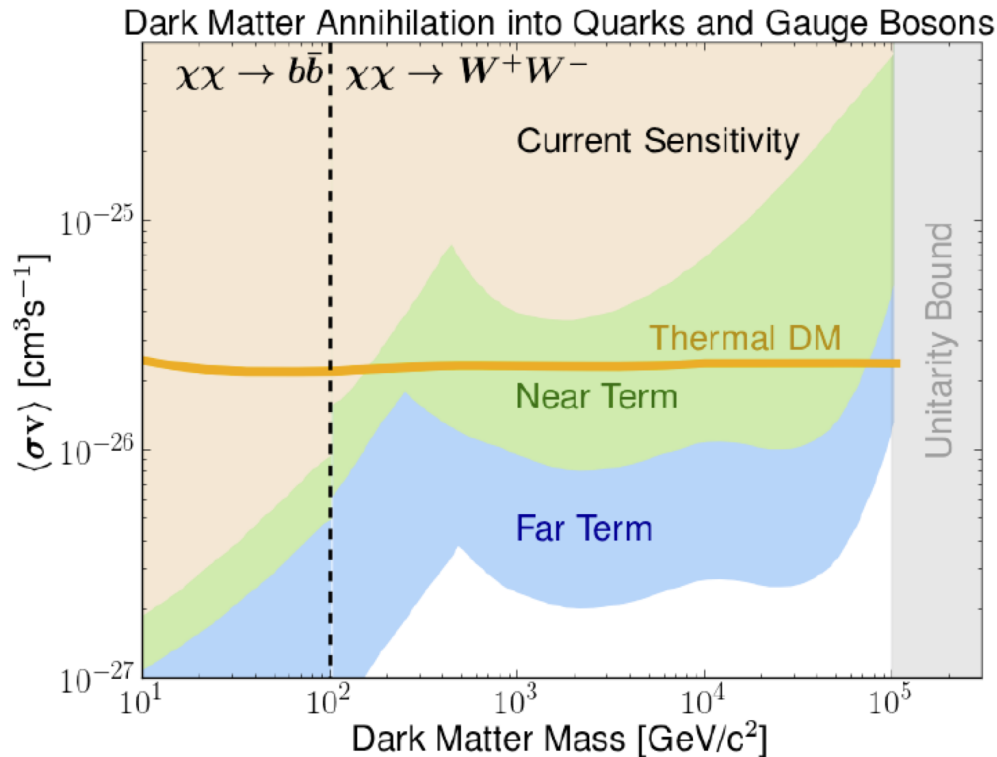
1166 reach of currently operating experiments is shaded in green, whereas proposed upgrades are shown in blue  
 1167 (including SuperCDMS, DarkSide-LowMass, SBC, XLZD, and ARGO) and the reach of proposals based on  
 1168 new technologies are shaded in yellow (Snowball and Planned $\times$  5). The region corresponding to the neutrino  
 1169 fog for a xenon target is shaded light grey.

1170 Similarly, Figure 5-17 shows the current, operating, and future projected 90% CL constraints for WIMP  
 1171 DM interacting spin-dependently [87]. Shaded in beige is the union of the currently excluded parameter  
 1172 space, led by LXe TPCs and freon-based bubble chambers [92, 107–109, 109–117], as collected by Ref [105].  
 1173 The expected reach of currently operating experiments is shaded green, and future proposed experiments  
 1174 are shaded blue (including PICO-500 and XLZD) and yellow (Snowball and PICO-100 ton). The region  
 1175 corresponding to the neutrino fog for a water or fluorine target is shaded light grey.

1176 Taken together, Figures 5-16 and 5-17 illustrate the exciting prospects for direct searches in the WIMP  
 1177 regime, with a suite of experiments based on different technologies capable of probing WIMP dark matter  
 1178 all the way down and into the neutrino fog.

### 1179 5.5.3.2 Indirect Searches

1180 Indirect searches for dark matter play a very important role in covering the parameter space. At lower  
 1181 masses, as discussed above, measurements of the CMB have a key role in constraining the viable portals to  
 1182 communicate with the SM. At the highest masses, indirect detection achieves sensitivity to WIMP parameter



**Figure 5-18.** Limits on WIMP annihilations into pairs of bottom quarks (for masses below  $\sim 100$  GeV and  $W$  bosons (for larger masses) based on null searches by  $\gamma$ -ray observatories. The beige regions indicate the current limits for each mass, whereas the green shaded region indicates near future gains based on planned missions, and the blue shading indicates the reach that would be enabled by long term investments in ground- and space-based observatories.

1183 space that is inaccessible to direct or collider searches, and in general provides a broadly model-agnostic probe  
 1184 of thermal freezeout scenarios. The most effective messengers of WIMP annihilation are gamma rays (which  
 1185 point back to their origin, giving an additional analysis handle) and energetic anti-matter signals in cosmic  
 1186 rays. Especially for higher DM masses, where abundant energy is available to produce the full range of SM  
 1187 particles, indirect signals are generically expected to be multi-messenger and multi-scale [118].

1188 The rate of annihilation depends on the underlying microphysics, which determines both the rate into each  
 1189 specific annihilation channel as well as its dependence on the relative velocity between the annihilating dark  
 1190 matter particles. Dark matter making up the halos of galactic structures is typically highly non-relativistic  
 1191 ( $v \sim 10^{-4} - 10^{-3}$ , and thus  $s$ -wave annihilations (for which  $\langle\sigma v\rangle$  is  $v$ -independent) typically dominate over  
 1192 annihilations at higher partial wave. Additional challenges involve determining the dark matter distribution  
 1193 along the line of sight of an indirect search, which strongly impacts the rate of annihilation, since  $\langle\sigma v\rangle \propto \rho_\chi^2$ .  
 1194 This systematic uncertainty benefits strongly from advances in simulation of galaxy formation. Modeling of  
 1195 astrophysical backgrounds with sufficient precision as to be able to distinguish a subtle signal of dark matter  
 1196 annihilation or decay from more prosaic astrophysical processes is also a key ingredient (see Refs. [1, 119, 120]  
 1197 for further discussion); this is already manifest in a number of puzzling excesses in indirect searches [119].

1198 Gamma rays may be detected both by space- and ground-based telescopes. The space-based Fermi-LAT  
 1199 is optimized to reconstruct gamma rays from 1 – 100 GeV and is effectively able to observe the entire

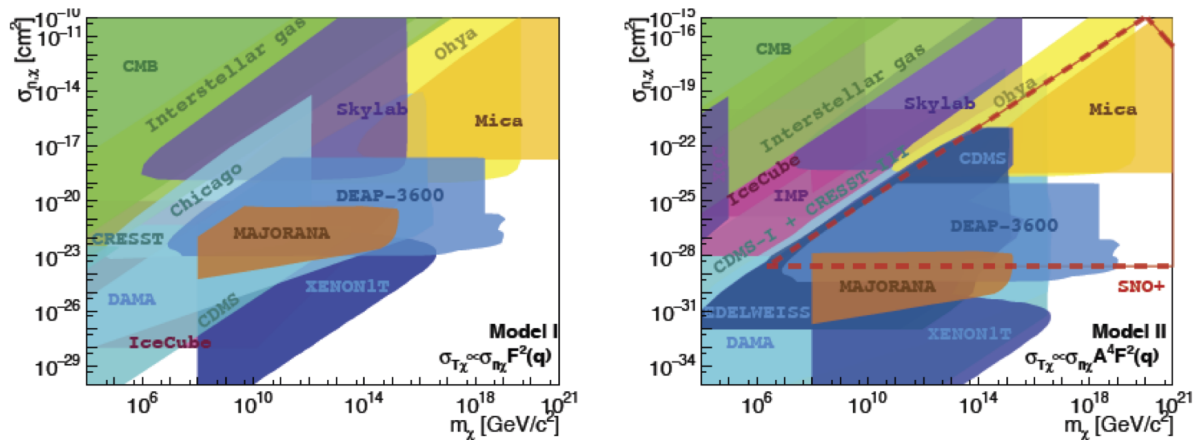
1200 sky as it orbits, whereas ground-based air and water Cherenkov telescopes are effective at energies above  
 1201  $\sim 100$  GeV. Most current Cherenkov telescopes are located in the Northern Hemisphere, but Southern  
 1202 Hemisphere locations are advantageous for observing both the Galactic Center and the central dense region  
 1203 of the Milky Way’s dark matter halo. For  $s$ -wave annihilation, Fermi-LAT has already strongly constrained  
 1204 the thermal freeze-out cross section for masses up to  $\sim 100$  GeV based on observations of the Milky Way’s  
 1205 dwarf spheroidal galaxies (e.g. [121, 122]), which are (as measured by stellar dynamics) dark matter rich and  
 1206 baryon poor, limiting the expected sources of background. These limits are expected to improve in the near  
 1207 future, as surveys such as Vera Rubin discover (about a factor of two) more dwarf spheroidal galaxies [123].

1208 In the near future, the SWGO [124] (water Cherenkov) and CTA [125] (air Cherenkov) telescopes, successors  
 1209 using similar technology to the successful HAWC and VERITAS Cherenkov telescopes but with larger  
 1210 installations and Southern Hemisphere sites, have the potential to probe the thermal freeze-out scenario  
 1211 up to 10s of TeV masses (depending on the annihilation channel) [126, 127], approaching the 100 TeV scale  
 1212 where we can begin to set unitarity-based limits on the capacity for freeze-out to generate the correct relic  
 1213 abundance. Further in the future, APT [128] is a concept for a space-based successor instrument to the  
 1214 Fermi-LAT (with a demonstrator suborbital mission scheduled for 2025), which aims to improve sensitivity  
 1215 at lower masses by an order of magnitude. Figure 5-18 shows the limits on WIMP annihilations into pairs  
 1216 of bottom quarks (for masses below  $\sim 100$  GeV) and  $W$  bosons (for larger masses) based on null searches  
 1217 by gamma-ray observatories. The beige regions indicate the current limits for each mass, whereas the green  
 1218 shaded region shows improvements expected in the near future based on planned Cherenkov observatories  
 1219 and including new populations of dwarf spheroidal galaxies expected to be discovered by Rubin / LSST. The  
 1220 blue shading indicates parameter space that could be probed by longer term investment into future large  
 1221 Cherenkov arrays and new space-based missions. The orange band indicates the benchmark cross section  
 1222 corresponding to the correct relic abundance from freeze-out. The gray “unitarity bound” region corresponds  
 1223 to the general mass range (the exact bound is model-dependent and depends on assumptions about long-  
 1224 range forces, compositeness of the dark matter, etc.) in which obtaining the correct dark matter abundance  
 1225 via freeze-out becomes inconsistent with unitarity in the early universe, under standard assumptions for the  
 1226 cosmological history.

1227 Cosmic-ray experiments (primarily AMS-02) currently set competitive constraints on heavy annihilating  
 1228 or decaying DM, with systematic uncertainties independent from gamma-ray probes. In the near future,  
 1229 GAPS [129] will provide the first dedicated search for low-energy anti-deuterons, which are expected to have  
 1230 only tiny astrophysical backgrounds and could serve as a very clean discovery channel, and HELIX [130]  
 1231 will provide new constraints on cosmic-ray propagation. Far-future proposals ALADInO [131, 132] &  
 1232 AMS-100 [133] involving superconducting magnetic spectrometers seek greatly improved sensitivity to anti-  
 1233 deuterons and anti-helium; GRAMS [70] would employ a liquid-argon TPC to search for both gamma  
 1234 rays and charged cosmic rays, and ADHD [56] is a proposal to pursue a novel delayed-annihilation signal  
 1235 from anti-nuclei in a helium detector. Fixed-target accelerator experiments including NA61/SHINE [134],  
 1236 ALICE [135], LHCb [136], and AMBER [137] can also provide complementary measurements to constrain  
 1237 cosmic-ray production and propagation, helping to reduce systematic uncertainties in both signals and  
 1238 backgrounds.

#### 1239 5.5.4 Ultra-Heavy Dark Matter ( $> 1$ PeV to $100 M_{\odot}$ )

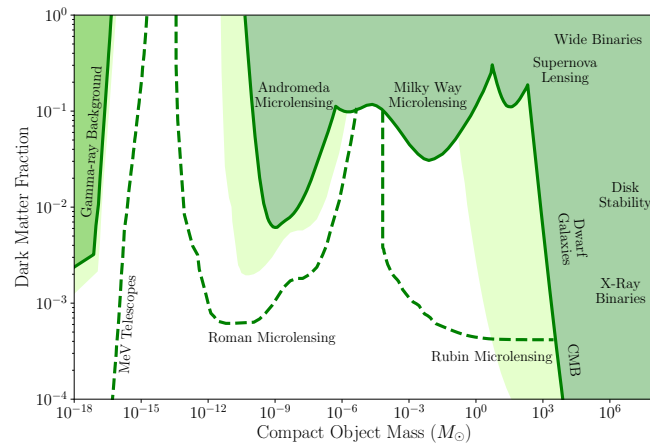
1240 Measurements of baryon acoustic oscillations and the cosmic microwave background set some of the strongest  
 1241 limits on the total baryonic fraction of matter in the universe. The sound waves of the primordial plasma  
 1242 have compressions due to gravitational attraction and the restoring force due to plasma pressure causes the  
 1243 following rarefactions. The higher amplitude in the odd-numbered power spectrum peaks indicate that most



**Figure 5-19.** Because the kinematics of nuclear recoil asymptotes to fixed recoil energy for large dark matter masses, direct detection experiments also have sensitivity to ultra heavy dark matter where sensitivity is primarily limited by the cross-sectional area of the detector.

1244 of the gravitating matter does not exhibit pressure when compressed and is hence not electromagnetically  
 1245 charged like protons or electrons. So even if dark matter took the form of non-luminous, cold baseballs (a  
 1246 model that is difficult to constrain observationally), these baseballs would still have to be made of some  
 1247 non-standard exotic particles.

1248 For masses beyond the unitarity limit, it becomes a challenge to produce dark matter via freeze-out. One  
 1249 interesting set of models produces dark matter in this mass range as composite blobs of more fundamental  
 1250 particles, held together by new dark forces. Cosmic observation of the CMB and interstellar gas set  
 1251 robust constraints at large cross sections, while observations of extreme astrophysical environments can  
 1252 also constrain the presence of UHDM. Both the searches for WIMPs described in the previous subsection  
 1253 as well as more specialized searches have sensitivity to this UHDM regime, depending on the strength  
 1254 and nature of its interactions with the SM (see Ref. [118] for a more in-depth discussion). For strongly  
 1255 interacting blobs, a typical event in a terrestrial experiment could involve multiple scattering, necessitating  
 1256 new analysis techniques and implying that the area of the detector is the most relevant factor rather than  
 1257 its volume. For indirect searches, high energy neutrinos and cosmic rays become especially important at  
 1258 high masses. Cascades of secondary particles from high-energy primaries often lead to observable indirect  
 1259 signatures at much lower energies than the DM mass, such that the gamma-ray telescopes described above  
 1260 have sensitivity to DM masses much higher than their target energy ranges. This situation represents  
 1261 an ongoing theoretical challenge, requiring new techniques for accurate predictions to take advantage of  
 1262 the opportunity to observe rich and complementary multi-wavelength and multi-messenger signals [118].  
 1263 For DM masses of order GUT-scale, however, the secondary photons *en route* to Earth travel unscathed.  
 1264 Accordingly, AugerPrime and next-generation cosmic ray observatories anchor unique DM indirect detection  
 1265 experiments which are free of astrophysical background: a clear detection of an extreme energy photon  
 1266 would be a momentous discovery [138]. Fig. 5-19 shows the current and projected experimental regions of  
 1267 ultra-heavy parameter space.



**Figure 5-20.** At even higher masses, some fraction of the dark matter could take the form of macroscopic, compact objects such as primordial black holes. Upcoming microlensing and gamma-ray searches will provide greater sensitivity in this regime (dashed lines).

#### 1268 5.5.4.1 Dark Matter Beyond the Planck Scale

1269 Dark matter at the Planck mass ( $\sim 10^{-5}$  g) and above is very difficult to detect directly due to its extremely  
 1270 low flux. Detection techniques rely on scattering to be mediated via long range forces so that effects can  
 1271 be seen in sparsely instrumented detectors with large collecting area. For example, if dark matter had a  
 1272 long range Yukawa force  $10^3$  times greater than gravity, then interactions of 10 kg mass dark matter with  
 1273 the 10 kg scale mirrors of the LIGO gravitational wave observatory could potentially be detected as the  
 1274 clumps traversed within the near field of the 4 km long interferometers [139]. Proposed larger observatories  
 1275 such as the 40 km long Cosmic Explorer or laser ranging between asteroids would provide a larger collection  
 1276 area [140].

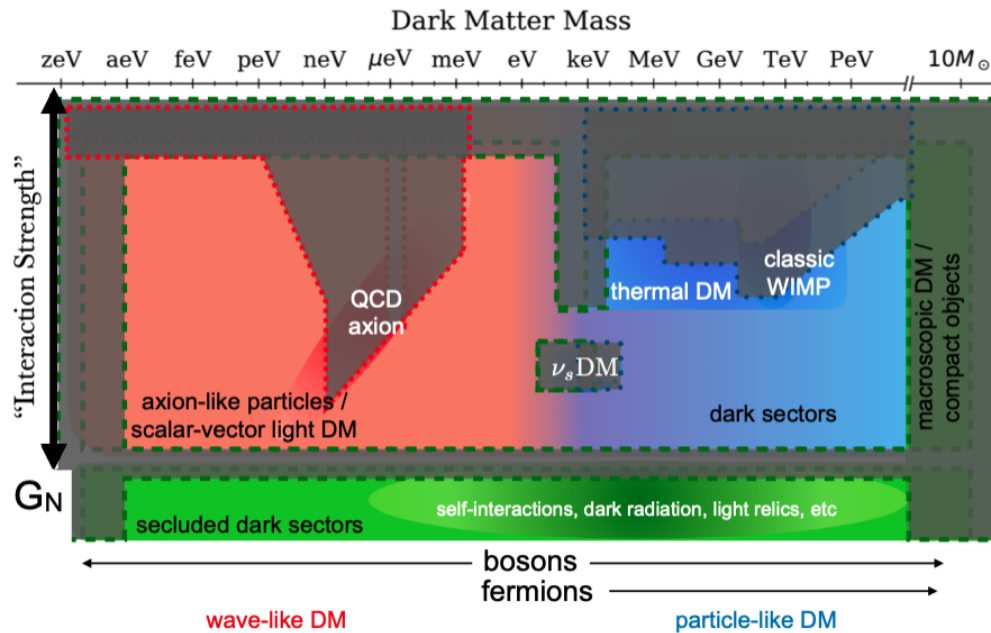
1277 Heavier compact constituents of dark matter including primordial black holes (PBHs), can be detected  
 1278 through cosmic probes (Figure 5-20). PBHs with masses of  $\sim 10^{-18} M_\odot$  ( $\sim 10^{15}$  g) should be evaporating  
 1279 today, creating short bursts of Hawking radiation in the GeV-TeV range that could be observed with  
 1280 space-based and ground-based gamma-ray telescopes, as well as neutrino observatories. For higher masses,  
 1281 gravitational micro-lensing provides the most sensitive probe of the abundances of PBHs. Near-future  
 1282 observatories such as Rubin LSST have the potential to radically increase the sensitivity of micro-lensing  
 1283 searches if scheduled optimally [141]. Sub-solar mass PBHs may be detected through microlensing with  
 1284 Rubin LSST and other observatories or through mergers with next-generation gravitational wave detectors  
 1285 including LISA and pulsar timing arrays. High redshift searches with optical or gravitational wave probes  
 1286 may also reveal the existence of primordial black holes.

#### 1287 5.5.5 Delve Deep, Search Wide

1288 The parameter space of dark matter models is vast, yet experiments which will provide significant inroads  
 1289 into this space are relatively inexpensive. The Cosmic Frontier plans to explore dark matter parameter  
 1290 space over the next decade through targeted searches to make deep progress on high-priority science targets,  
 1291 and an expanded portfolio of cosmic probes and small pathfinder experiments implementing new detector



1292 technologies to search wide and provide broad coverage of the range of possibilities. This strategy will cover  
 1293 large regions of model-space, and could easily lead to a transformational discovery within the next decade  
 1294 (see Figure 5-21).



1295 **Figure 5-21.** Cartoon (not to scale) of future dark matter sensitivity which could be obtained by  
 1296 implementing the Delve Deep, Search Wide strategy outlined in this report.

1295 The WIMP and the QCD axion remain highly motivated and clear experimental targets. Tremendous  
 1296 progress has been made in the past two decades in liquid noble detectors and a multi-national collabo-  
 1297 rative project would continue this momentum to push sensitivity down to the solar neutrino background.  
 1298 Concurrently, indirect searches with proposed TeV gamma-ray observatories would conclusively test the  
 1299 self-annihilation cross section which defines the WIMP model, up to the scale where unitarity bounds  
 1300 become relevant. In the axion field, a portfolio of complementary techniques including NMR, lumped  
 1301 element resonators, cavity resonators, quantum sensors, and novel scattering targets would cover most of  
 1302 the remaining parameter space of the QCD models that solve the strong-CP problem. Axion-like particles  
 1303 are also well-motivated with a parameter space that can be targeted over the next ten years and beyond.  
 1304 Similarly, a portfolio of new technologies including CCDs, novel condensed matter scattering targets, and  
 1305 ultra-low-noise readouts (specifically reducing the size and thus heat capacity of transition edge sensors  
 1306 to match them to the smaller phonon signal strengths expected in sub-eV band gap target materials) is  
 1307 being assembled to test a broad range of portal dark matter models. Pathfinder experiments using AMO  
 1308 techniques including atomic clocks, atom interferometers, and opto-mechanical oscillators are also being  
 1309 developed. Cosmic and indirect probes will complement and inform all of these searches, while broadly  
 1310 exploring parameter space and phenomena that are currently inaccessible to terrestrial experiments. With  
 1311 the exception of the medium-scale multi-national WIMP experiments (and of the large-scale cosmic survey  
 1312 experiments), all of the other terrestrial dark matter experiments would be best characterized as “small.”

1313 All combined, a coordinated U.S. experimental program to delve deep and search wide for dark matter might  
 1314 reasonably be expected to cost as much as a single large experiment. Such a comprehensive and coordinated  
 1315 program would be very reasonably justified by its broad science reach in addition to its broader impacts



1316 in technology development, contributing to the quantum ecosystem, and training the next generation of  
1317 experimental hardware specialists.

1318 Development of new experimental techniques for both low-mass portal dark matter, axions, and more general  
1319 bosonic waves has been slow to ramp up as most DOE/NSF dark matter funding has been invested in the  
1320 larger dark matter experiments as per the recommendations of the previous P5 report. Identifying this  
1321 gap in the U.S. dark matter strategy, private foundations including the Simons Foundation, the Heising-  
1322 Simons Foundation, and the Moore Foundation have begun to invest in university programs to fund dark  
1323 matter instrumentation R&D including pathfinder experiments at the \$10M level. Some of the technology  
1324 development and instrumentation work has also been enabled by leveraging various quantum programs, given  
1325 that ultraweak dark matter interactions are a natural target for new quantum sensor technology. Cross-  
1326 disciplinary engagement with neighboring scientific communities including atomic, molecular, and optical  
1327 physics, condensed matter physics, and quantum information science have been facilitated by DOE-OHEP's  
1328 QuantISED program and by the National Quantum Initiative Science Research Centers. This collaborative  
1329 research enables the direct transfer of novel detection concepts as well as mature detector technologies from  
1330 neighboring fields to HEP applications as well as the direct engagement of hardware experts who are able  
1331 offer fresh perspectives on the detection challenges.

1332 A positive development has been two DOE Basic Research Needs workshops, the first creating the Dark  
1333 Matter New Initiatives (DMNI) program to fund small dark matter experiments of a scale smaller than the  
1334 medium-size LZ and SuperCDMS experiments, and the second aimed at identifying instrumentation R&D  
1335 needs including quantum sensing for fundamental physics. The DMNI program has funded the technical  
1336 design studies of a limited portfolio of small experiments as recommended in the previous P5 report. These  
1337 include the axion experiments ADMX-Extended-Frequency-Range and DM-Radio, the 1 eV threshold Oscura  
1338 silicon CCD experiment, the sub-eV threshold TESSERACT program that utilizes novel condensed matter  
1339 targets, as well as beam dump experiments LDMX and CCM that probe portal dark matter models via light  
1340 mediators. While these experimental concepts have not yet proceeded to project status, the DMNI program  
1341 provides a prototype for implementing a broad and coordinated dark matter program with a portfolio of  
1342 complementary experiments to target different regions of unprobed parameter space.

1343 Similar support could significantly strengthen cosmic and indirect probes of dark matter. These approaches  
1344 have the potential to yield large scientific dividends by searching wide ranges of dark matter parameter  
1345 space. For example, Rubin LSST has enormous potential to discover new physics beyond the prevailing cold  
1346 dark matter paradigm [20, 38]. Rubin LSST can measure the distribution of dark matter on unprecedentedly  
1347 small scales, thereby probing microscopic properties of dark matter, including thermal particle mass, self-  
1348 interactions, interactions with radiation, and quantum wave features. Microlensing measurements will  
1349 directly probe primordial black holes as a component of dark matter. The planned CMB-S4 [10] and Spec-  
1350 S5 [22] projects will extend access to rich dark matter particle physics. For instance, detection of additional  
1351 relativistic degrees of freedom by CMB-S4 would imply the existence of a dark sector. However, building the  
1352 infrastructure to perform dark matter analyses with these experimental facilities requires devoted support  
1353 for experimental collaborations, theorist, and large numerical simulations (which are crucial to distinguish  
1354 novel dark matter physics from baryonic astrophysics). Rubin LSST and other cosmology facilities should  
1355 be explicitly identified as a dark matter facilities to enable what promises to be an exciting decade of dark  
1356 matter research.

## 1357 5.6 Neutrinos

1358 Cosmic probes can provide crucial information about the neutrino sector and exhibit a high degree of  
1359 complementarity to terrestrial probes [142]. Cosmology is sensitive to the number and masses of the neutrinos

1360 through their impact on the evolution of the Universe, providing information that is currently inaccessible  
1361 through other means.

### 1362 5.6.1 Cosmic Measurements of Neutrino Masses

1363 The current cosmological limits on the sum of the neutrino masses is  $\sum m_\nu < 0.12$  eV as obtained using  
1364 cosmic microwave background, baryon acoustic oscillations, supernova Ia, and large scale structure measure-  
1365 ments. Additional information on impact of cosmological neutrinos on the growth of structure both will be  
1366 obtained from high redshift galaxy surveys. CMB-S4 will reach sensitivity  $\sigma(\sum m_\nu) < 0.02$  eV and if the data  
1367 indicate a mass sum  $< 0.1$  eV, this would rule out the inverted neutrino mass hierarchy which predicts higher  
1368 mass sum. A Stage V Spectroscopic Facility would reach similar constraints from an independent experiment  
1369 at  $2 < z < 5$ . Further gains in accuracy on  $\sigma(m_\nu)$  beyond what CMB-S4 and Spec-S5 each will achieve will  
1370 be limited by the current  $\tau$  prior, unless that parameter is constrained better. Proposals to achieve such  
1371 improvements are described in the CF4 report [4]. Higher statistics observations of the matter distribution  
1372 at high redshifts could be measured by line intensity mapping surveys (in 21-cm and mm-wave) and other  
1373 probes, such as future high resolution CMB imaging, and would reach resolution  $\sigma(\sum m_\nu) < 0.01$  eV.

### 1374 5.6.2 New Opportunity: High-energy neutrinos

1375 The discovery of TeV-PeV astrophysical neutrinos opens up unique opportunities to probe the neutrino sector  
1376 at energy scales not accessible with laboratory neutrino beams. High-energy neutrinos from IceCube were  
1377 used to discover the Glashow resonance and high-energy neutrino-nucleon cross section. Future statistics  
1378 will allow further tests of the Standard Model neutrino physics between 1 TeV and 10 PeV. In addition,  
1379 ultrahigh energy neutrinos (UHE;  $\gtrsim 100$  PeV neutrinos) have been long-predicted but remain undetected.  
1380 They provide a path to probe weak-scale physics at center-of-mass energies above 50 TeV. Next-generation  
1381 UHE neutrino experiments have the potential to detect these neutrinos and push the forefront of neutrino  
1382 physics [143].

1383 Observations of neutrino flavors and neutrino-antineutrino ratios may probe BSM physics [144]. The search  
1384 for BSM neutrino interactions with other neutrinos or with dark matter may shed light on the UV theory  
1385 of neutrinos and the origins of their low mass. Such new interactions could be observed for example via  
1386 spectral distortions in the cosmic neutrino flux in the PeV-EeV range as they scatter on the cosmic neutrino  
1387 background. Deviations from equal admixtures of the 3 neutrino species might occur at various energy  
1388 thresholds when new interactions turn on. Neutrino production measurements from CERN's Forward Physics  
1389 Facility will provide important information to calibrate the atmospheric neutrino background produced by  
1390 cosmic rays scattering on nuclei in the atmosphere and thus allow cleaner detection of cosmic neutrinos [145,  
1391 146]. Finally, at GeV energies, the upcoming IceCube Upgrade and other future cosmic neutrino experiments  
1392 will provide neutrino oscillation sensitivity complementary to long baseline experiments.

## 1393 5.7 Exploring the Unknown: New Particles, New Fields, New 1394 Principles of Nature

### 1395 5.7.1 Dark Radiation

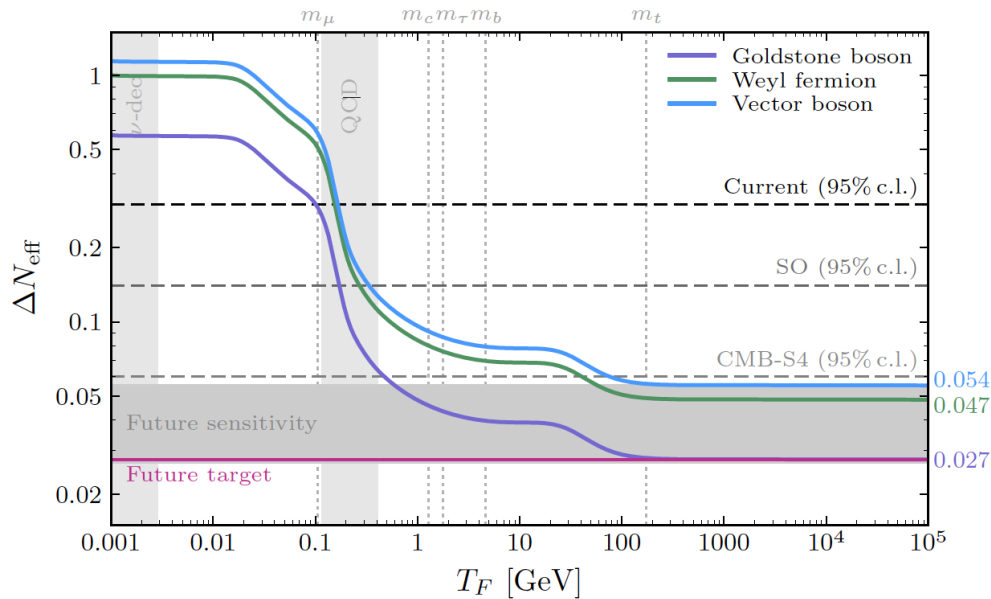
1396 Dark radiation can be observed in various cosmological epochs while it is still relativistic, before its kinetic  
1397 energy redshifts away. Within the Standard Model (SM), neutrinos form an important component of  
1398 dark radiation, and constraints from their impact on the cosmic evolution at the epochs of Big Bang  
1399 Nucleosynthesis and the CMB provide powerful constraints on the number of neutrino species, typically  
1400 reported as  $N_{\text{eff}}$ , the “effective number of neutrinos”. In addition, determinations of the dark radiation  
1401 content of the Universe provide unique opportunities to search for particles produced in the early universe,  
1402 even when such particles are extremely weakly interacting with the SM. A few well-motivated examples  
1403 include light particles invoked by models that aim to explain the physics of a dark sector, address the strong  
1404 CP problem, solve the weak hierarchy problem, account for short baseline neutrino anomalies, and/or models  
1405 of warm inflation.

1406 Figure 5-22 shows the current measurement of  $\Delta N_{\text{eff}}$ , and future prospects from operation of the Simons  
1407 Observatory and CMB-S4. The projections indicate the contribution to  $\Delta N_{\text{eff}}$  from a single species of one  
1408 of three different types of relic particles (Goldstone or vector bosons and Weyl fermions) for a particle which  
1409 was initially in chemical equilibrium with the SM plasma, but whose interactions decoupled at freeze out  
1410 temperature  $T_F$ . CMB-S4, by reaching a precision of  $\Delta N_{\text{eff}} \lesssim 0.05$  will be able to rule out or establish  
1411 the existence of such weakly-interacting particles which have frozen out back to the temperature of the  
1412 QCD phase transition ( $\sim 0.5$  GeV). A Stage V Spectroscopic Facility would reach similar constraints from  
1413 an independent experiment at  $2 < z < 5$  which could be combined to probe higher precision. A future  
1414 measurement reaching a precision of 0.027 would be sensitive to such particles freezing out all the way back  
1415 to the electroweak scale.

1416 Quintessence models promote the dark energy from a cosmological constant to a dynamically evolving field  
1417 which is slowly rolling in a scalar potential, triggering a new epoch of cosmic inflation. A natural question is  
1418 then whether this quintessence field is slowly rolling in a flat potential, or if it is rolling in a steeper potential  
1419 but slowed down by dynamical friction. In this warm dark energy scenario, the ongoing interactions during  
1420 the slow roll in the present day universe would populate the universe with dark radiation in the form of  
1421 relativistic, low mass particles with temperature and energy density typically characterized by the milli-eV  
1422 scale (10 K) of the dark energy density [147]. This late injection of entropy is not constrained by bounds on  
1423 CMB distortion at earlier times. Searches for this 10 K bath of exotic particles would nicely complement the  
1424 cosmic surveys studies of the evolution of cosmic acceleration. Wave dark matter detectors and quantum  
1425 sensors could provide the required low energy threshold of  $10^{-3}$  eV to see absorption of this dark radiation.

### 1426 5.7.2 The Highest Energy Particles

1427 Cosmic particles above 100 TeV provide a unique window into into fundamental particle physics at energy  
1428 scales beyond those reachable by terrestrial accelerators. In addition to providing indirect probes of dark  
1429 matter in various annihilation or decay channels and the measurements of the neutrino sector described  
1430 elsewhere in this report, observatories of cosmic rays, neutrinos, gamma rays, and gravitational waves  
1431 together constitute a rich, multi-messenger program to explore the unknown through the study of the highest  
1432 energy particle astrophysical phenomena in the universe.



**Figure 5-22.** The presence of additional relativistic BSM particles beyond the three known neutrino species increases the Hubble expansion rate and thus the power spectrum of acoustic oscillations in the primordial plasma. Measurements of the CMB can be used to discern the presence or absence of these exotic degrees of freedom, as parameterized as  $N_{\text{eff}}$ , the effective number of relativistic species active during this epoch. Some representative models are shown in this plot along with the current constraints on additional particles and projections for the Simons Observatory and CMB-S4.

1433 In the area of hadronic shower physics, a long standing puzzle is the mysterious excess of muons in UHE  
 1434 cosmic ray ( $> 10^{18}$  eV cosmic-ray) air showers beyond what is expected from extrapolations of hadronization  
 1435 models calibrated with LHC and other collider data [148]. Ongoing upgrades of the Pierre Auger Observatory  
 1436 and IceCube will provide greater resolution of the electromagnetic and muonic flux in air showers as well as  
 1437 measurements of the depth of shower maximum which provides complementary information on the initial  
 1438 particle species. The proposed Forward Physics Facility at the LHC [145, 146] would also provide muon and  
 1439 electron neutrino shower data which could be used to calibrate kaon production in high energy hadronic  
 1440 showers. Higher than expected kaon production rates would enhance the muon production by reducing the  
 1441 production of neutral pions which siphon energy into the electromagnetic portion of the shower. Whether  
 1442 the puzzle is resolved by a better understanding of the fragmentation process or potentially by BSM physics  
 1443 at high center-of-mass energies, these data will better inform the modeling and design of future high energy  
 1444 colliders. The high boosts and large, cosmological propagation scales of UHE cosmic particles allow powerful  
 1445 tests of Lorentz and CPT symmetries by searches for spectral distortions in the flux of UHE cosmic rays,  
 1446 very-high-energy gamma-rays (0.1 – 100 TeV photons), UHE gamma-rays ( $> 100$  TeV photons), and cosmic  
 1447 neutrinos [149].

1448 High energy gamma rays and neutrinos additionally offer probes of new physics including the decays of  
 1449 super-heavy relics left behind from the Big Bang, cosmic strings, and axion-photon conversion in large-  
 1450 scale magnetic fields. Next-generation gamma-ray telescopes such as the Southern Wide-field Gamma-ray  
 1451 Observatory (SWGGO), the Cherenkov Telescope Array (CTA), and the All-sky Medium-Energy Gamma-  
 1452 ray Observatory (AMEGO) will open up the access to new sky regions and energy ranges, and advance  
 1453 fundamental physics studies with significantly improved sensitivities. Next-generation UHE neutrino ( $>$   
 1454 100 PeV neutrino) observatories led by the U.S. community such as the IceCube-Gen2, The Radio Neutrino

1455 Observatory in Greenland (RNO-G), Probe of Extreme Multi-Messenger Astrophysics (POEMMA), Beam  
1456 forming Elevated Array for COsmic Neutrinos (BEACON), the Payload for Ultrahigh Energy Observations  
1457 (PUEO), and the Extreme Neutrino Observatory (Trinity), have the potential to discover UHE neutrinos  
1458 and use them to probe fundamental physics.

1459 The concept of multi-messenger astronomy is finally being realized with co-detection of gamma rays and  
1460 gravitational waves in a binary neutron star merger and the co-detection of neutrinos and gamma rays in a  
1461 blazar flare [150]. Studies of these violently energetic astrophysical events may reveal new physics including  
1462 the nature of quark matter in neutron stars, the possible accumulation or scattering of dark matter on  
1463 these objects, and the production of exotic particles in large astrophysical large magnetic fields. Various  
1464 tests of Einstein gravity may lead to discoveries of relevance to cosmology including modifications to the  
1465 understanding of cosmic acceleration and the distribution of matter. As remarked in the dark matter section  
1466 of this report, while a 20-year strategic plan is being developed for most cosmic particle channels across  
1467 broad ranges in energy, an ongoing concern is the lack of coverage for MeV-GeV gamma rays beyond that  
1468 provided by the aging Fermi telescope.

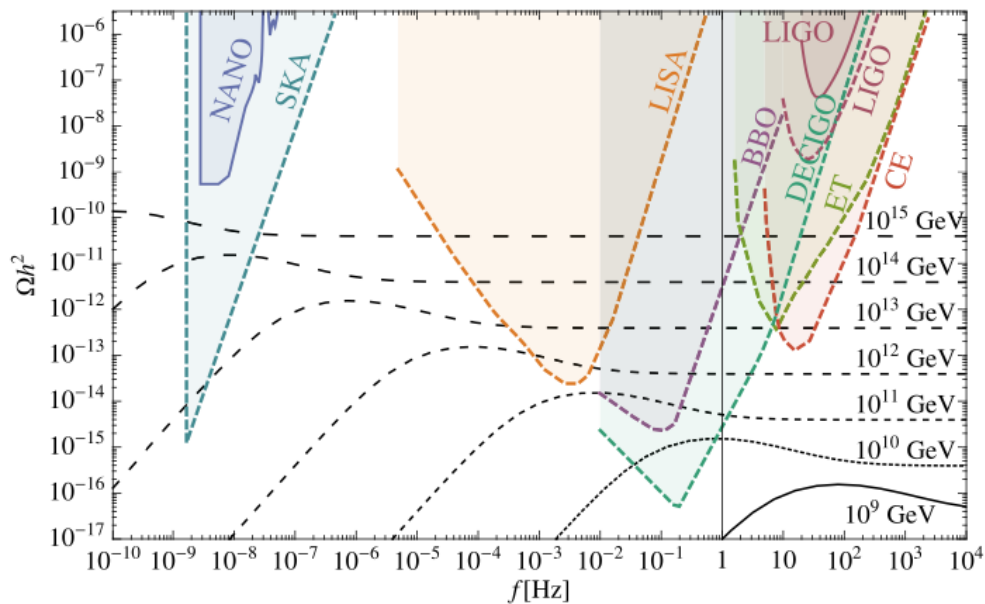
### 1469 5.7.3 Probes of Fundamental Physics with Gravitational Waves

1470 One of the most significant scientific developments of the past decade is the emergence of gravitational  
1471 wave detection as a powerful new tool to study fundamental physics. In addition to the extensive menu of  
1472 gravitational physics that will be studied by current and future gravitational wave observatories, the new  
1473 GW probes will also provide unique and independent information in a number of areas of interest to the high  
1474 energy physics community. For example, neutron star binary mergers will probe the neutron star equation  
1475 of state and may also be sensitive to accumulations of various types of dark matter within the neutron stars.  
1476 Once the neutron star physics is understood, the binary mergers may provide a new standard candle for  
1477 measurements of cosmic acceleration with systematics independent of those of the currently used optical  
1478 probes. Alternatively, the neutron star merger's gravitational wave luminosity distance may be combined  
1479 with redshift information from the host galaxy to provide new measurements of the Hubble parameter at  
1480 large redshifts.

1481 Gravitational wave observatories will also provide unique probes of early universe phase transitions which  
1482 may have happened at times prior to big bang nucleosynthesis – a period about which not much information  
1483 is currently available. These phase transitions may produce a stochastic background of gravitational waves  
1484 from boiling of the vacuum during strongly first order phase transitions or from the oscillations and decay of  
1485 cosmic strings and other topological defects which may form during these phase transitions. The resulting  
1486 gravitational wave power spectrum will provide information about these new fundamental energy scales  
1487 which, once discovered will provide concrete new information for BSM model building.

1488 The signal will be visible even through the cosmic microwave background which obscures photon probes of  
1489 earlier cosmological epochs. Even phase transitions which might happen in secluded dark sectors which have  
1490 no BSM couplings to standard model particles can still be probed due to the universal coupling of gravity to  
1491 all matter and energy. As example, figure 5-23 shows how a broad program of new gravitational observatories  
1492 targeting different frequency ranges can probe deeply into the predicted power spectrum of phase transitions  
1493 at energy scales from  $10^{10} - 10^{15}$  GeV for various models of thermal leptogenesis. The observatories range  
1494 from pulsar timing arrays at the lowest frequencies to space-based experiments at sub-Hz frequencies to next  
1495 generation ground-based interferometers at the highest frequencies.

1496 New physics of early universe phase transitions may even be observed with the current generation of  
1497 gravitational wave observatories. As a case study, LIGO may already be sensitive to the stochastic power



**Figure 5-23.** Future gravitational wave observatories will be sensitive to the stochastic spectrum of GWs created by cosmic string decay or by boiling of the vacuum in first order phase transitions. This plot shows an example of the high energy symmetry breaking scales that can be probed in a particular model of thermal leptogenesis. In other models, the predicted spectrum may be peaked at certain frequencies, in which case observatories optimized for specific frequency ranges may have greater sensitivity. Figure shown in the Seattle meeting by H. Murayama and reproduced from reference [151].

1498 spectrum from Peccei-Quinn phase transition in QCD axion models if a particularly violent phase transition  
 1499 occurs at relatively low energies around  $10^8$  GeV where the red-shifted frequency spectrum of the GW  
 1500 emission becomes matched to the LIGO band. If the energy scale of the phase transition can be determined  
 1501 in this way, then the axion model gives a firm prediction for the mass of the dark matter axion and will narrow  
 1502 the search window for the direct detection experiments from 10 decades in mass down to a single decade. A  
 1503 confirmation of an dark matter signal would then provide corroborating evidence for the axion solution to  
 1504 the strong-CP problem. Alternatively, if the dark matter signal is discovered first in the post-inflationary  
 1505 scenario with mass between  $10^{-5} - 10^{-2}$  eV, then this provides a firm target for the sensitivity and frequency  
 1506 range needed for future gravitational wave observatories to observe the phase transition. Similar strategies  
 1507 may be employed to search for new fundamental energy scales and test see-saw models in other contexts.

1508 The collection of currently operating gravitational wave observatories include the LIGO facilities in Liv-  
 1509 ington, Louisiana and Hanford, Washington, along with the Virgo facility in Italy and the recently con-  
 1510 structed, underground KAGRA observatory in Japan. LIGO-India will also be online at the end of the  
 1511 decade, and the space-based LISA observatory is planned for the next decade. Pulsar timing arrays operating  
 1512 at lower frequencies include NANOGrav in North America, the Parkes Pulsar Timing Array in Australia,  
 1513 and the European Pulsar Timing Array, while plans are being made for a future Square Kilometer Array.  
 1514 Near term plans for the U.S. gravitational wave community include the LIGO Voyager upgrade in which  
 1515 the existing fused silica mirrors may be replaced with crystalline silicon mirrors which can be more easily  
 1516 thermalized to avoid geometric distortions from heating, and the Cosmic Explorer proposal which would  
 1517 increase the interferometer arm length and hence antenna size from 4 km to 40 km. The AMO community  
 1518 is also developing long baseline atom interferometry as a potential new technique targeting lower frequency  
 1519 gravitational waves as well as oscillatory forces from dark matter in the 1 Hz band. The MAGIS-100

1520 pathfinder experiment will utilize a 100 m vertical beamline access shaft at Fermilab with possible future  
1521 expansion to a 2 km drop at SURF and eventually a space mission to avoid terrestrial Newtonian noise.

1522 Participation of the HEP community may be key to the success of these ambitious projects and to ensure  
1523 that optimizations for HEP science can be integrated into the designs, data pipelines, and operations plans.  
1524 As an example, searches for stochastic gravitational wave power can be performed within a narrow band  
1525 to reject noise as opposed to measurements of time domain waveforms which intrinsically require broader  
1526 bandwidth template searches. Similarly, searches for kg-mass dark matter passing within an antenna length  
1527 of the interferometers will need dedicated data pipelines and analyses. More direct engagement by national  
1528 laboratories may also be critical for the construction of future large scale GWO projects.



---

---

# Bibliography

- 1529 [1] J. Cooley et al., *Report of the Topical Group on Particle Dark Matter for Snowmass 2021*, [2209.07426](#).
- 1530 [2] J. Jaeckel, G. Rybka and L. Winslow, *Report of the Topical Group on Wave Dark Matter for Snowmass*  
1531 *2021*, [2209.08125](#).
- 1532 [3] A. Drlica-Wagner et al., *Report of the Topical Group on Cosmic Probes of Dark Matter for Snowmass*  
1533 *2021*, [2209.08215](#).
- 1534 [4] J. Annis, J.A. Newman and A. Slosar, *Snowmass2021 Cosmic Frontier: Report of the CF04 Topical*  
1535 *Group on Dark Energy and Cosmic Acceleration in the Modern Universe*, [2209.08049](#).
- 1536 [5] C.L. Chang et al., *Report of the Topical Group on Cosmic Frontier 5 Dark Energy and Cosmic*  
1537 *Acceleration: Cosmic Dawn and Before for Snowmass 2021*, in *2022 Snowmass Summer Study*, 9,  
1538 2022 [[2209.08265](#)].
- 1539 [6] B. Flaugher et al., *Report of the Topical Group on Dark Energy and Cosmic Acceleration:*  
1540 *Complementarity of Probes and New Facilities for Snowmass 2021*, in *2022 Snowmass Summer*  
1541 *Study*, 9, 2022 [[2209.08654](#)].
- 1542 [7] R.X. Adhikari et al., *Report of the Topical Group on Cosmic Probes of Fundamental Physics for for*  
1543 *Snowmass 2021*, [2209.11726](#).
- 1544 [8] A. Albrecht et al., *Report of the Dark Energy Task Force*, [astro-ph/0609591](#).
- 1545 [9] E. Abdalla et al., *Cosmology intertwined: A review of the particle physics, astrophysics, and cosmology*  
1546 *associated with the cosmological tensions and anomalies*, *JHEAp* **34** (2022) 49 [[2203.06142](#)].
- 1547 [10] CMB-S4 collaboration, *Snowmass 2021 CMB-S4 White Paper*, [2203.08024](#).
- 1548 [11] C.L. Chang et al., *Snowmass2021 Cosmic Frontier: Cosmic Microwave Background Measurements*  
1549 *White Paper*, [2203.07638](#).
- 1550 [12] G. Bertone, D. Hooper and J. Silk, *Particle dark matter: Evidence, candidates and constraints*, *Phys.*  
1551 *Rept.* **405** (2005) 279 [[hep-ph/0404175](#)].
- 1552 [13] G. Bertone and T. Tait, M. P., *A new era in the search for dark matter*, *Nature* **562** (2018) 51  
1553 [[1810.01668](#)].
- 1554 [14] J.L. Feng et al., *Planning the Future of U.S. Particle Physics (Snowmass 2013): Chapter 4: Cosmic*  
1555 *Frontier*, in *Community Summer Study 2013: Snowmass on the Mississippi*, 1, 2014 [[1401.6085](#)].
- 1556 [15] M. Battaglieri et al., *US Cosmic Visions: New Ideas in Dark Matter 2017: Community Report*, in  
1557 *U.S. Cosmic Visions: New Ideas in Dark Matter*, 7, 2017 [[1707.04591](#)].
- 1558 [16] K.R. Dienes and B. Thomas, *More is Different: Non-Minimal Dark Sectors and their Implications for*  
1559 *Particle Physics, Astrophysics, and Cosmology - 13 Take-Away Lessons for Snowmass 2021*, in *2022*  
1560 *Snowmass Summer Study*, 3, 2022 [[2203.17258](#)].

- 1561 [17] A. Banerjee et al., *Snowmass2021 Cosmic Frontier White Paper: Cosmological Simulations for Dark*  
1562 *Matter Physics*, in *2022 Snowmass Summer Study*, 3, 2022 [[2203.07049](#)].
- 1563 [18] S. Chakrabarti et al., *Snowmass2021 Cosmic Frontier White Paper: Observational Facilities to Study*  
1564 *Dark Matter*, in *2022 Snowmass Summer Study*, 3, 2022 [[2203.06200](#)].
- 1565 [19] M. Valluri et al., *Snowmass2021 Cosmic Frontier White Paper: Prospects for obtaining Dark Matter*  
1566 *Constraints with DESI*, [2203.07491](#).
- 1567 [20] Y.-Y. Mao et al., *Snowmass2021: Vera C. Rubin Observatory as a Flagship Dark Matter Experiment*,  
1568 [2203.07252](#).
- 1569 [21] C. Dvorkin et al., *Dark Matter Physics from the CMB-S4 Experiment*, in *2022 Snowmass Summer*  
1570 *Study*, 3, 2022 [[2203.07064](#)].
- 1571 [22] DESI collaboration, *A Spectroscopic Road Map for Cosmic Frontier: DESI, DESI-II, Stage-5*,  
1572 [2209.03585](#).
- 1573 [23] PLANCK collaboration, *Planck 2018 results. VI. Cosmological parameters*, *Astron. Astrophys.* **641**  
1574 (2020) A6 [[1807.06209](#)].
- 1575 [24] B.W. Lee and S. Weinberg, *Cosmological Lower Bound on Heavy Neutrino Masses*, *Phys. Rev. Lett.*  
1576 **39** (1977) 165.
- 1577 [25] J.L. Feng and J. Kumar, *The WIMPlless Miracle: Dark-Matter Particles without Weak-Scale Masses*  
1578 *or Weak Interactions*, *Phys. Rev. Lett.* **101** (2008) 231301 [[0803.4196](#)].
- 1579 [26] G. Gelmini, P. Gondolo, A. Soldatenko and C.E. Yaguna, *The Effect of a late decaying scalar on the*  
1580 *neutralino relic density*, *Phys. Rev. D* **74** (2006) 083514 [[hep-ph/0605016](#)].
- 1581 [27] G.B. Gelmini and P. Gondolo, *Neutralino with the right cold dark matter abundance in (almost) any*  
1582 *supersymmetric model*, *Phys. Rev. D* **74** (2006) 023510 [[hep-ph/0602230](#)].
- 1583 [28] S. Hamdan and J. Unwin, *Dark Matter Freeze-out During Matter Domination*, *Mod. Phys. Lett. A* **33**  
1584 (2018) 1850181 [[1710.03758](#)].
- 1585 [29] B. Barman, P. Ghosh, F.S. Queiroz and A.K. Saha, *Scalar multiplet dark matter in a fast expanding*  
1586 *Universe: Resurrection of the desert region*, *Phys. Rev. D* **104** (2021) 015040 [[2101.10175](#)].
- 1587 [30] D. Berger, S. Ipek, T.M.P. Tait and M. Waterbury, *Dark Matter Freeze Out during an Early*  
1588 *Cosmological Period of QCD Confinement*, *JHEP* **07** (2020) 192 [[2004.06727](#)].
- 1589 [31] K.R. Dienes, L. Heurtier, F. Huang, D. Kim, T.M.P. Tait and B. Thomas, *Stasis in an expanding*  
1590 *universe: A recipe for stable mixed-component cosmological eras*, *Phys. Rev. D* **105** (2022) 023530  
1591 [[2111.04753](#)].
- 1592 [32] J.N. Howard, S. Ipek, T.M.P. Tait and J. Turner, *Dark matter freeze-out during  $SU(2)_L$  confinement*,  
1593 *JHEP* **02** (2022) 047 [[2112.09152](#)].
- 1594 [33] N. Craig, *Naturalness: A Snowmass White Paper*, in *2022 Snowmass Summer Study*, 5, 2022  
1595 [[2205.05708](#)].
- 1596 [34] C.B. Adams et al., *Axion Dark Matter*, in *2022 Snowmass Summer Study*, 3, 2022 [[2203.14923](#)].
- 1597 [35] S. Tulin and H.-B. Yu, *Dark Matter Self-interactions and Small Scale Structure*, *Phys. Rept.* **730**  
1598 (2018) 1 [[1705.02358](#)].

- 1599 [36] S. Adhikari et al., *Astrophysical Tests of Dark Matter Self-Interactions*, [2207.10638](#).
- 1600 [37] J.S. Bullock and M. Boylan-Kolchin, *Small-Scale Challenges to the  $\Lambda$ CDM Paradigm*, *Ann. Rev. Astron. Astrophys.* **55** (2017) 343 [[1707.04256](#)].
- 1601
- 1602 [38] LSST DARK MATTER GROUP collaboration, *Probing the Fundamental Nature of Dark Matter with the Large Synoptic Survey Telescope*, [1902.01055](#).
- 1603
- 1604 [39] K. Bechtol et al., *Snowmass2021 Cosmic Frontier White Paper: Dark Matter Physics from Halo Measurements*, in *2022 Snowmass Summer Study*, 3, 2022 [[2203.07354](#)].
- 1605
- 1606 [40] M. Baryakhtar et al., *Dark Matter In Extreme Astrophysical Environments*, in *2022 Snowmass Summer Study*, 3, 2022 [[2203.07984](#)].
- 1607
- 1608 [41] P. Carena, T. Fischer, M. Giannotti, G. Guo, G. Martínez-Pinedo and A. Mirizzi, *Improved axion emissivity from a supernova via nucleon-nucleon bremsstrahlung*, *JCAP* **10** (2019) 016 [[1906.11844](#)].
- 1609
- 1610 [42] M. Buschmann, C. Dessert, J.W. Foster, A.J. Long and B.R. Safdi, *Upper Limit on the QCD Axion Mass from Isolated Neutron Star Cooling*, *Phys. Rev. Lett.* **128** (2022) 091102 [[2111.09892](#)].
- 1611
- 1612 [43] A. Ayala, I. Domínguez, M. Giannotti, A. Mirizzi and O. Straniero, *Revisiting the bound on axion-photon coupling from Globular Clusters*, *Phys. Rev. Lett.* **113** (2014) 191302 [[1406.6053](#)].
- 1613
- 1614 [44] M.J. Dolan, F.J. Hiskens and R.R. Volkas, *Advancing Globular Cluster Constraints on the Axion-Photon Coupling*, [2207.03102](#).
- 1615
- 1616 [45] F. Capozzi and G. Raffelt, *Axion and neutrino bounds improved with new calibrations of the tip of the red-giant branch using geometric distance determinations*, *Phys. Rev. D* **102** (2020) 083007 [[2007.03694](#)].
- 1617
- 1618
- 1619 [46] R. Brito, V. Cardoso and P. Pani, *Superradiance: New Frontiers in Black Hole Physics*, *Lect. Notes Phys.* **906** (2015) pp.1 [[1501.06570](#)].
- 1620
- 1621 [47] V. Gluscevic and K.K. Boddy, *Constraints on Scattering of keV–TeV Dark Matter with Protons in the Early Universe*, *Phys. Rev. Lett.* **121** (2018) 081301 [[1712.07133](#)].
- 1622
- 1623 [48] K.K. Boddy, V. Gluscevic, V. Poulin, E.D. Kovetz, M. Kamionkowski and R. Barkana, *Critical assessment of CMB limits on dark matter-baryon scattering: New treatment of the relative bulk velocity*, *Phys. Rev. D* **98** (2018) 123506 [[1808.00001](#)].
- 1624
- 1625
- 1626 [49] E.O. Nadler, V. Gluscevic, K.K. Boddy and R.H. Wechsler, *Constraints on Dark Matter Microphysics from the Milky Way Satellite Population*, *Astrophys. J. Lett.* **878** (2019) 32 [[1904.10000](#)].
- 1627
- 1628 [50] DES collaboration, *Milky Way Satellite Census. III. Constraints on Dark Matter Properties from Observations of Milky Way Satellite Galaxies*, *Phys. Rev. Lett.* **126** (2021) 091101 [[2008.00022](#)].
- 1629
- 1630 [51] K.K. Rogers, C. Dvorkin and H.V. Peiris, *Limits on the Light Dark Matter–Proton Cross Section from Cosmic Large-Scale Structure*, *Phys. Rev. Lett.* **128** (2022) 171301 [[2111.10386](#)].
- 1631
- 1632 [52] D. Carney et al., *Mechanical Quantum Sensing in the Search for Dark Matter*, *Quantum Sci. Technol.* **6** (2021) 024002 [[2008.06074](#)].
- 1633
- 1634 [53] F. Monteiro, G. Afek, D. Carney, G. Krnjaic, J. Wang and D.C. Moore, *Search for composite dark matter with optically levitated sensors*, *Phys. Rev. Lett.* **125** (2020) 181102 [[2007.12067](#)].
- 1635

- 1636 [54] G. Afek, D. Carney and D.C. Moore, *Coherent Scattering of Low Mass Dark Matter from Optically*  
1637 *Trapped Sensors*, *Phys. Rev. Lett.* **128** (2022) 101301 [[2111.03597](#)].
- 1638 [55] M. Cirelli, N. Fornengo, B.J. Kavanagh and E. Pinetti, *Integral X-ray constraints on sub-GeV Dark*  
1639 *Matter*, *Physical Review D* **103** (2021) 063022 [[2007.11493](#)].
- 1640 [56] T. Aramaki et al., *Snowmass2021 Cosmic Frontier: The landscape of cosmic-ray and high-energy*  
1641 *photon probes of particle dark matter*, [2203.06894](#).
- 1642 [57] M. Tashiro, H. Maejima, K. Toda, R. Kelley, L. Reichenthal, L. Hartz et al., *Status of x-ray*  
1643 *imaging and spectroscopy mission (XRISM)*, in *Society of Photo-Optical Instrumentation Engineers*  
1644 *(SPIE) Conference Series*, vol. 11444 of *Society of Photo-Optical Instrumentation Engineers (SPIE)*  
1645 *Conference Series*, p. 1144422, Dec., 2020, DOI.
- 1646 [58] XRISM Science Team, *Science with the X-ray Imaging and Spectroscopy Mission (XRISM)*, *arXiv*  
1647 *e-prints* (2020) arXiv:2003.04962 [[2003.04962](#)].
- 1648 [59] J.S. Adams, A.J. Anderson, R. Baker, S.R. Bandler, N. Bastidon, D. Castro et al., *Micro-X Sounding*  
1649 *Rocket: Transitioning from First Flight to a Dark Matter Configuration*, *Journal of Low Temperature*  
1650 *Physics* **199** (2020) 1072 [[1908.09010](#)].
- 1651 [60] K. Nandra, D. Barret, X. Barcons, A. Fabian, J.-W. den Herder, L. Piro et al., *The Hot and Energetic*  
1652 *Universe: A White Paper presenting the science theme motivating the Athena+ mission*, *arXiv e-prints*  
1653 (2013) arXiv:1306.2307 [[1306.2307](#)].
- 1654 [61] K. Madsen, R. Hickox, M. Bachetti, D. Stern, N.C. Gellert, J. García et al., *Hex-P: The High-Energy*  
1655 *X-ray Probe*, *Bulletin of the AAS* **51** (2019) .
- 1656 [62] V. Schonfelder, H. Aarts, K. Bennett, H. Deboer, J. Clear, W. Collmar et al., *Instrument description*  
1657 *and performance of the imaging gamma-ray telescope COMPTEL aboard the Compton Gamma-Ray*  
1658 *Observatory*, *The Astrophysical Journal Supplement Series* (1993) .
- 1659 [63] COSI collaboration, *The Compton Spectrometer and Imager Project for MeV Astronomy*, *PoS*  
1660 **ICRC2021** (2021) 652 [[2109.10403](#)].
- 1661 [64] H. Fleischhack, *AMEGO-X: MeV gamma-ray Astronomy in the Multi-messenger Era*, *PoS ICRC2021*  
1662 (2021) 649 [[2108.02860](#)].
- 1663 [65] E-ASTROGAM collaboration, *Science with e-ASTROGAM: A space mission for MeV–GeV gamma-*  
1664 *ray astrophysics*, *JHEAp* **19** (2018) 1 [[1711.01265](#)].
- 1665 [66] A. Takada, T. Takemura, K. Yoshikawa, Y. Mizumura, T. Ikeda, Y. Nakamura et al., *First observation*  
1666 *of mev gamma-ray universe with true imaging spectroscopy using the electron-tracking compton*  
1667 *telescope aboard smile-2+*, [2107.00180](#).
- 1668 [67] T. Tanimori et al., *MeV Gamma-ray imaging spectroscopic observation for Galactic Centre and Cosmic*  
1669 *Background MeV gammas by SMILE-2+ Balloon Experiment*, in *Journal of Physics: Conference*  
1670 *Series*, vol. 1468, p. 012046, IOP Publishing, 2020.
- 1671 [68] K. Hamaguchi, T. Tanimori, A. Takada, J.F. Beacom, S. Gunji, M. Mori et al., *A space-based all-sky*  
1672 *mev gamma-ray survey with the electron tracking compton camera*, [1907.06658](#).
- 1673 [69] A. Takada, T. Tanimori, Y. Mizumura, T. Takemura, K. Yoshikawa, Y. Nakamura et al., *Smile-3: sky*  
1674 *survey in mev gamma-ray using the electron-tracking compton telescope loaded on balloons*, in *Space*  
1675 *Telescopes and Instrumentation 2020: Ultraviolet to Gamma Ray*, vol. 11444, pp. 1017–1022, SPIE,  
1676 2020.

- 1677 [70] T. Aramaki, P. Hansson Adrian, G. Karagiorgi and H. Odaka, *Dual MeV Gamma-Ray and Dark*  
1678 *Matter Observatory - GRAMS Project*, *Astroparticle Physics* **114** (2020) 107 [[1901.03430](#)].
- 1679 [71] A. Moiseev, S. Profumo and A. Coogan, *Snowmass2021-Letter of Interest Searching for Dark Matter*  
1680 *and New Physics with GECCO*, .
- 1681 [72] A. Coogan, L. Morrison and S. Profumo, *Precision gamma-ray constraints for sub-GeV dark matter*  
1682 *models*, *JCAP* **08** (2021) 044 [[2104.06168](#)].
- 1683 [73] L. Roszkowski, E.M. Sessolo and S. Trojanowski, *WIMP dark matter candidates and searches—current*  
1684 *status and future prospects*, *Rept. Prog. Phys.* **81** (2018) 066201 [[1707.06277](#)].
- 1685 [74] G. Servant and T.M.P. Tait, *Is the lightest Kaluza-Klein particle a viable dark matter candidate?*,  
1686 *Nucl. Phys. B* **650** (2003) 391 [[hep-ph/0206071](#)].
- 1687 [75] J. Hubisz and P. Meade, *Phenomenology of the lightest Higgs with T-parity*, *Phys. Rev. D* **71** (2005)  
1688 [035016](#) [[hep-ph/0411264](#)].
- 1689 [76] M. Cirelli, N. Fornengo and A. Strumia, *Minimal dark matter*, *Nucl. Phys. B* **753** (2006) 178  
1690 [[hep-ph/0512090](#)].
- 1691 [77] M. Cahill-Rowley, R. Cotta, A. Drlica-Wagner, S. Funk, J. Hewett, A. Ismail et al., *Complementarity*  
1692 *of dark matter searches in the phenomenological MSSM*, *Phys. Rev. D* **91** (2015) 055011 [[1405.6716](#)].
- 1693 [78] I. Garcia Garcia, R. Lasenby and J. March-Russell, *Twin Higgs WIMP Dark Matter*, *Phys. Rev. D*  
1694 **92** (2015) 055034 [[1505.07109](#)].
- 1695 [79] R.K. Leane, T.R. Slatyer, J.F. Beacom and K.C.Y. Ng, *GeV-scale thermal WIMPs: Not even slightly*  
1696 *ruled out*, *Phys. Rev. D* **98** (2018) 023016 [[1805.10305](#)].
- 1697 [80] J. Arakawa and T.M.P. Tait, *Is a Miracle-less WIMP Ruled out?*, *SciPost Phys.* **11** (2021) 019  
1698 [[2101.11031](#)].
- 1699 [81] J. Goodman, M. Ibe, A. Rajaraman, W. Shepherd, T.M.P. Tait and H.-B. Yu, *Constraints on Dark*  
1700 *Matter from Colliders*, *Phys. Rev. D* **82** (2010) 116010 [[1008.1783](#)].
- 1701 [82] A.L. Fitzpatrick, W. Haxton, E. Katz, N. Lubbers and Y. Xu, *The Effective Field Theory of Dark*  
1702 *Matter Direct Detection*, *JCAP* **02** (2013) 004 [[1203.3542](#)].
- 1703 [83] M.I. Gresham and K.M. Zurek, *Effect of nuclear response functions in dark matter direct detection*,  
1704 *Phys. Rev. D* **89** (2014) 123521 [[1401.3739](#)].
- 1705 [84] R.J. Hill and M.P. Solon, *Standard Model anatomy of WIMP dark matter direct detection II: QCD*  
1706 *analysis and hadronic matrix elements*, *Phys. Rev. D* **91** (2015) 043505 [[1409.8290](#)].
- 1707 [85] A.K. Drukier, K. Freese and D.N. Spergel, *Detecting Cold Dark Matter Candidates*, *Phys. Rev. D* **33**  
1708 (1986) 3495.
- 1709 [86] S.E. Vahsen, C.A.J. O'Hare and D. Loomba, *Directional Recoil Detection*, *Ann. Rev. Nucl. Part. Sci.*  
1710 **71** (2021) 189 [[2102.04596](#)].
- 1711 [87] D.S. Akerib et al., *Snowmass2021 Cosmic Frontier Dark Matter Direct Detection to the Neutrino Fog,*  
1712 *in 2022 Snowmass Summer Study*, 3, 2022 [[2203.08084](#)].
- 1713 [88] E. Alfonso-Pita et al., *Snowmass 2021 Scintillating Bubble Chambers: Liquid-noble Bubble Chambers*  
1714 *for Dark Matter and CEνNS Detection*, in *2022 Snowmass Summer Study*, 7, 2022 [[2207.12400](#)].



- 1715 [89] XENON collaboration, *Search for Coherent Elastic Scattering of Solar  $^8B$  Neutrinos in the XENON1T*  
1716 *Dark Matter Experiment*, *Phys. Rev. Lett.* **126** (2021) 091301 [2012.02846].
- 1717 [90] SUPERCDMS collaboration, *Search for Low-Mass Dark Matter with CDMSlite Using a Profile*  
1718 *Likelihood Fit*, *Phys. Rev. D* **99** (2019) 062001 [1808.09098].
- 1719 [91] G. Adhikari et al., *An experiment to search for dark-matter interactions using sodium iodide detectors*,  
1720 *Nature* **564** (2018) 83 [1906.01791].
- 1721 [92] CRESST collaboration, *First results from the CRESST-III low-mass dark matter program*, *Phys.*  
1722 *Rev. D* **100** (2019) 102002 [1904.00498].
- 1723 [93] R. Bernabei et al., *First model independent results from DAMA/LIBRA-phase2*, *Universe* **4** (2018)  
1724 116.
- 1725 [94] DARKSIDE collaboration, *Low-Mass Dark Matter Search with the DarkSide-50 Experiment*, *Phys. Rev.*  
1726 *Lett.* **121** (2018) 081307 [1802.06994].
- 1727 [95] DEAP collaboration, *Search for dark matter with a 231-day exposure of liquid argon using DEAP-3600*  
1728 *at SNOLAB*, *Phys. Rev. D* **100** (2019) 022004 [1902.04048].
- 1729 [96] EDELWEISS collaboration, *Improved EDELWEISS-III sensitivity for low-mass WIMPs using a*  
1730 *profile likelihood approach*, *Eur. Phys. J. C* **76** (2016) 548 [1607.03367].
- 1731 [97] LUX collaboration, *Results from a search for dark matter in the complete LUX exposure*, *Phys. Rev.*  
1732 *Lett.* **118** (2017) 021303 [1608.07648].
- 1733 [98] NEWS-G collaboration, *First results from the NEWS-G direct dark matter search experiment at the*  
1734 *LSM*, *Astropart. Phys.* **97** (2018) 54 [1706.04934].
- 1735 [99] PANDAX-II collaboration, *Dark Matter Results From 54-Ton-Day Exposure of PandaX-II Experiment*,  
1736 *Phys. Rev. Lett.* **119** (2017) 181302 [1708.06917].
- 1737 [100] PICO collaboration, *Improved dark matter search results from PICO-2L Run 2*, *Phys. Rev. D* **93**  
1738 (2016) 061101 [1601.03729].
- 1739 [101] PICO collaboration, *Dark Matter Search Results from the PICO-60 C<sub>3</sub>F<sub>8</sub> Bubble Chamber*, *Phys.*  
1740 *Rev. Lett.* **118** (2017) 251301 [1702.07666].
- 1741 [102] XENON collaboration, *Search for Light Dark Matter Interactions Enhanced by the Migdal Effect or*  
1742 *Bremsstrahlung in XENON1T*, *Phys. Rev. Lett.* **123** (2019) 241803 [1907.12771].
- 1743 [103] PANDAX-4T collaboration, *Dark Matter Search Results from the PandaX-4T Commissioning Run*,  
1744 *Phys. Rev. Lett.* **127** (2021) 261802 [2107.13438].
- 1745 [104] LZ collaboration, *First Dark Matter Search Results from the LUX-ZEPLIN (LZ) Experiment*,  
1746 2207.03764.
- 1747 [105] C.A.J. O'Hare, *New Definition of the Neutrino Floor for Direct Dark Matter Searches*, *Phys. Rev.*  
1748 *Lett.* **127** (2021) 251802 [2109.03116].
- 1749 [106] P. Cushman et al., *Working Group Report: WIMP Dark Matter Direct Detection*, in *Community*  
1750 *Summer Study 2013: Snowmass on the Mississippi*, 10, 2013 [1310.8327].
- 1751 [107] COUPP COLLABORATION collaboration, *First Dark Matter Search Results from a 4-kg CF<sub>3</sub>I Bubble*  
1752 *Chamber Operated in a Deep Underground Site*, *Phys. Rev. D* **86** (2012) 052001 [1204.3094].

- 1753 [108] KIMS collaboration, *Limits on Interactions between Weakly Interacting Massive Particles and*  
1754 *Nucleons Obtained with NaI(Tl) crystal Detectors*, *JHEP* **03** (2019) 194 [[1806.06499](#)].
- 1755 [109] PICASSO COLLABORATION collaboration, *Constraints on Low-Mass WIMP Interactions on  $^{19}\text{F}$*   
1756 *from PICASSO*, *Phys. Lett. B* **711** (2012) 153 [[1202.1240](#)].
- 1757 [110] PICO COLLABORATION collaboration, *Improved dark matter search results from PICO-2L Run 2*,  
1758 *Phys. Rev. D* **93** (2016) 061101 [[1601.03729](#)].
- 1759 [111] PICO collaboration, *Dark Matter Search Results from the PICO-60  $\text{C}_3\text{F}_8$  Bubble Chamber*, *Phys.*  
1760 *Rev. Lett.* **118** (2017) 251301 [[1702.07666](#)].
- 1761 [112] CDMS-II collaboration, *Dark Matter Search Results from the CDMS II Experiment*, *Science* **327**  
1762 (2010) 1619 [[0912.3592](#)].
- 1763 [113] LUX collaboration, *Limits on spin-dependent WIMP-nucleon cross section obtained from the complete*  
1764 *LUX exposure*, *Phys. Rev. Lett.* **118** (2017) 251302 [[1705.03380](#)].
- 1765 [114] PANDAX-II collaboration, *PandaX-II Constraints on Spin-Dependent WIMP-Nucleon Effective*  
1766 *Interactions*, *Phys. Lett. B* **792** (2019) 193 [[1807.01936](#)].
- 1767 [115] XENON100 collaboration, *XENON100 Dark Matter Results from a Combination of 477 Live Days*,  
1768 *Phys. Rev. D* **94** (2016) 122001 [[1609.06154](#)].
- 1769 [116] XENON collaboration, *Search for Light Dark Matter Interactions Enhanced by the Migdal Effect or*  
1770 *Bremsstrahlung in XENON1T*, *Phys. Rev. Lett.* **123** (2019) 241803 [[1907.12771](#)].
- 1771 [117] XENON collaboration, *Constraining the spin-dependent WIMP-nucleon cross sections with*  
1772 *XENON1T*, *Phys. Rev. Lett.* **122** (2019) 141301 [[1902.03234](#)].
- 1773 [118] D. Carney et al., *Snowmass2021 Cosmic Frontier White Paper: Ultraheavy particle dark matter*,  
1774 [2203.06508](#).
- 1775 [119] R.K. Leane et al., *Snowmass2021 Cosmic Frontier White Paper: Puzzling Excesses in Dark Matter*  
1776 *Searches and How to Resolve Them*, [2203.06859](#).
- 1777 [120] S. Ando et al., *Snowmass2021 Cosmic Frontier: Synergies between dark matter searches and*  
1778 *multiwavelength/multimessenger astrophysics*, in *2022 Snowmass Summer Study*, 3, 2022 [[2203.06781](#)].
- 1779 [121] FERMI-LAT, DES collaboration, *Searching for Dark Matter Annihilation in Recently Discovered*  
1780 *Milky Way Satellites with Fermi-LAT*, *Astrophys. J.* **834** (2017) 110 [[1611.03184](#)].
- 1781 [122] S. Ando, A. Geringer-Sameth, N. Hiroshima, S. Hoof, R. Trotta and M.G. Walker, *Structure formation*  
1782 *models weaken limits on WIMP dark matter from dwarf spheroidal galaxies*, *Phys. Rev. D* **102** (2020)  
1783 [061302](#) [[2002.11956](#)].
- 1784 [123] J.R. Hargis, B. Willman and A.H.G. Peter, *Too Many, Too Few, or Just Right? The Predicted*  
1785 *Number and Distribution of Milky Way Dwarf Galaxies*, *The Astrophysical Journal* **795** (2014) L13  
1786 [[1407.4470](#)].
- 1787 [124] P. Abreu et al., *The Southern Wide-Field Gamma-Ray Observatory (SWG0): A Next-Generation*  
1788 *Ground-Based Survey Instrument for VHE Gamma-Ray Astronomy*, [1907.07737](#).
- 1789 [125] Cherenkov Telescope Array Consortium, B.S. Acharya, I. Agudo, I. Al Samarai, R. Alfaro, J. Alfaro  
1790 et al., *Science with the Cherenkov Telescope Array* (2019), [10.1142/10986](#).



- 1791 [126] CTA collaboration, *Sensitivity of the Cherenkov Telescope Array to a dark matter signal from the*  
1792 *Galactic centre*, *Journal of Cosmology and Astroparticle Physics* **01** (2021) 057 [2007.16129].
- 1793 [127] A. Viana, H. Schoorlemmer, A. Albert, V. de Souza, J.P. Harding and J. Hinton, *Searching for Dark*  
1794 *Matter in the Galactic Halo with a Wide Field of View TeV Gamma-ray Observatory in the Southern*  
1795 *Hemisphere*, *Journal of Cosmology and Astroparticle Physics* **12** (2019) 061 [1906.03353].
- 1796 [128] J. Buckley, L. Bergström, W. Binns, J. Buhler, W. Chen, S. Cherry et al., *Astro2020 APC White*  
1797 *Paper: The Advanced Particle-astrophysics Telescope (APT)*, .
- 1798 [129] T. Aramaki, C.J. Hailey, S.E. Boggs, P. von Doetinchem, H. Fuke, S.I. Mognet et al., *Antideuteron*  
1799 *sensitivity for the GAPS experiment*, *Astroparticle Physics* **74** (2016) 6 [1506.02513].
- 1800 [130] P. Allison, J. Beatty, L. Beaufore, Y. Chen, S. Coutu, E. Ellingwood et al., *Cosmic-ray isotope*  
1801 *measurements with helix*, *Proceedings of Science* **358** (2019) .
- 1802 [131] R. Battiston et al., *High precision particle astrophysics as a new window on the universe with an*  
1803 *Antimatter Large Acceptance Detector In Orbit (ALADInO)*, *Exper. Astron.* **51** (2021) 1299.
- 1804 [132] O. Adriani, C. Altomare, G. Ambrosi, P. Azzarello, F.C.T. Barbato, R. Battiston et al., *Design of an*  
1805 *Antimatter Large Acceptance Detector In Orbit (ALADInO)*, *Instruments* **6** (2022) .
- 1806 [133] S. Schael, A. Atanasyan, J. Berdugo, T. Bretz, M. Czapalla, B. Dachwald et al., *AMS-100: The*  
1807 *next generation magnetic spectrometer in space - An international science platform for physics and*  
1808 *astrophysics at Lagrange point 2*, *Nuclear Instruments and Methods in Physics Research A* **944** (2019)  
1809 162561 [1907.04168].
- 1810 [134] A. Aduszkiewicz et al., *Measurements of  $\pi^+$ ,  $K^+$ ,  $p$  and  $p\bar{p}$  spectra in proton-proton interactions*  
1811 *at 20, 31, 40, 80 and 158 GeV/s with the NA61/SHINE spectrometer at the CERN SPS*, *European*  
1812 *Physical Journal C* **77** (2017) 671 [1705.02467].
- 1813 [135] ALICE collaboration, *Measurement of the low-energy antideuteron inelastic cross section*, *Physical*  
1814 *Review Letters* **125** (2020) 162001 [2005.11122].
- 1815 [136] R. Aaij et al., *Measurement of Antiproton Production in pHe Collisions at  $\sqrt{s_{NN}} = 110$  GeV*, *Physical*  
1816 *Review Letters* **121** (2018) 222001 [1808.06127].
- 1817 [137] B. Adams et al., *Letter of Intent: A New QCD facility at the M2 beam line of the CERN SPS*  
1818 *(COMPASS++/AMBER)*, 1808.00848.
- 1819 [138] L.A. Anchordoqui et al., *Hunting super-heavy dark matter with ultra-high energy photons*, *Astropart.*  
1820 *Phys.* **132** (2021) 102614 [2105.12895].
- 1821 [139] E.D. Hall, R.X. Adhikari, V.V. Frolov, H. Müller, M. Pospelov and R.X. Adhikari, *Laser*  
1822 *Interferometers as Dark Matter Detectors*, *Phys. Rev. D* **98** (2018) 083019 [1605.01103].
- 1823 [140] S. Baum, M.A. Fedderke and P.W. Graham, *Searching for dark clumps with gravitational-wave*  
1824 *detectors*, *Phys. Rev. D* **106** (2022) 063015 [2206.14832].
- 1825 [141] S. Bird et al., *Snowmass2021 Cosmic Frontier White Paper: Primordial Black Hole Dark Matter*,  
1826 2203.08967.
- 1827 [142] K.N. Abazajian et al., *Synergy between cosmological and laboratory searches in neutrino physics: a*  
1828 *white paper*, 2203.07377.

- 1829 [143] M. Ackermann et al., *High-energy and ultra-high-energy neutrinos: A Snowmass white paper*, *JHEAp*  
1830 **36** (2022) 55 [2203.08096].
- 1831 [144] C.A. Argüelles et al., *Snowmass White Paper: Beyond the Standard Model effects on Neutrino Flavor*,  
1832 in *2022 Snowmass Summer Study*, 3, 2022 [2203.10811].
- 1833 [145] L.A. Anchordoqui et al., *The Forward Physics Facility: Sites, experiments, and physics potential*,  
1834 *Phys. Rept.* **968** (2022) 1 [2109.10905].
- 1835 [146] J.L. Feng et al., *The Forward Physics Facility at the High-Luminosity LHC*, 2203.05090.
- 1836 [147] K.V. Berghaus, P.W. Graham, D.E. Kaplan, G.D. Moore and S. Rajendran, *Dark energy radiation*,  
1837 *Phys. Rev. D* **104** (2021) 083520 [2012.10549].
- 1838 [148] PIERRE AUGER collaboration, *Testing Hadronic Interactions at Ultrahigh Energies with Air Showers*  
1839 *Measured by the Pierre Auger Observatory*, *Phys. Rev. Lett.* **117** (2016) 192001 [1610.08509].
- 1840 [149] A. Coleman et al., *Ultra-High-Energy Cosmic Rays: The Intersection of the Cosmic and Energy*  
1841 *Frontiers*, 2205.05845.
- 1842 [150] K. Engel, T. Lewis, M.S. Muzio and T.M. Venters, *Advancing the Landscape of Multimessenger*  
1843 *Science in the Next Decade*, in *2022 Snowmass Summer Study*, 3, 2022 [2203.10074].
- 1844 [151] J.A. Dror, T. Hiramatsu, K. Kohri, H. Murayama and G. White, *Testing the Seesaw Mechanism and*  
1845 *Leptogenesis with Gravitational Waves*, *Phys. Rev. Lett.* **124** (2020) 041804 [1908.03227].